

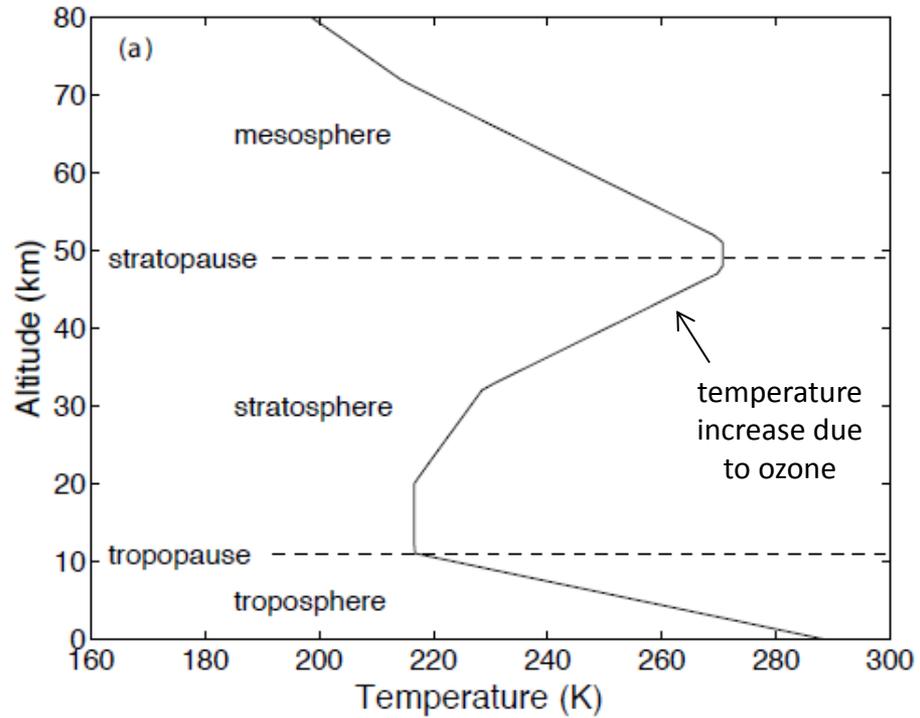
## Overview of this week's lectures:

- Global circulation and transport
- Satellite observations of stratospheric temperature and water vapor
- Global upper troposphere - lower stratosphere (UTLS)
- Circulation and transport near the tropical tropopause layer (TTL)
- UTLS monsoon circulations and water vapor isotopes
- Research seminar: tropical dynamics with GPS radio occultation data

## Lecture 1: Global atmospheric circulation and satellite observations

- large-scale circulation of the troposphere and stratosphere
  - Climatology and variability of the stratosphere
  - wave forcing of the zonal mean flow
  - Rossby waves: propagation and dissipation (critical layers)
  - Tropospheric baroclinic wave life cycles
  - large-scale tropical circulations
  - QBO and ENSO
- Large-scale transport

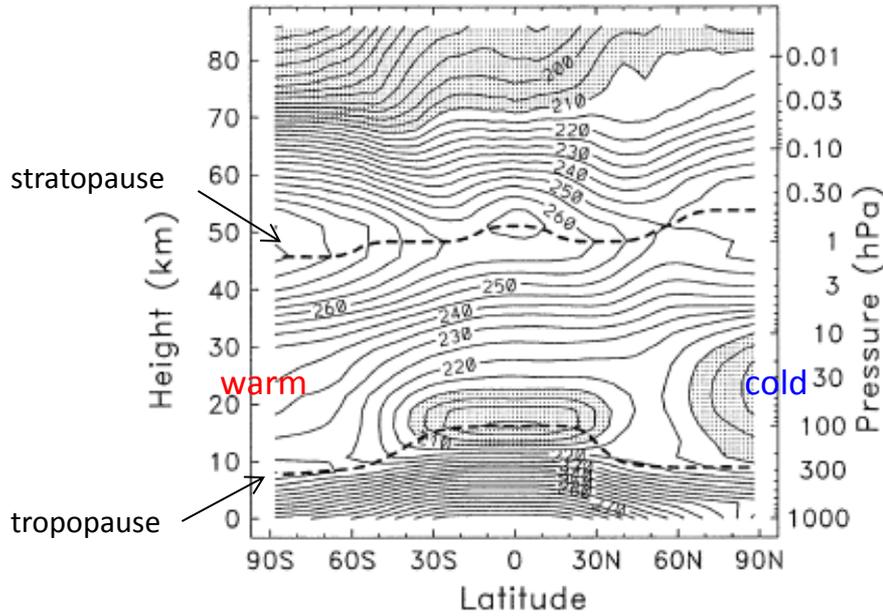
# The Standard Atmosphere



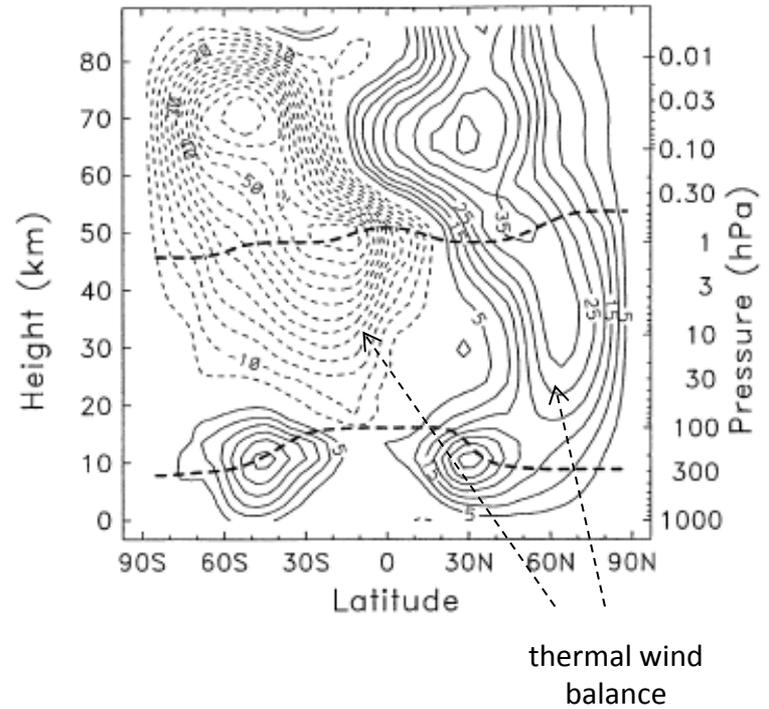
# Climatological temperatures and zonal winds in January



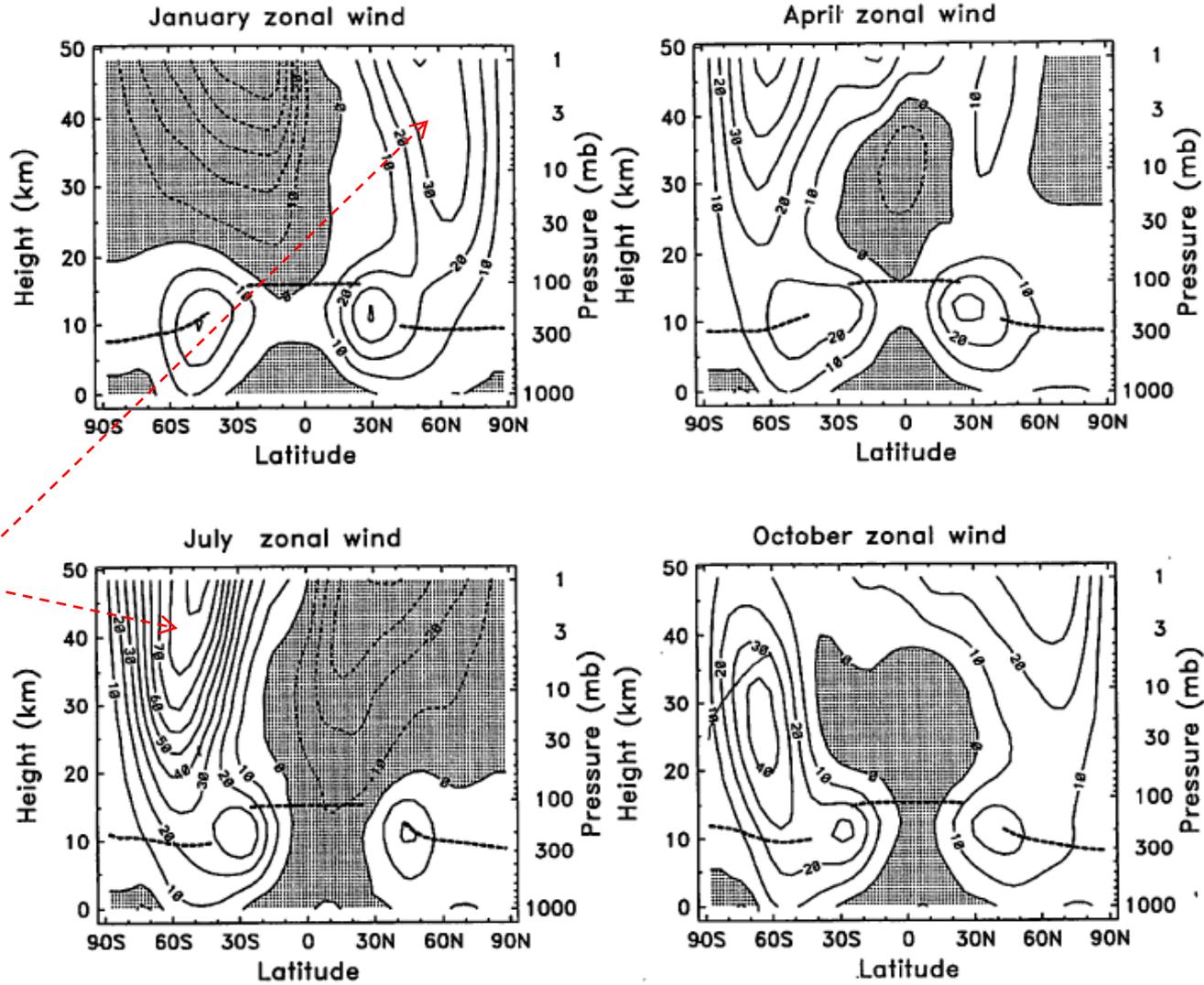
January temperature



January zonal wind

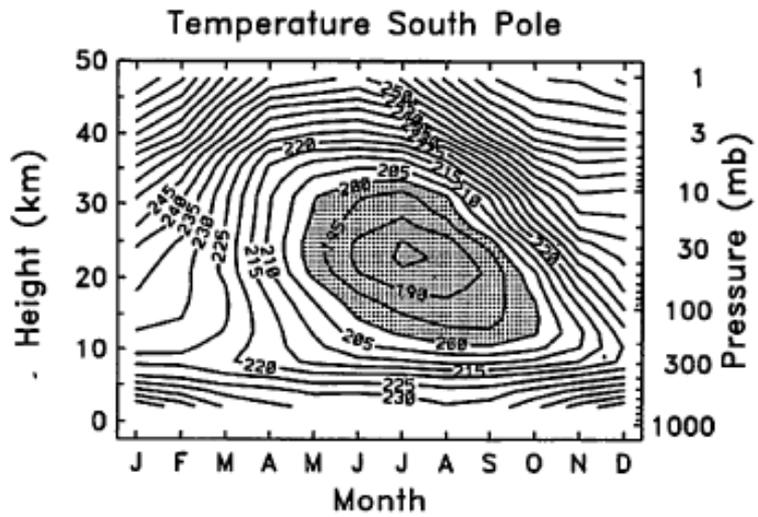
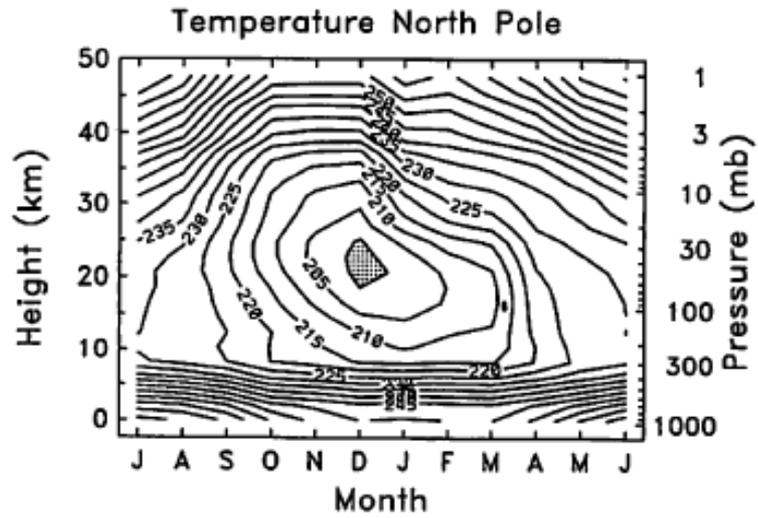


# Seasonal cycle of zonal mean zonal winds

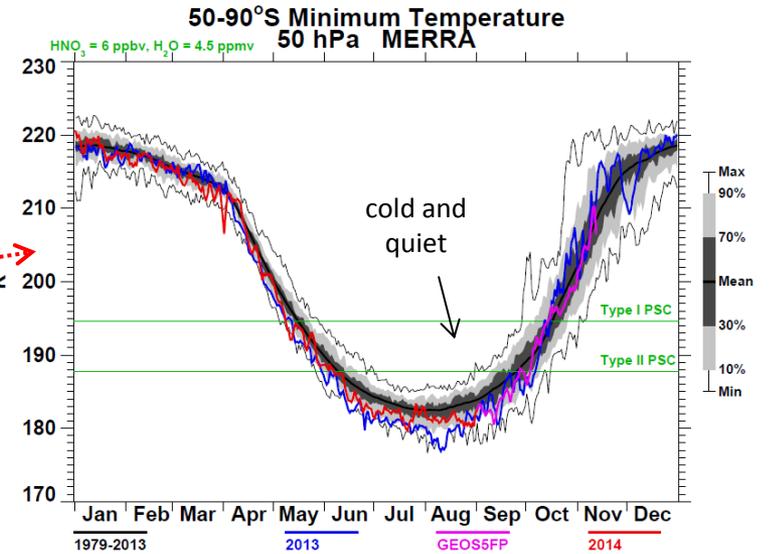
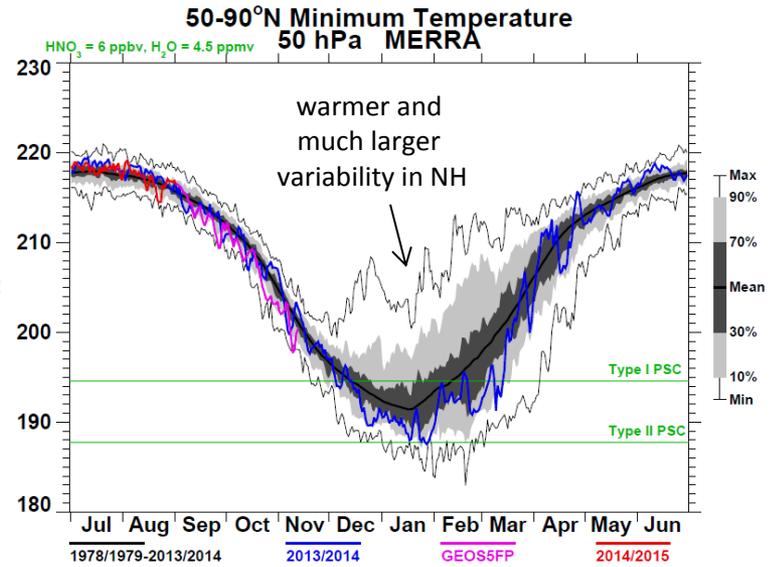
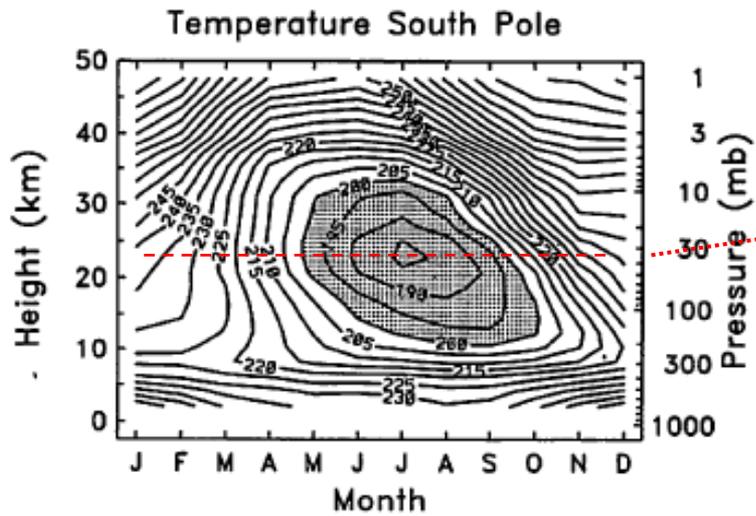
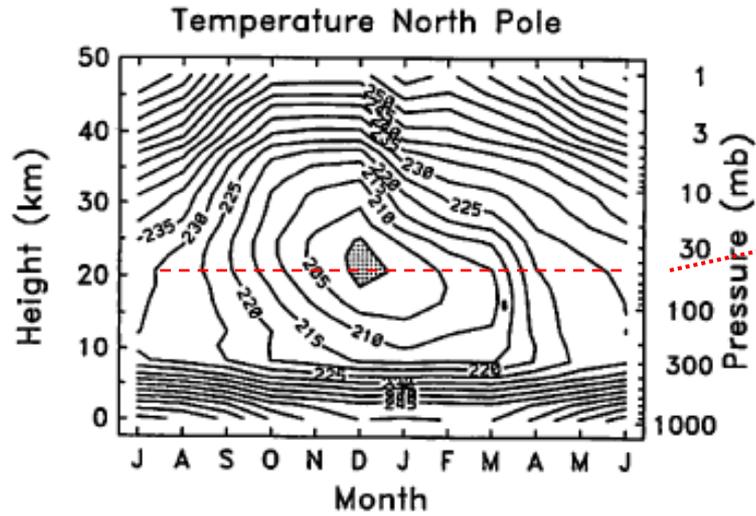


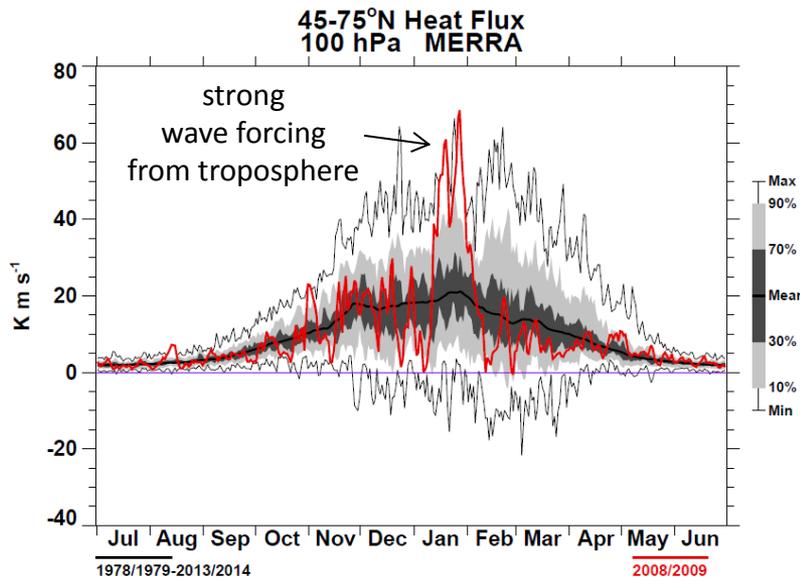
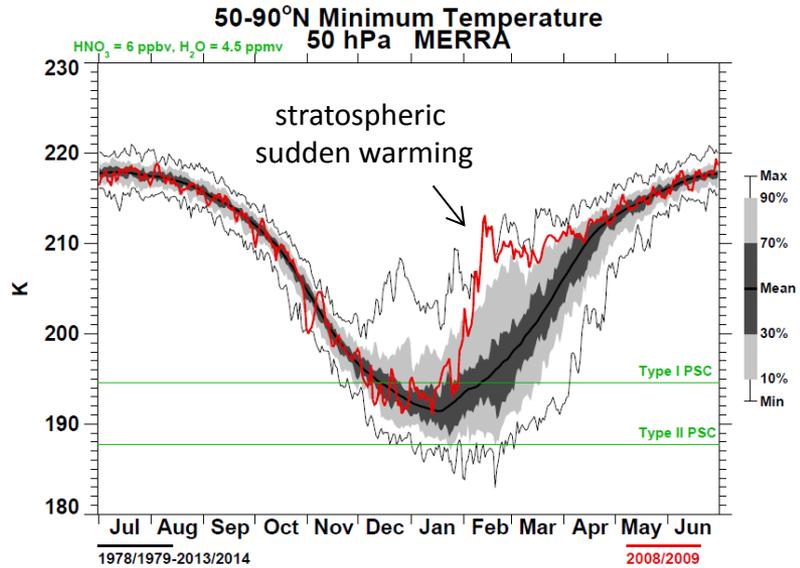
polar night jet  
stronger in SH  
than in NH

# climatological polar temperature

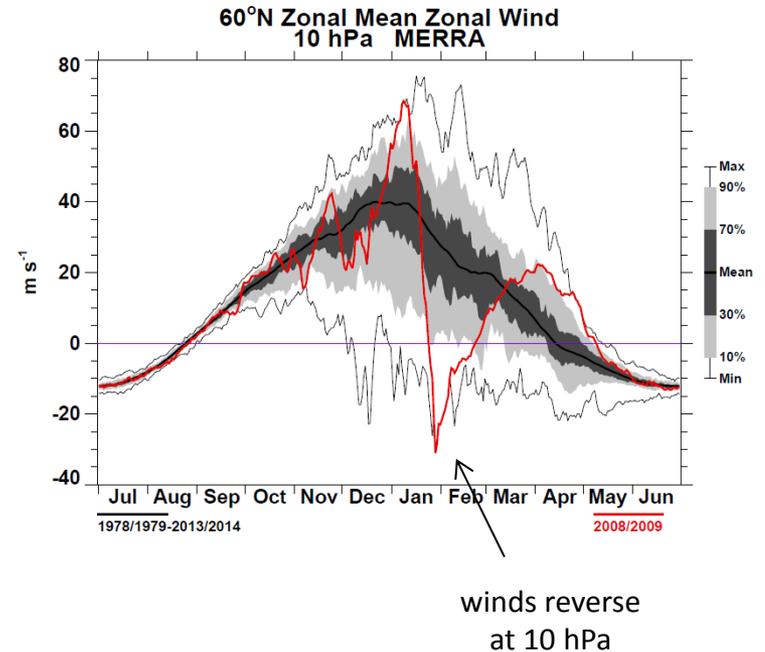


# climatological polar temperature





- Variability in NH winter stratosphere tied to large-scale forcing from troposphere.
- Episodic forcing produces 'stratospheric sudden warming' events.
- Largest observed stratosphere sudden warming in January 2009



# A Major Stratospheric Sudden Warming Event in January 2009

YAYOI HARADA, ATSUSHI GOTO, HIROSHI HASEGAWA, AND NORIHISA FUJIKAWA

Climate Prediction Division, Japan Meteorological Agency, Tokyo, Japan

JAS 2010

EP flux



(I) 18-20Jan. 2009, U & EPF

(Z at 10 hPa)

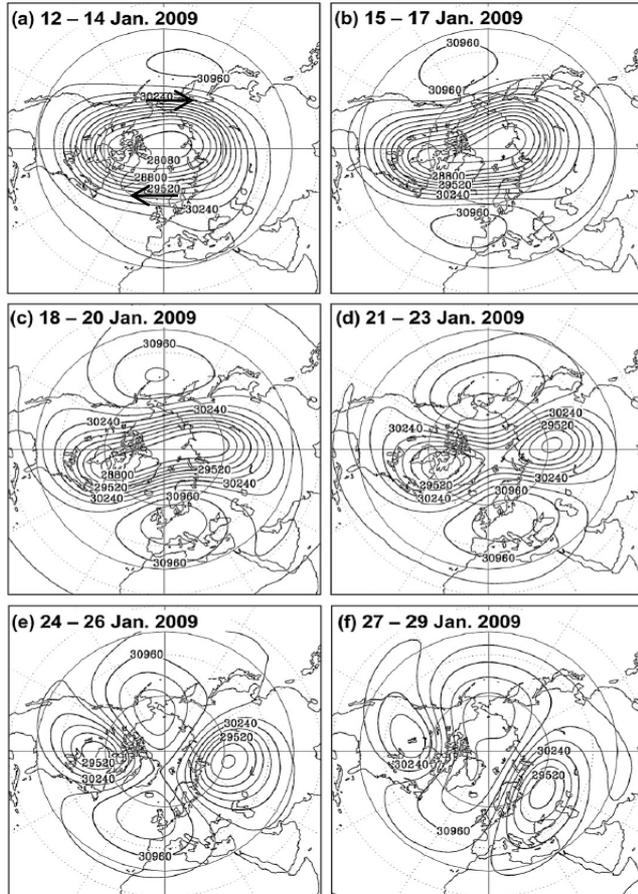
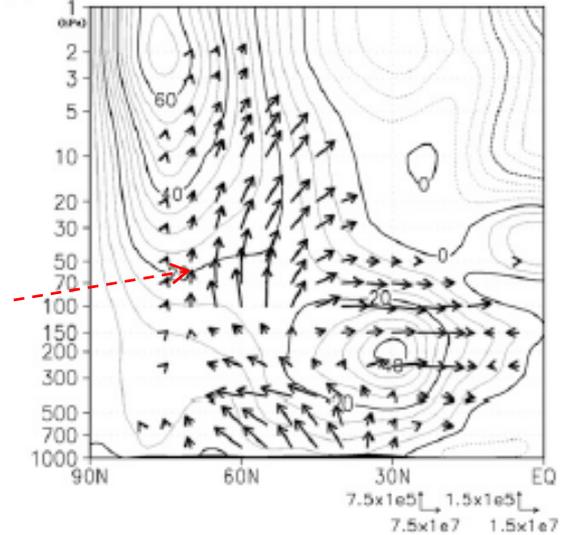


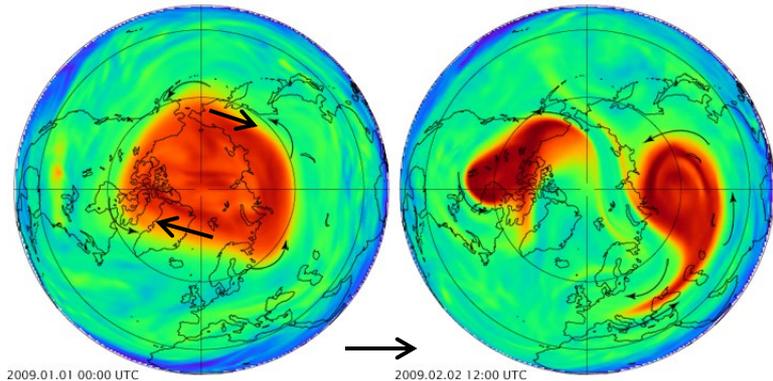
FIG. 3. The 10-hPa geopotential heights for six successive 3-day means in January 2009. The contour interval is 240 m.

wave forcing from troposphere



polar vortex near 30 km

potential vorticity



split of polar vortex

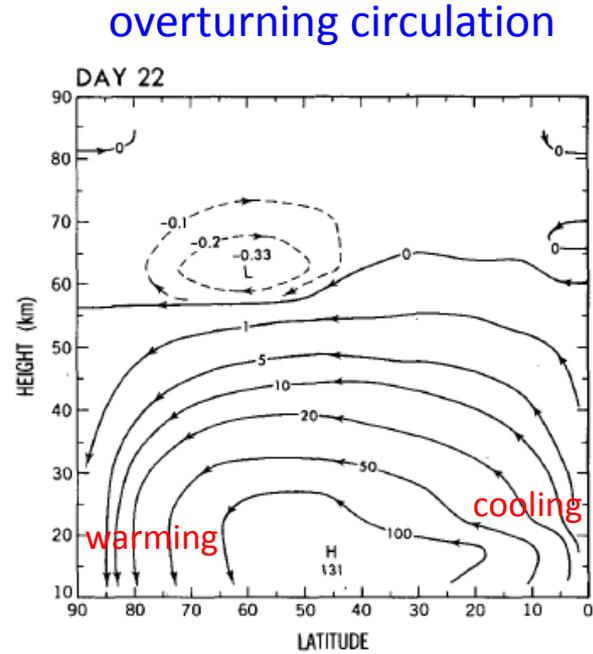
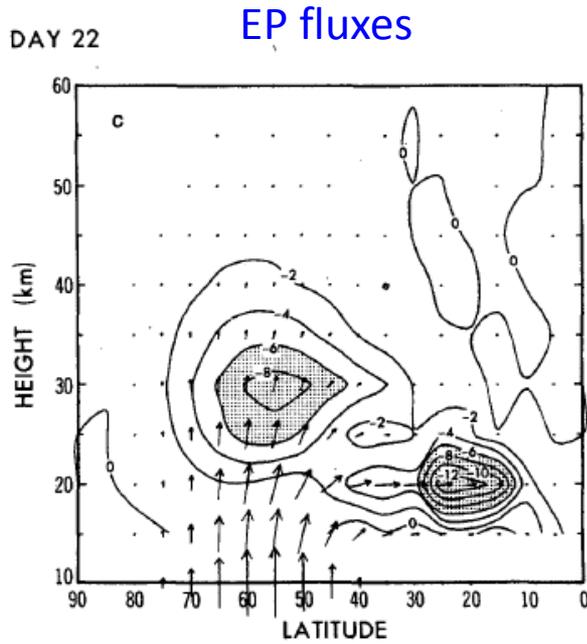
# A Dynamical Model of the Stratospheric Sudden Warming

TAROH MATSUNO<sup>1</sup>

*Geophysical Fluid Dynamics Laboratory, NOAA, Princeton University, Princeton, N. J.*

(Manuscript received 29 March 1971, in revised form 16 August 1971)

← solution to puzzle of stratospheric warmings



## Some Eulerian and Lagrangian Diagnostics for a Model Stratospheric Warming<sup>1</sup>

T. DUNKERTON, C.-P. F. HSU<sup>2</sup> AND M. E. MCINTYRE<sup>3</sup>

*Department of Atmospheric Sciences, University of Washington, Seattle 98195*

(Manuscript received 30 May 1980, in final form 11 December 1980)

Governing equations for the zonal mean flow (Transformed Eulerian mean)

EP flux divergence (wave forcing)

zonal momentum balance

$$\frac{\partial \bar{u}}{\partial t} - \hat{f} \bar{v}^* = \text{DF}$$

thermodynamic balance

$$\frac{\partial \bar{T}}{\partial t} + \bar{v}^* \frac{1}{a} \frac{\partial \bar{T}}{\partial \phi} + \bar{w}^* S = \bar{Q},$$

continuity equation

$$(a \cos \phi)^{-1} \frac{\partial}{\partial \phi} (\bar{v}^* \cos \phi) + e^{z/H} \frac{\partial}{\partial z} (\bar{w}^* e^{-z/H}) = 0,$$

geostrophic thermal wind

$$f \frac{\partial \bar{u}}{\partial z} + \frac{R}{aH} \frac{\partial \bar{T}}{\partial \phi} = 0.$$

Andrews et al, 1987

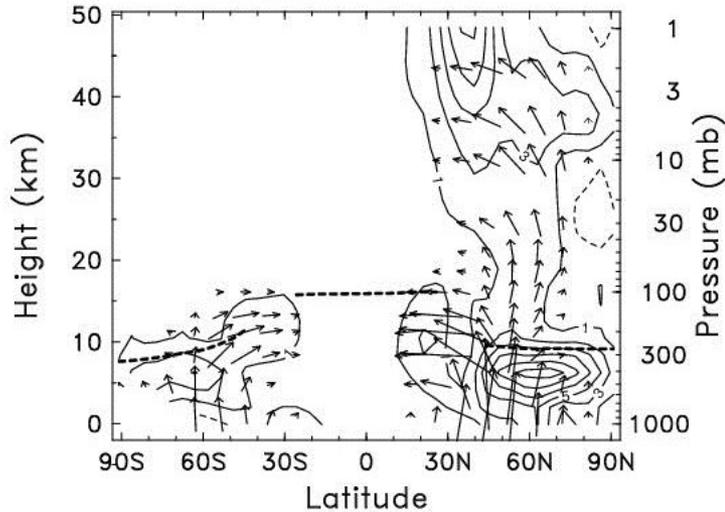
Eliassen-Palm fluxes:

EP flux divergence (wave forcing)

$$\frac{\partial \bar{u}}{\partial t} - \hat{f} \bar{v}^* = DF \quad DF = \frac{\exp(z/H)}{a \cos \phi} \nabla \cdot \mathbf{F}$$

climatology

January EP flux



components:

latitudinal flux

$$F_{\phi} = \exp(-z/H) a \cos \phi \left[ -\overline{u'v'} \right]$$

momentum flux

vertical flux

$$F_z = \exp(-z/H) a \cos \phi \left[ \overline{v'T'} \right]$$

heat flux

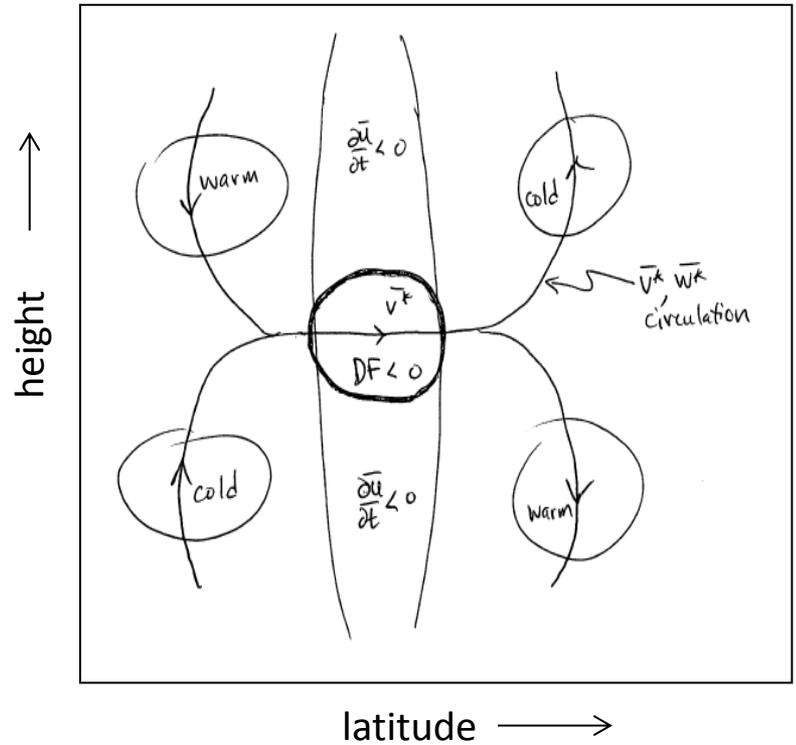
Important points:

- DF quantifies zonal momentum forcing
- $\mathbf{F}$  proportional to 'wave activity' flux (DF shows sources and sinks of waves)
- $F_{\phi}$  and  $F_z$  indicate direction of wave propagation

# Response of a balanced vortex to localized EP flux forcing (DF)

$$\frac{\partial \bar{u}}{\partial t} - f \bar{v}^* = DF$$

- response is balanced between  $\frac{\partial \bar{u}}{\partial t}$  and  $f \bar{v}^*$
- $\bar{v}^*$ ,  $\bar{w}^*$  and  $\frac{\partial \bar{u}}{\partial t}$  act to extend DF forcing non-locally
- overall circulation maintains thermal wind balance



# Circulation response depends on frequency of forcing:

Combine equations:

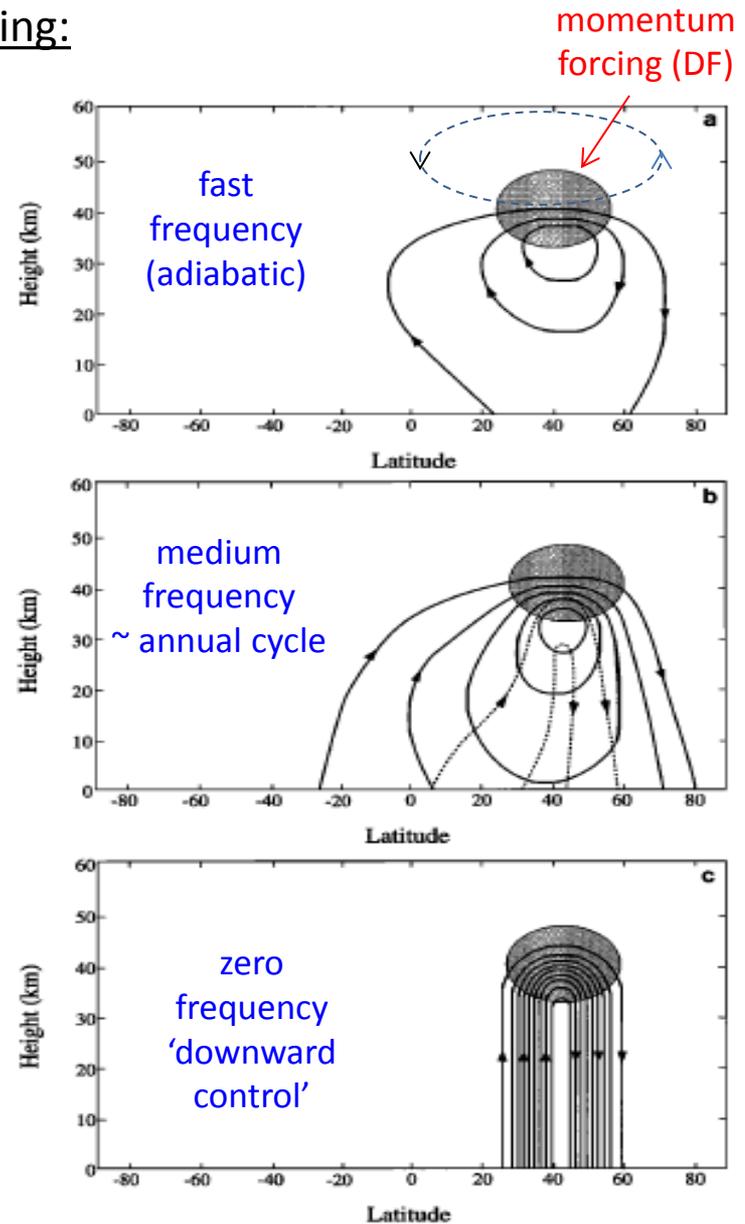
$$\frac{\partial}{\partial z} \left( \frac{1}{\rho_0} \frac{\partial(\rho_0 \hat{w})}{\partial z} \right) \quad \text{time dependence}$$

$$+ \left( \frac{i\sigma}{i\sigma + \alpha} \right) \frac{N^2}{4\Omega^2 a^2 \cos \phi} \frac{\partial}{\partial \phi} \left( \frac{\cos \phi}{\sin^2 \phi} \frac{\partial \hat{w}}{\partial \phi} \right)$$

$$= \left( \frac{i\sigma}{i\sigma + \alpha} \right) \frac{(R/H)}{4\Omega^2 a^2 \cos \phi} \frac{\partial}{\partial \phi} \left( \frac{\cos \phi}{\sin^2 \phi} \frac{\partial \hat{Q}}{\partial \phi} \right) \quad \leftarrow \text{diabatic heating}$$

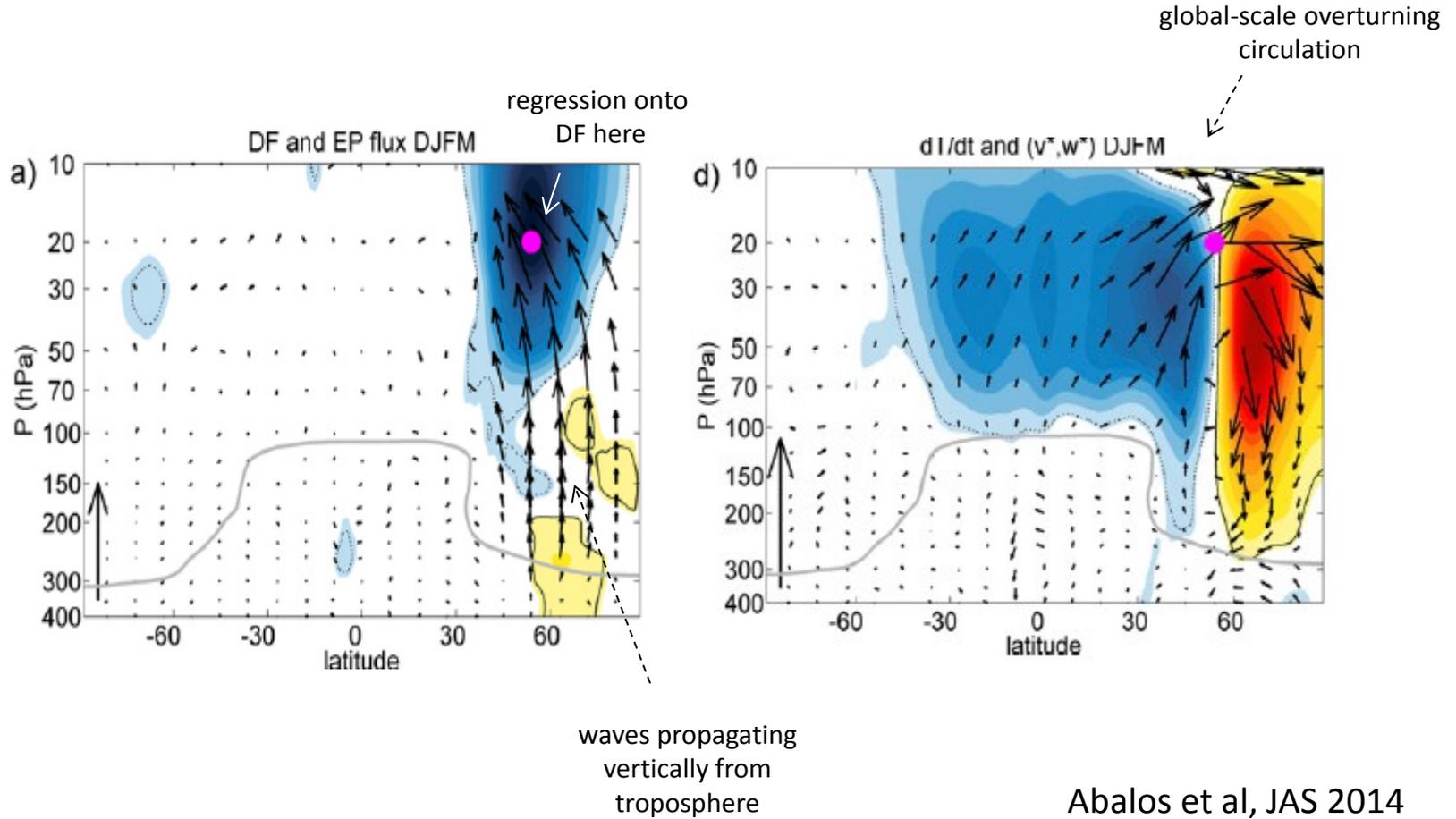
$$+ \frac{1}{2\Omega a \cos \phi} \frac{\partial}{\partial \phi} \left( \frac{\cos \phi}{\sin \phi} \frac{\partial \hat{G}}{\partial z} \right) \quad \leftarrow \text{momentum forcing (DF)}$$

In general both Q and DF drive the mean circulation. These plots show the response to isolated forcing from Rossby wave EP flux divergence. The lower cell becomes more important for slower forcing.

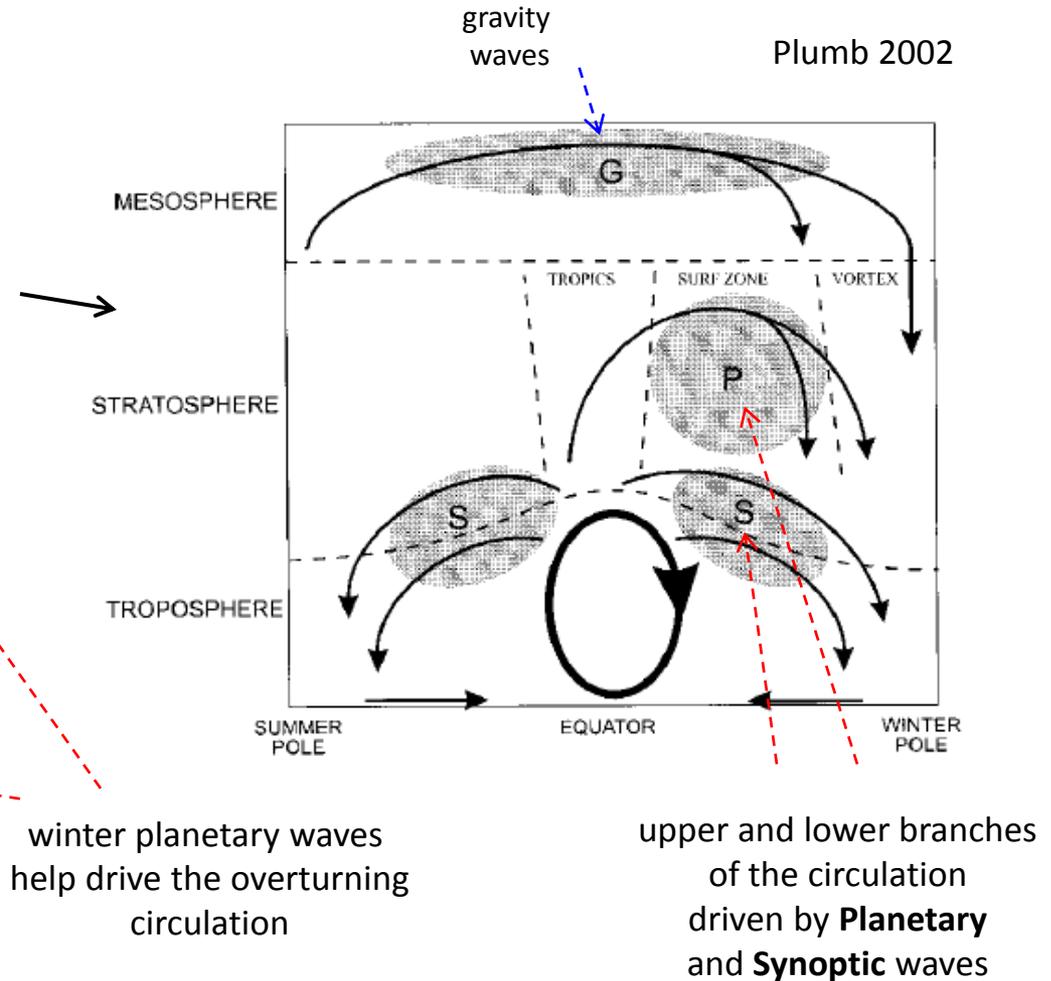
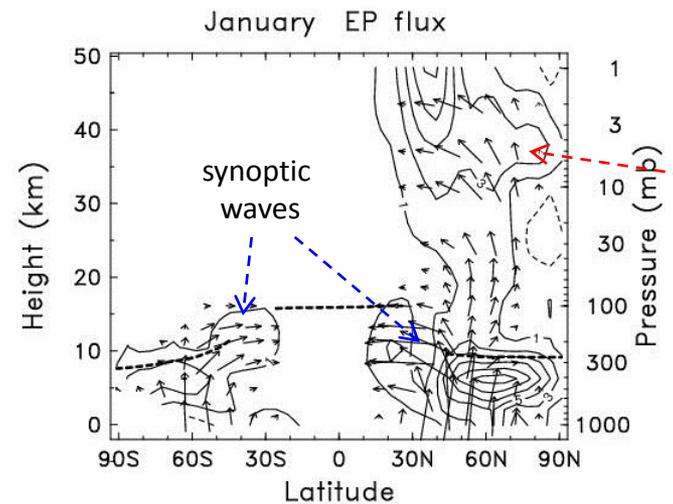
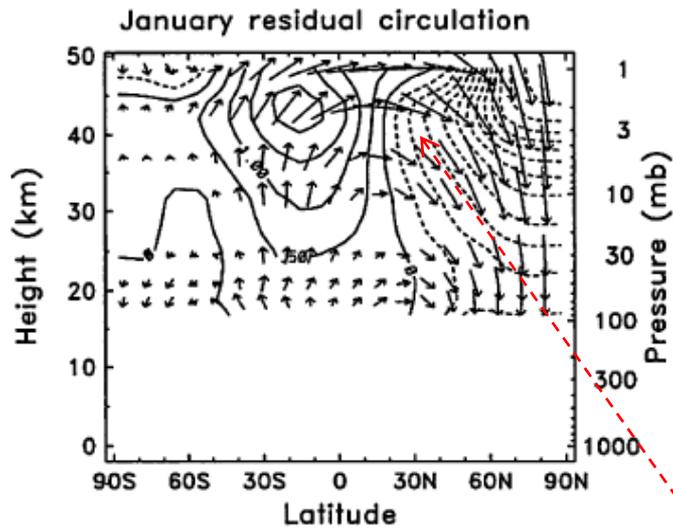


Haynes et al 1991  
Holton et al 1995

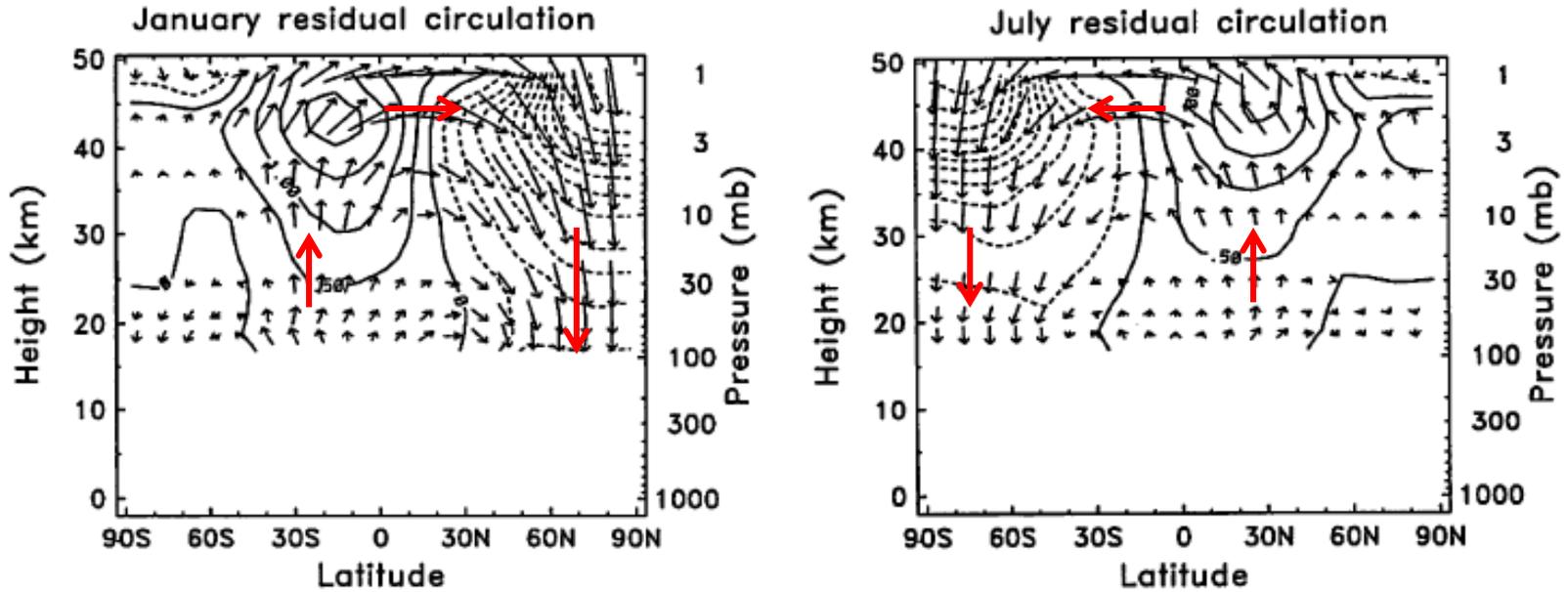
# Wave – mean flow patterns from observations (reanalysis data)



# Climatology of stratospheric overturning circulation



The overturning circulation reverses between solstice seasons



↑  
circulation is stronger during NH winter, related to stronger wave forcing from troposphere

The stratospheric overturning circulation is often termed the Brewer-Dobson circulation (closely related to the Lagrangian or transport circulation)

deduced by Brewer (1949) studying stratospheric water vapor and Dobson (1956) studying stratospheric ozone

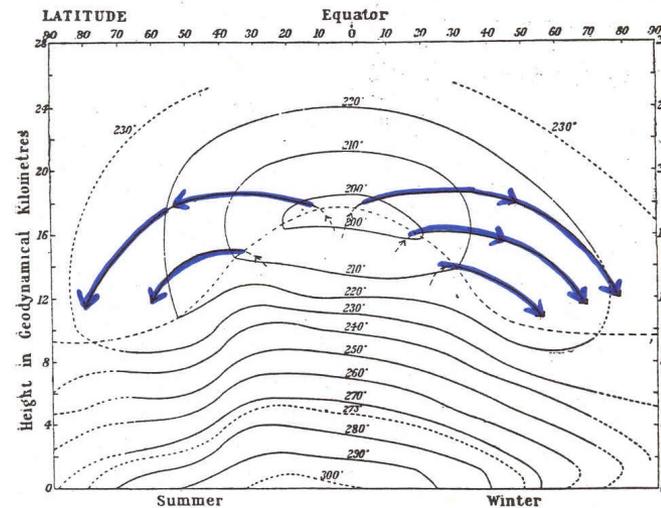


FIG. 5. Isotherms over the Globe  
A supply of dry air is maintained by a slow mean circulation from the equatorial tropopause.

see recent review by Butchart 2014

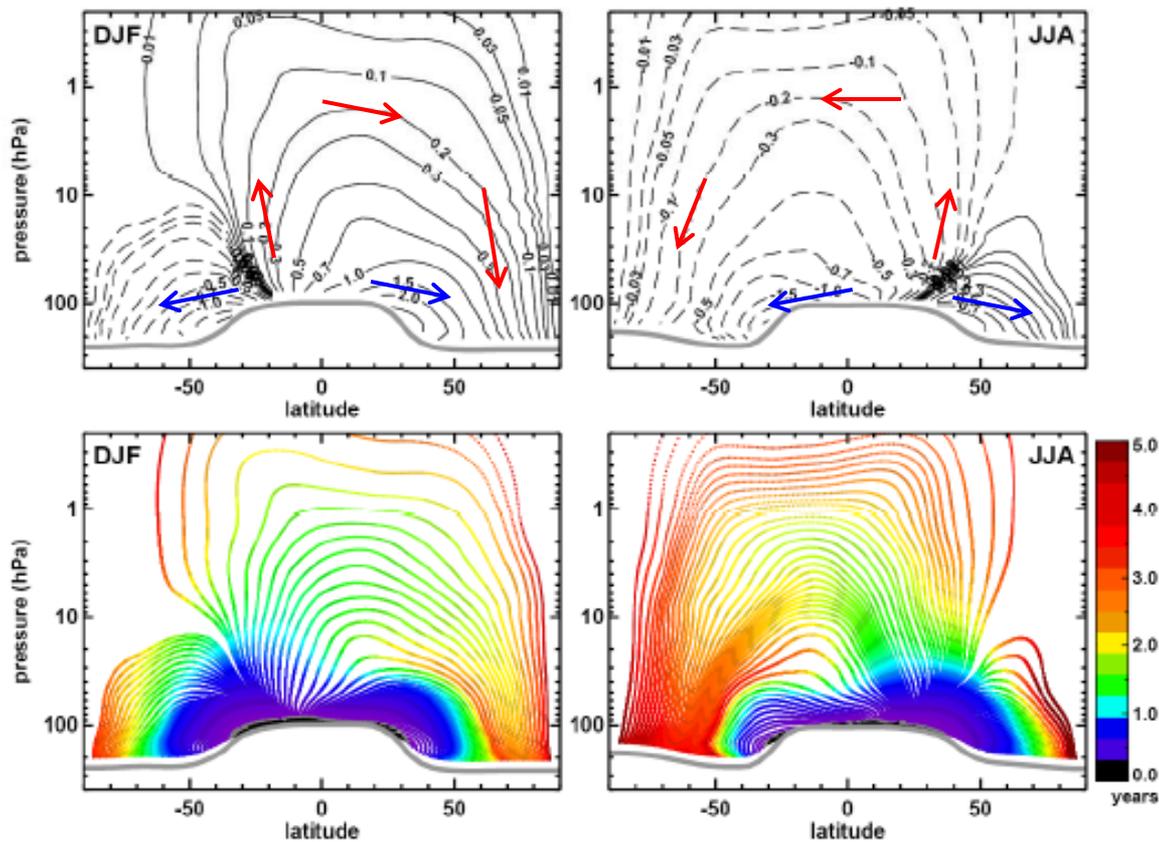
# Residual circulation trajectories and transit times into the extratropical lowermost stratosphere

T. Birner<sup>1</sup> and H. Bönisch<sup>2</sup>

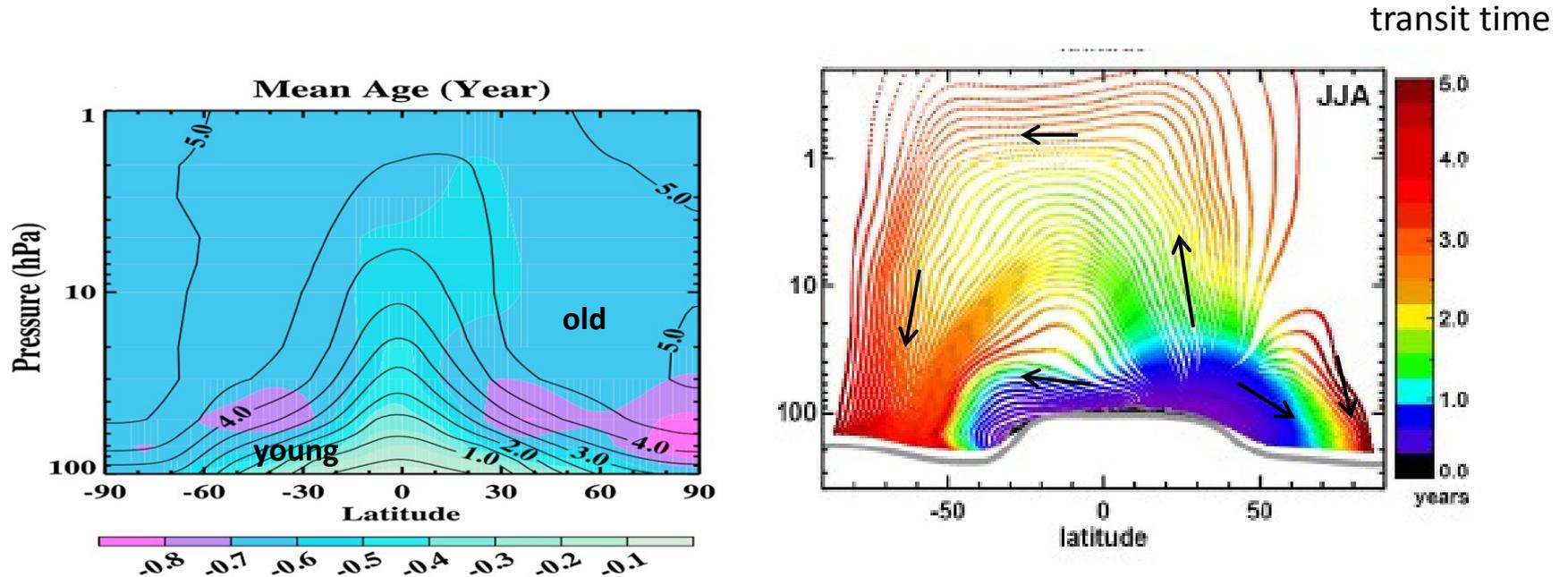
ACP 2011

renewed appreciation that there are upper and lower branches of the BDC

Red: deep branch (slow)  
Blue: shallow branch (fast)



Transit time is closely related to 'mean age' (time since air entered stratosphere)

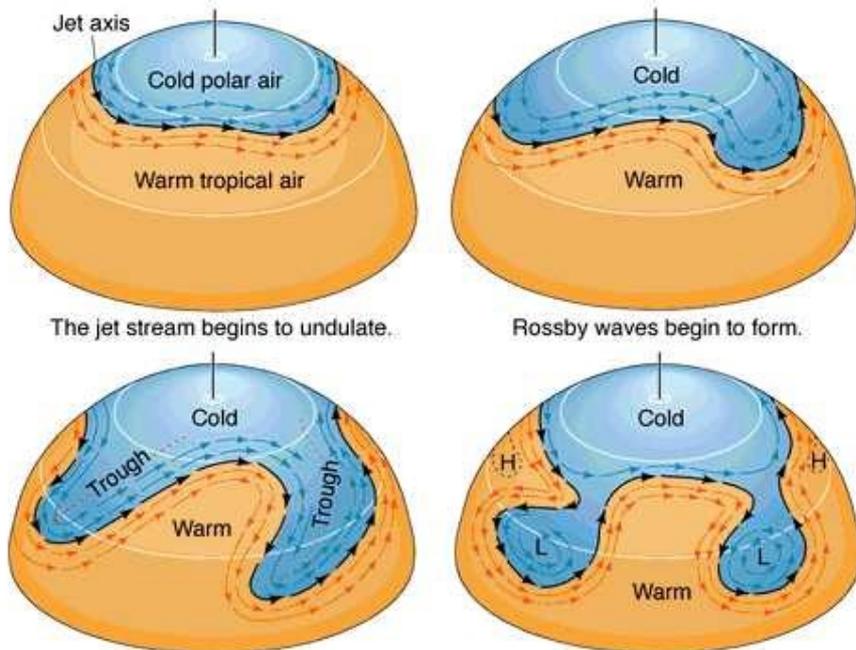


Air at any particular location is characterized by a distribution of transit times and ages (so-called age spectrum)

## Key points:

- Asymmetry in winter stratosphere circulations: more disturbed in the NH, cold and quiet in the SH
- The stratosphere is forced by waves from the troposphere (stronger forcing in NH; episodic stratospheric sudden warmings)
- Dynamical response of balanced vortex to wave forcing (non-local temperature and wind changes)
- Eliassen-Palm (EP) fluxes quantify wave forcing
- Brewer-Dobson transport circulation (deep and shallow branches)

# Rossby waves



The jet stream begins to undulate.

Rossby waves begin to form.

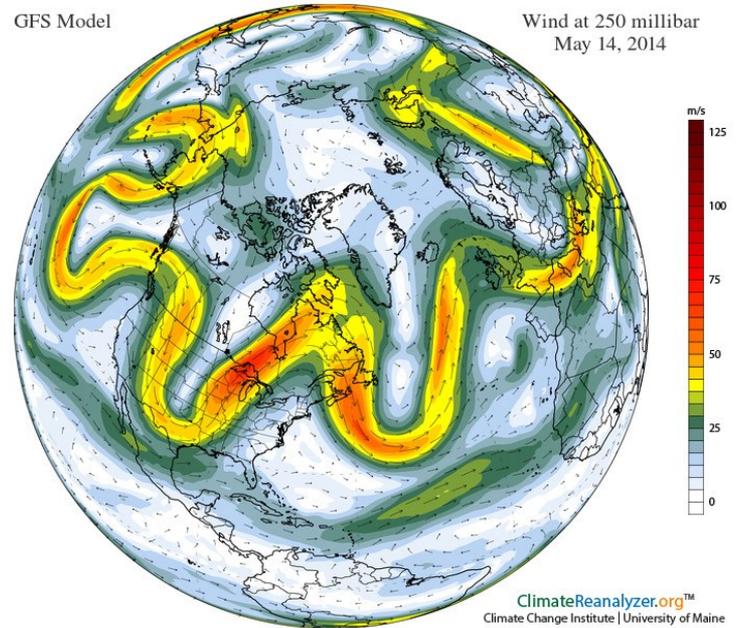
Waves are strongly developed. The cold air occupies troughs of low pressure.

When the waves are pinched off, they form cyclones of cold air.

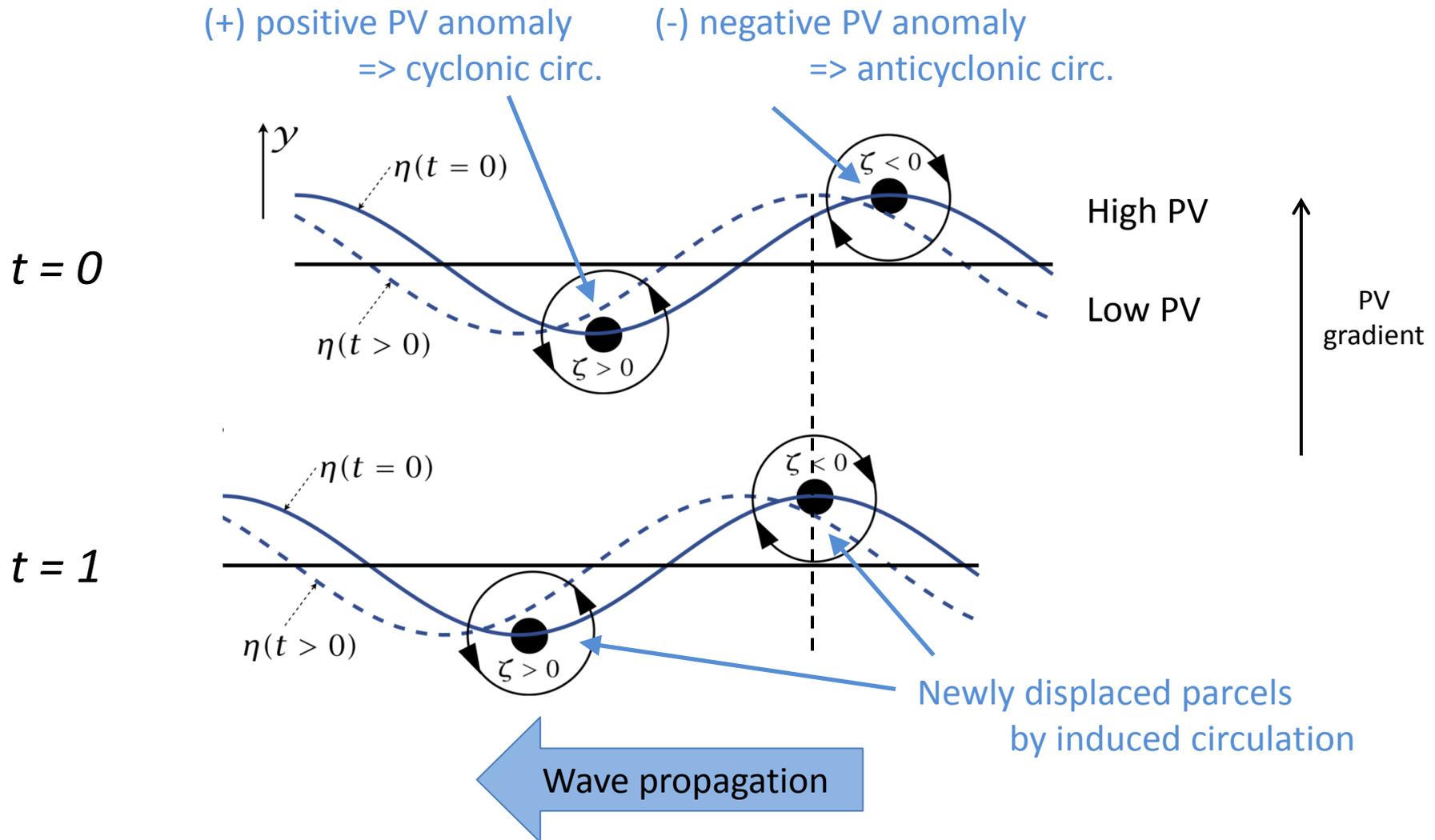
Copyright © A.N. Strahler.

GFS Model

Wind at 250 millibar  
May 14, 2014



# Rossby wave (potential vorticity wave)



## Rossby wave propagation: quasi-geostrophic linearized PV equation

$$\left( \frac{\partial}{\partial t} + \frac{\bar{u}}{a \cos \phi} \frac{\partial}{\partial \lambda} \right) q'_{(M)} + a^{-1} \bar{q}_\phi v' = 0,$$

↑  
eddy PV

↑  
background PV gradient

wave solution:

$$\Phi' = e^{z/2H} \operatorname{Re} \Psi(\phi, z) e^{is\lambda}$$

$$\bar{q}_\phi = 2\Omega \cos \phi - \left[ \frac{(\bar{u} \cos \phi)_\phi}{a \cos \phi} \right]_\phi - \frac{a}{\rho_0} \left( \frac{\rho_0 f^2}{N^2} \bar{u}_z \right)_z.$$

$$\frac{f^2}{a^2 \cos \phi} \left( \frac{\cos \phi}{f^2} \Psi_\phi \right)_\phi + \frac{f^2}{N^2} \Psi_{zz} + n_s^2 \Psi = 0$$

**wave equation:**  
propagation for  
 $n_s^2 > 0$

$$n_s^2 = \frac{\bar{q}_\phi}{a\bar{u}} - \frac{s^2}{a^2 \cos^2 \phi} - \frac{f^2}{4N^2 H^2}$$

refractive index

# Propagation of Planetary-Scale Disturbances from the Lower into the Upper Atmosphere

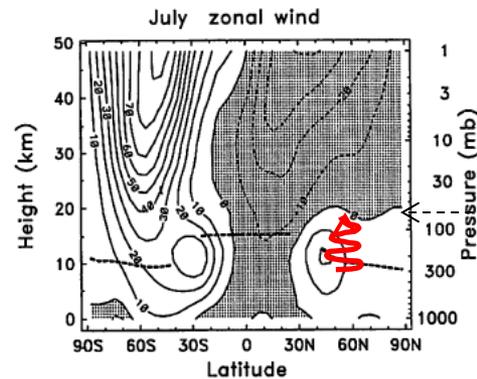
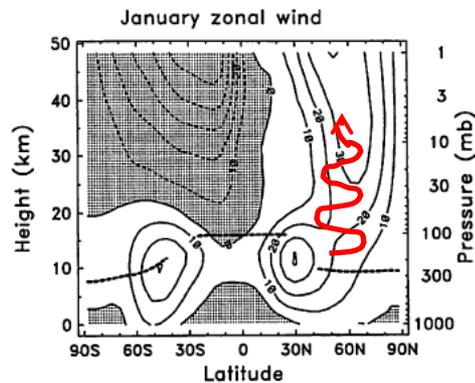
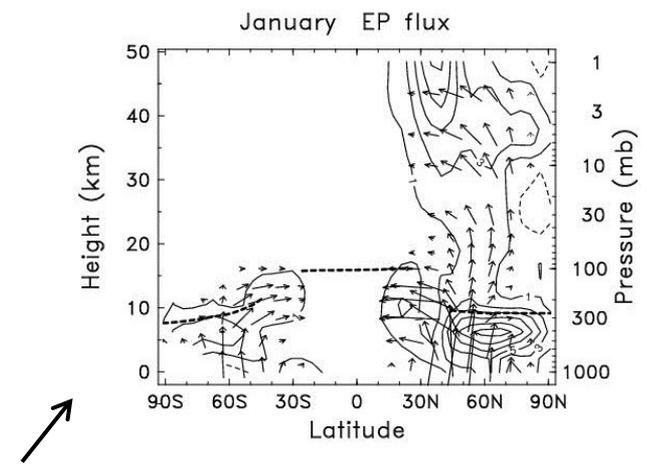
J. G. CHARNEY AND P. G. DRAZIN<sup>1</sup>

JGR 1961

$$n_s^2 = \frac{\bar{q}_\phi}{a\bar{u}} \sim \cos(\text{lat}) + U \text{ terms} - \frac{s^2}{a^2 \cos^2 \phi} - \frac{f^2}{4N^2 H^2}$$

## 2 key points:

- $n_s^2$  proportional to  $\sim \cos(\text{lat})$  (Rossby wave refraction towards low latitudes)
- vertical propagation for  $U > 0$  and small zonal wavenumbers (planetary waves propagate to stratosphere only during winter)

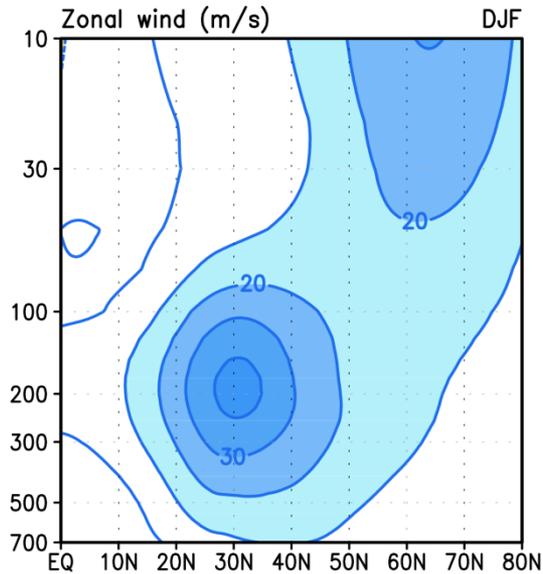


no propagation in summer easterlies

# Refractive index squared (Matsuno 1970)

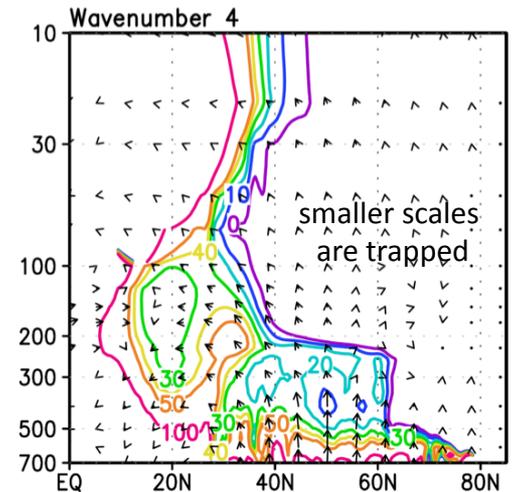
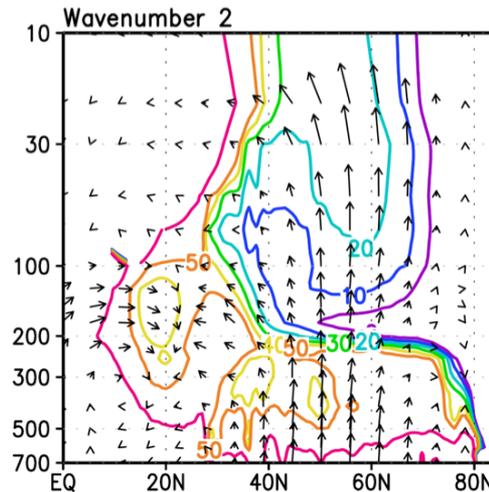
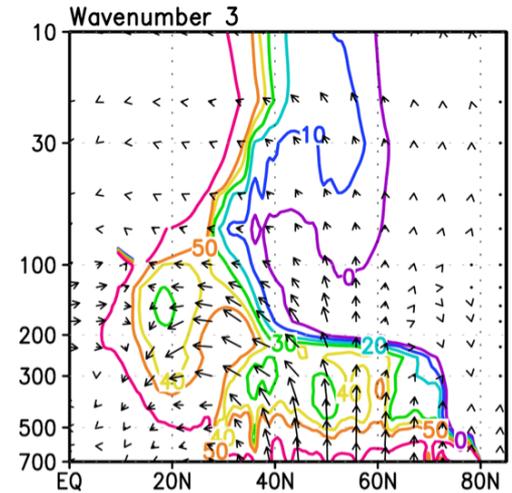
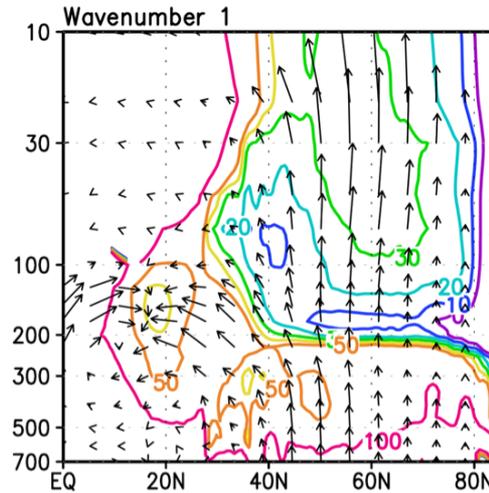
calculated  $n_s^2$  and  
observed EP fluxes

$$n_s^2 = \frac{\bar{q}_\phi}{a\bar{u}} - \frac{f^2}{4H^2N^2} - \frac{s^2}{a^2 \cos^2 \phi}$$

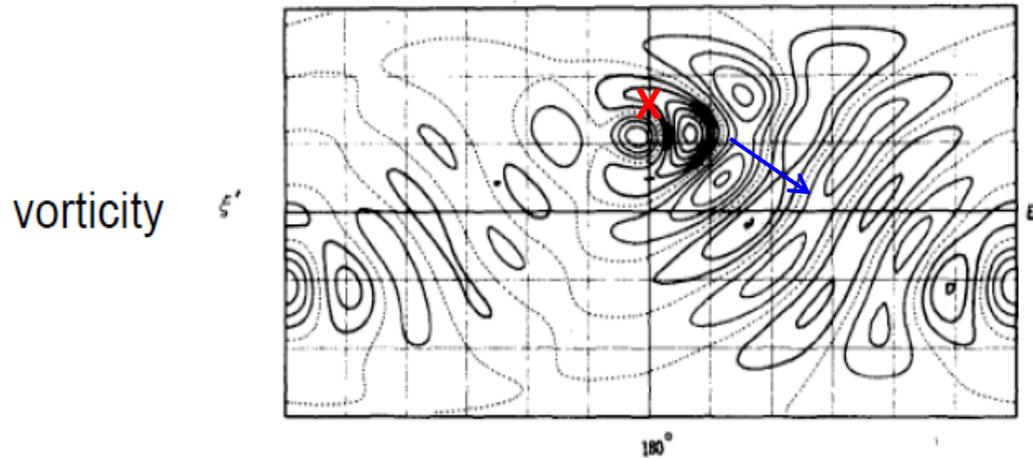


thanks to Joowan Kim

zonal wave 1 and 2 can propagate  
vertically in climatological basic state

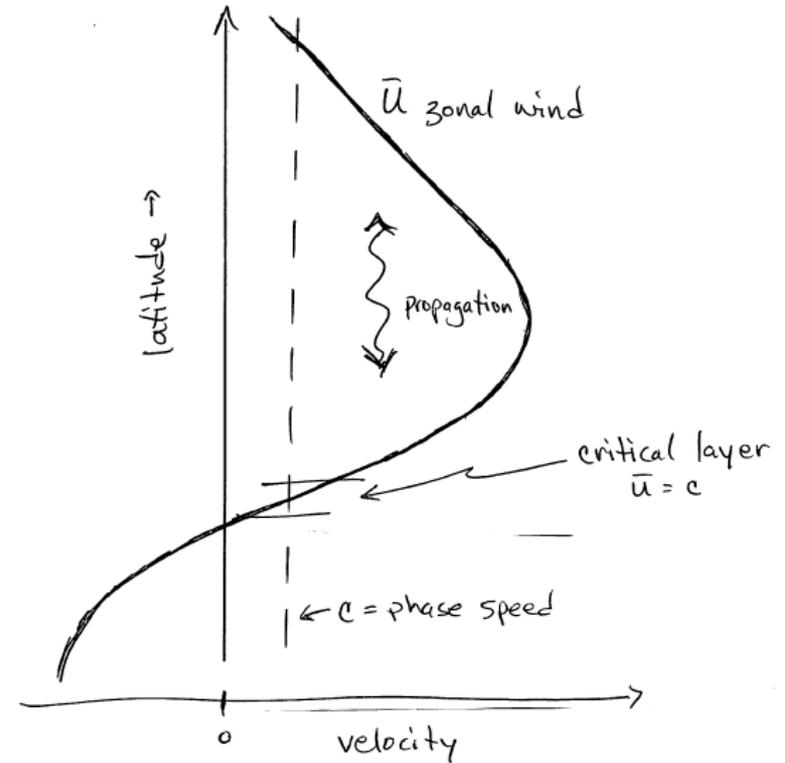
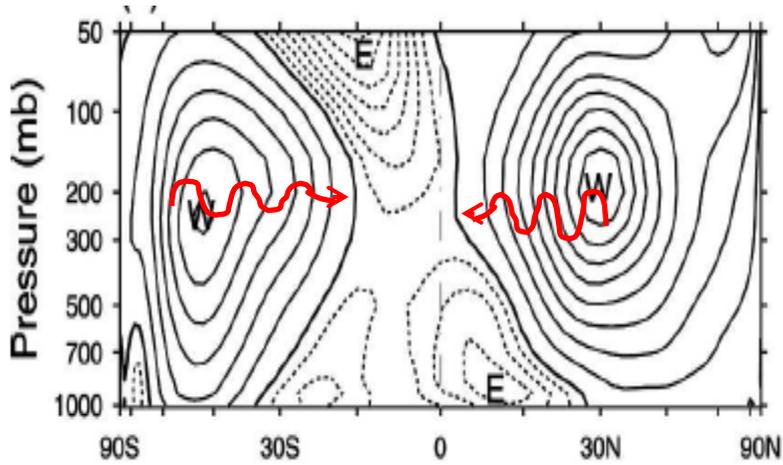


# Rossby wave dispersion on the sphere: response to isolated topography

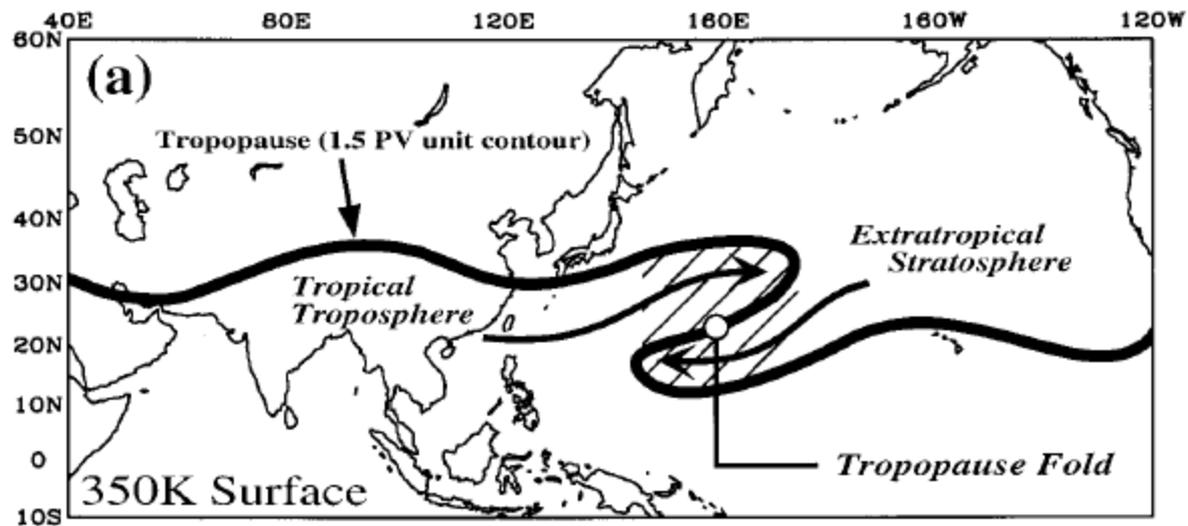


Steady state solution of the shallow water equations linearized about a constant angular velocity flow and perturbed by a circular mountain at  $30^\circ N, 180^\circ E$ . The model includes linear drag. After B. J. Hoskins, *Horizontal Wave Propagation on the Sphere*, in *The General Circulation. Theory, Modeling and Observations*. NCAR Summer Colloquium 1978.

Rossby waves cannot propagate into regions where  $U < 0$  (i.e. across equator)



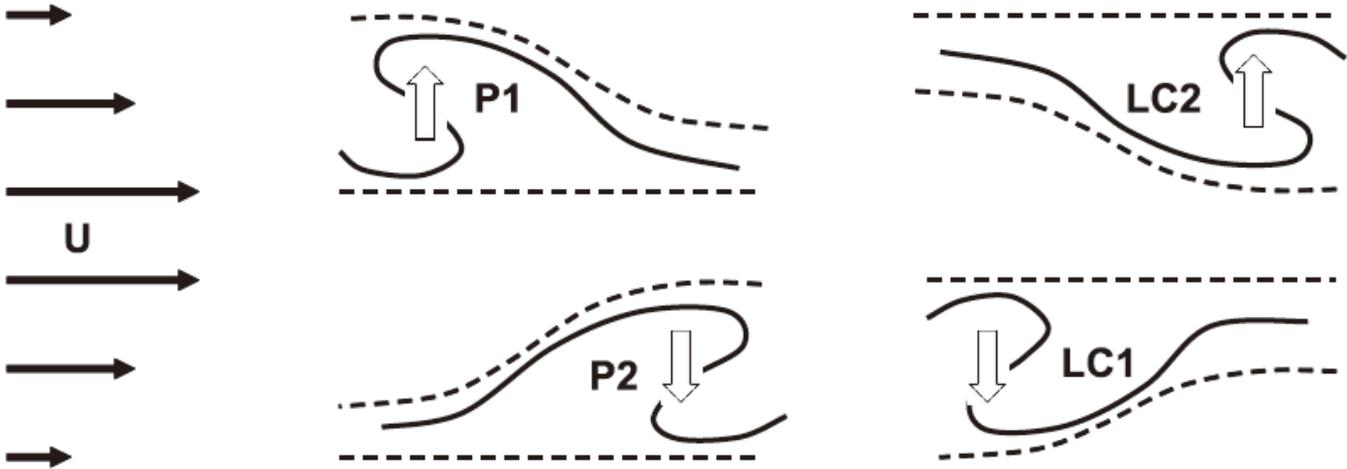
# Breaking Rossby waves: overturning of PV contours



Postel and Hitchman 1999  
Homeyer et al 2013

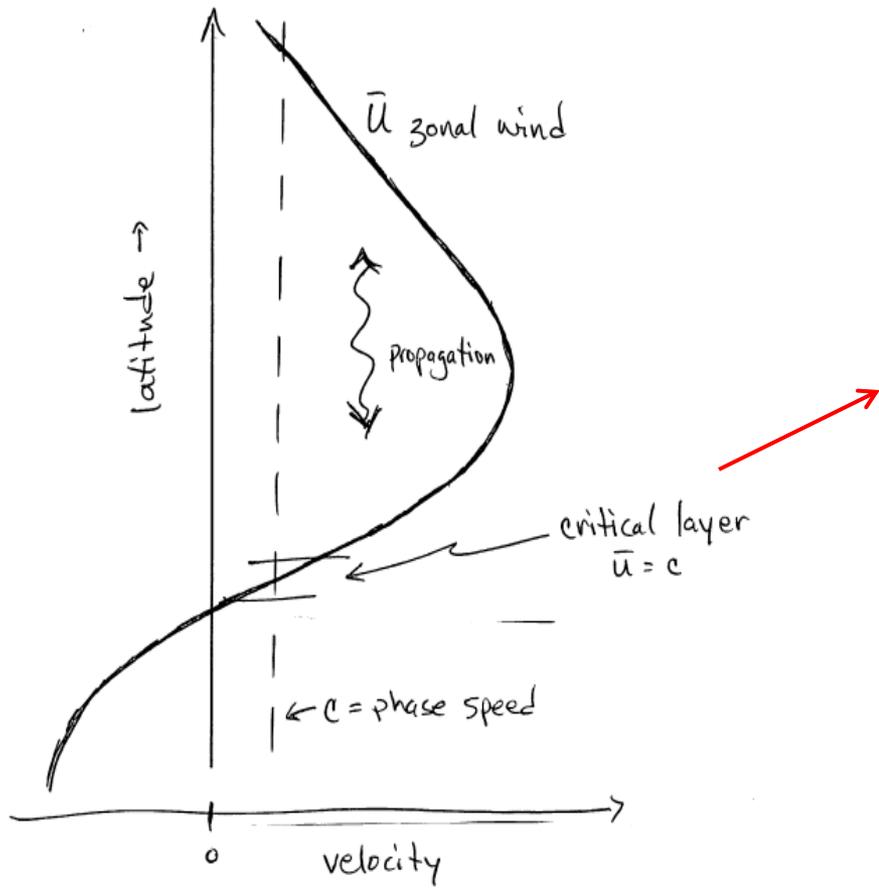
Two types of wave breaking, depending on shear of background winds

poleward breaking (cyclonic)

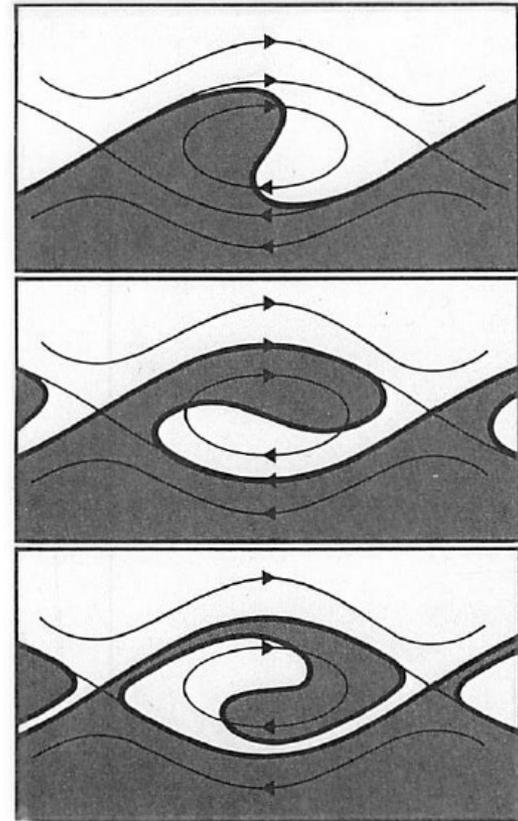


equatorward breaking (anticyclonic)

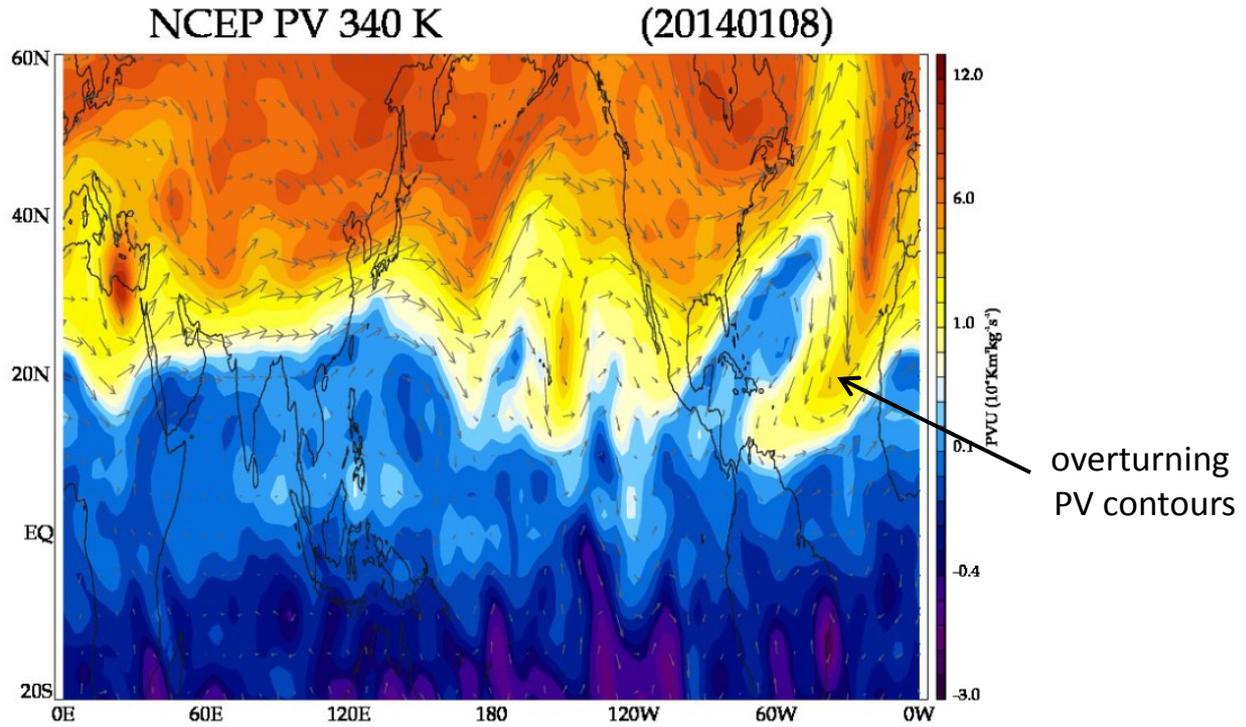
# Rossby wave critical layer interactions (critical layer: $U = c$ )



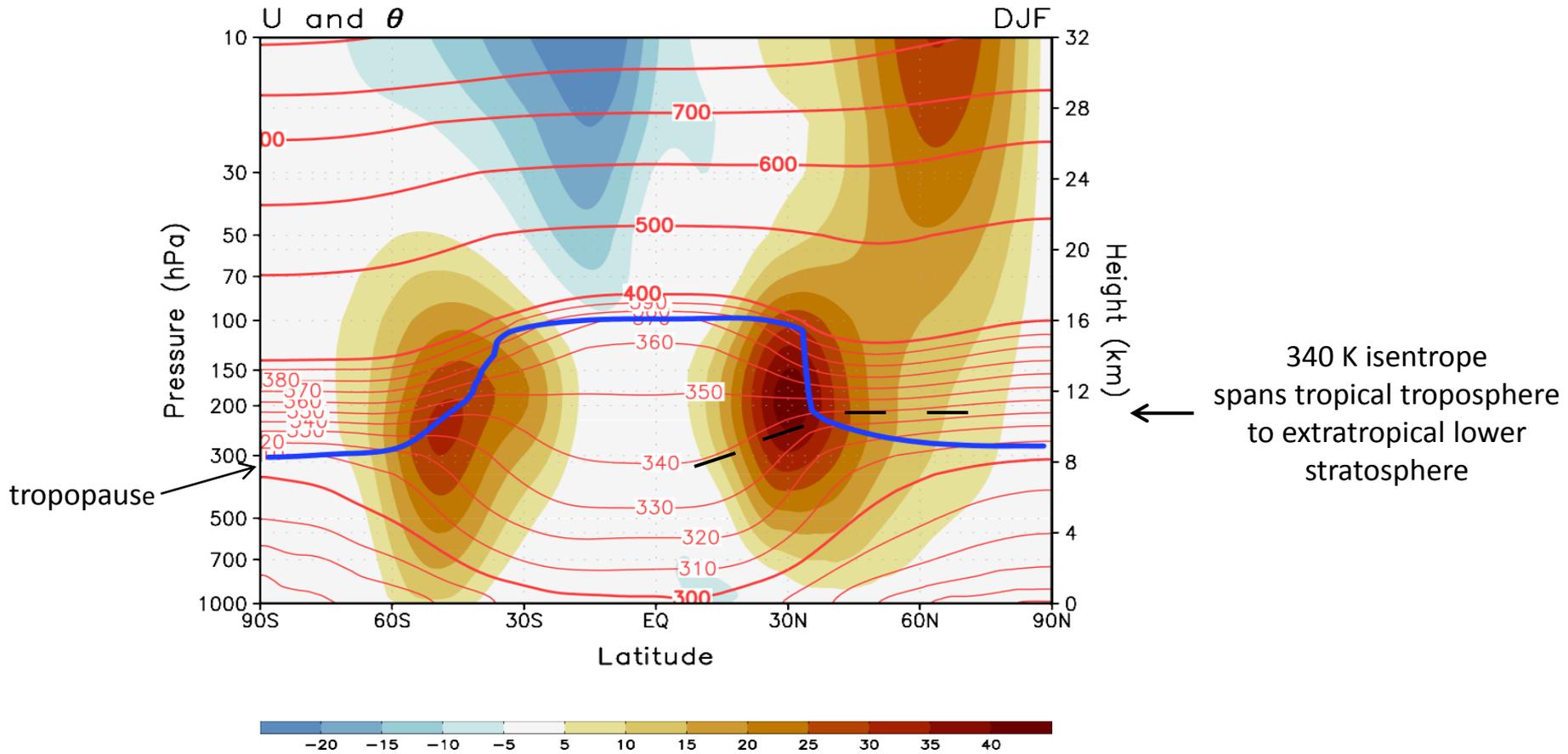
nonlinear overturning at critical layer  
(irreversible transport and mixing)



# Example of a large-scale breaking Rossby wave



fast, synoptic flow mainly along isentropes:



# Rossby waves during January-May at 340 K

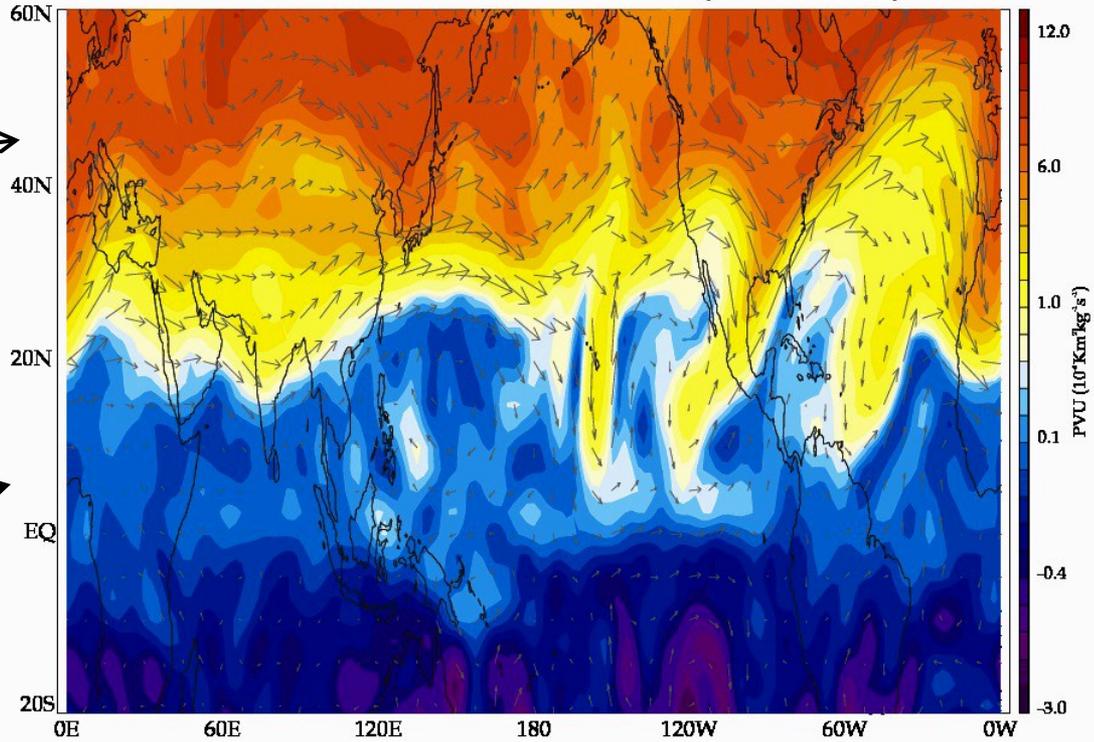
NCEP PV 340 K

(20140130)

extra-tropical  
lower stratosphere



tropical  
troposphere



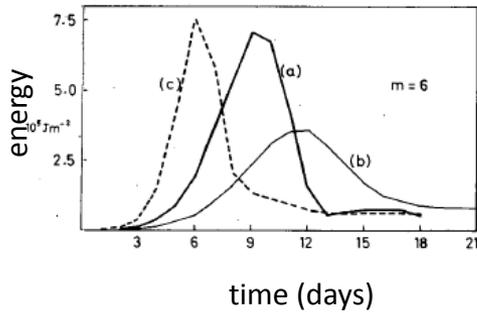
Key points:

- Rossby waves: conserve PV
- General refraction of Rossby waves towards low latitudes
- Latitudinal or vertical propagation for  $U > 0$  (more generally  $U > c$ )
- Rossby wave breaking near critical lines ( $U = c$ )
- Poleward or equatorward breaking depending on background  $U$  shear
- Key mechanism for dissipation, transporting PV or trace species

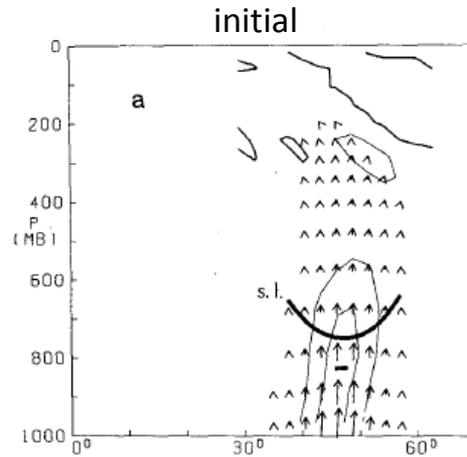


# Extratropical EP flux patterns are related to baroclinic wave life cycles

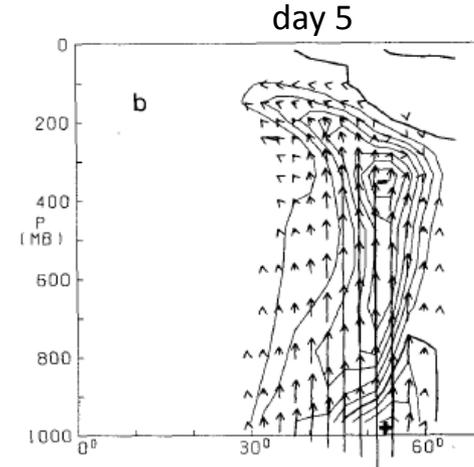
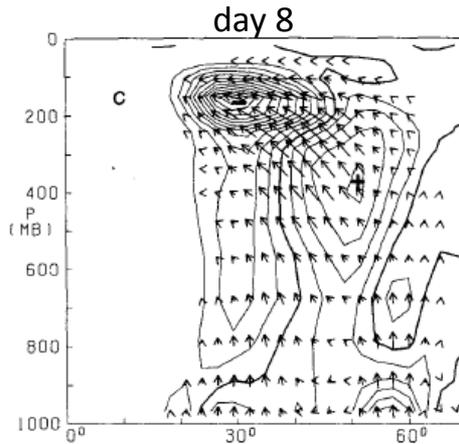
idealized  
zonal wave 6  
baroclinic eddy  
life cycle



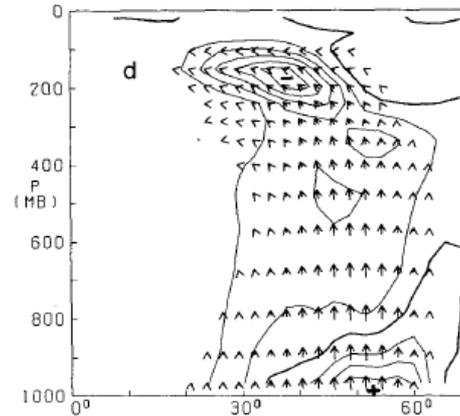
barotropic  
decay



TOTAL E-P FLUX DIVERGENCE  
DAY -0.00



TOTAL E-P FLUX DIVERGENCE  
DAY 5.00  
life cycle average



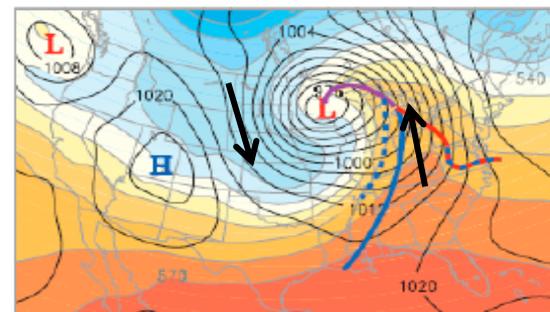
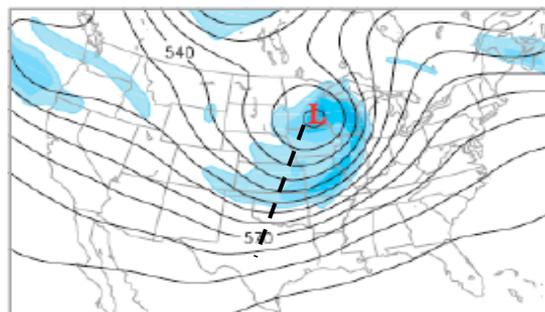
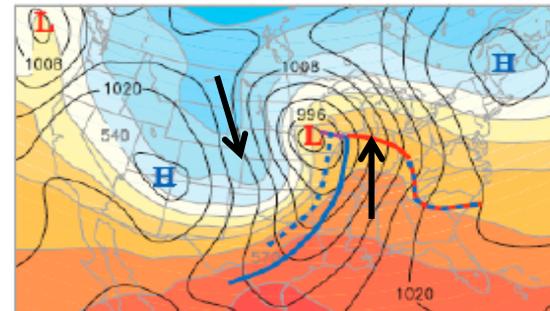
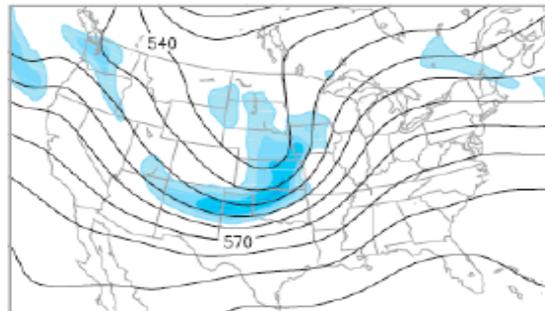
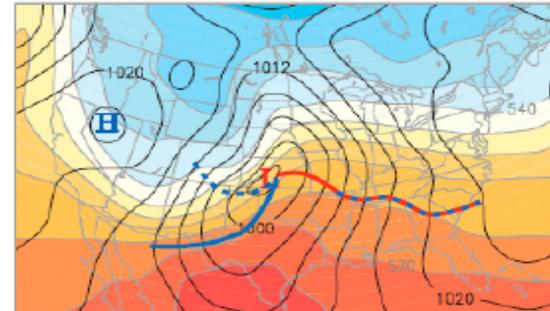
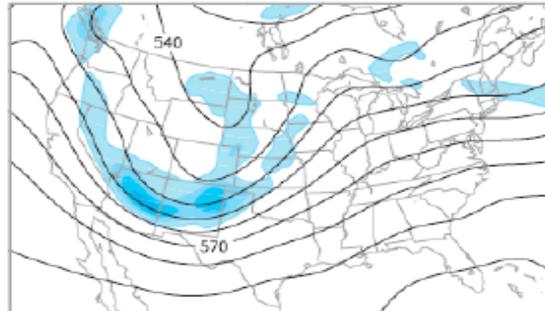
baroclinic  
growth

Simmons and Hoskins 1980  
Edmon et al 1980

synoptic views  
of developing  
baroclinic  
wave

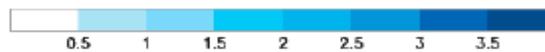
500 hPa

surface



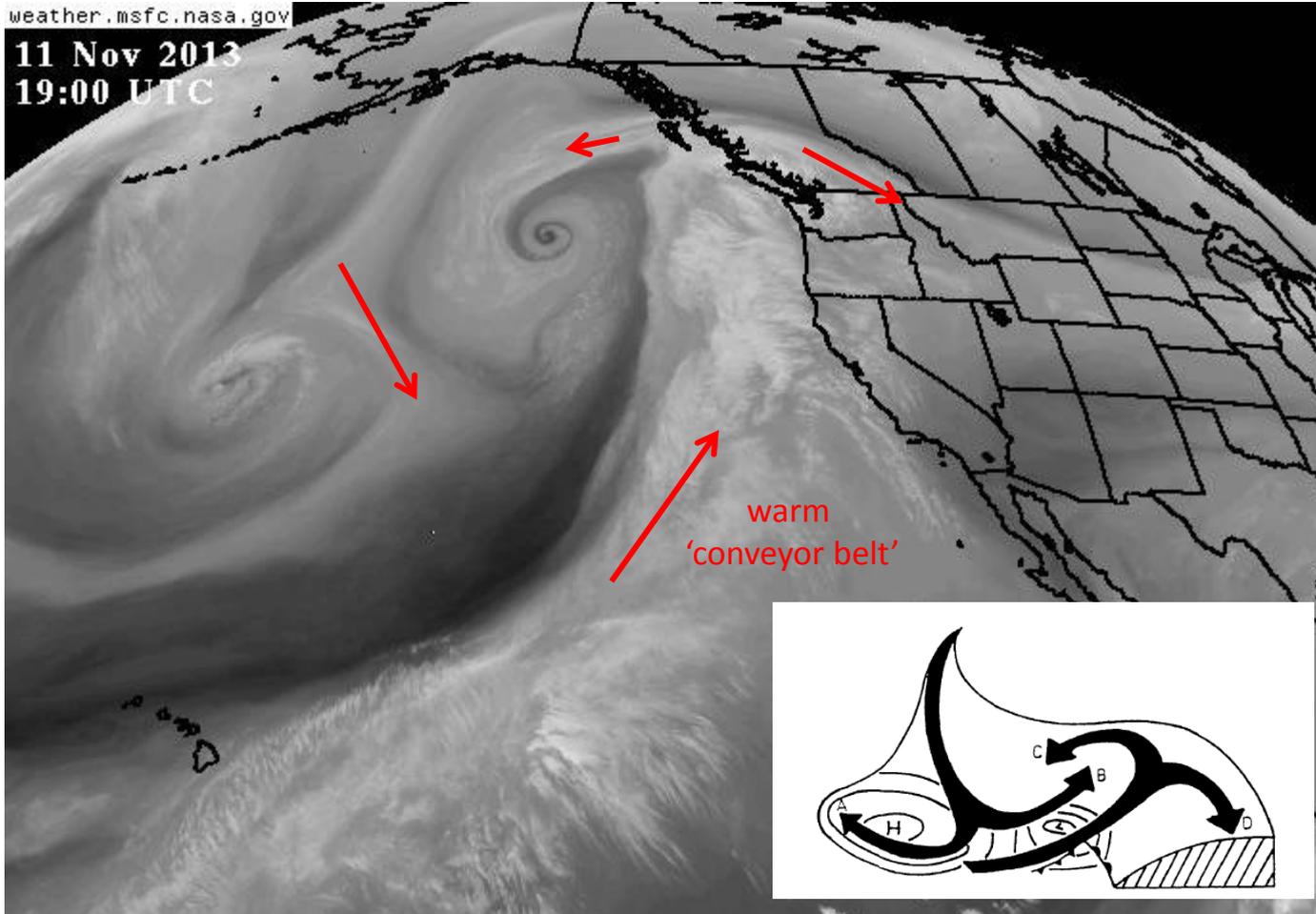
poleward  
heat flux

phase tilt:  
equatorward  
momentum flux



weather.msfc.nasa.gov

11 Nov 2013  
19:00 UTC



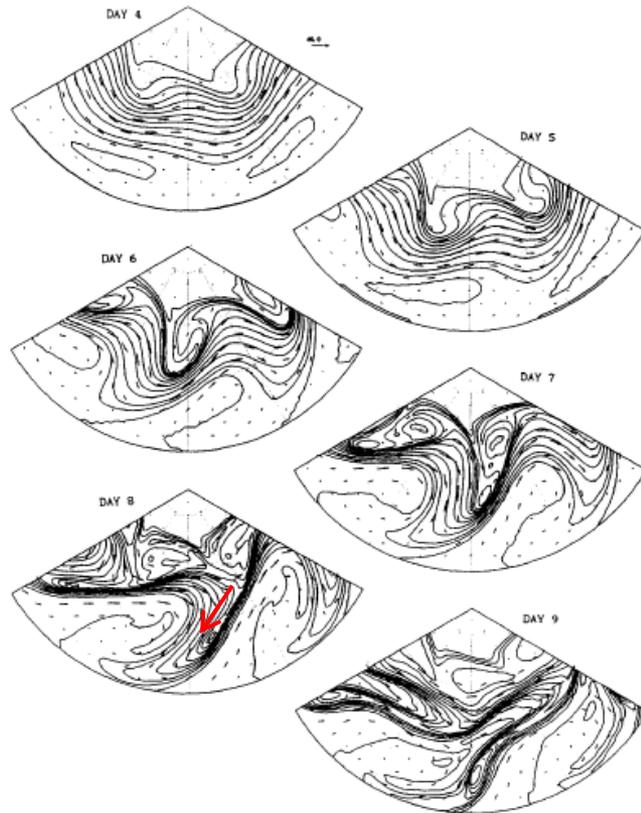
## Two paradigms of baroclinic-wave life-cycle behaviour

By C. D. THORNCROFT<sup>1\*</sup>, B. J. HOSKINS<sup>1</sup> and M. E. McINTYRE<sup>2</sup>

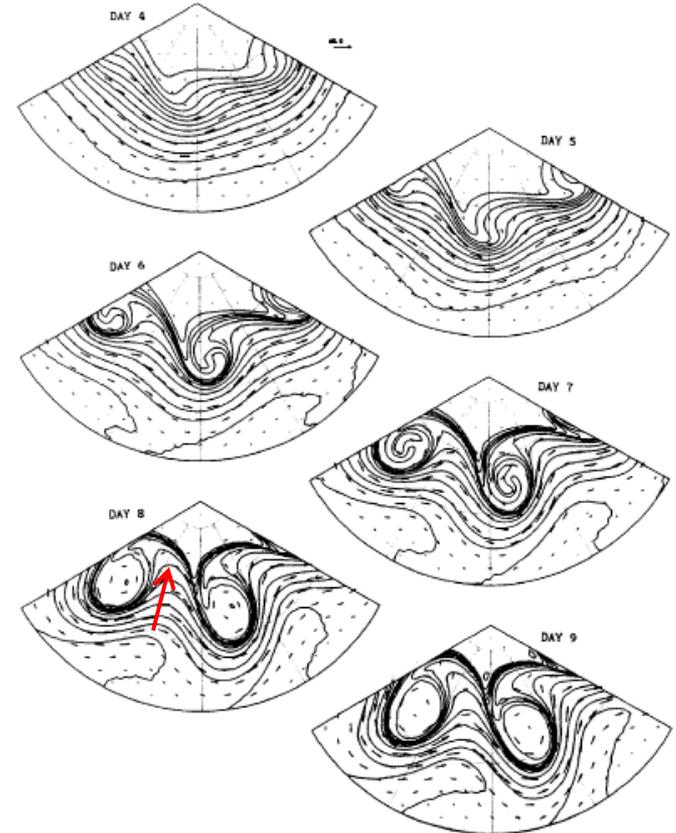
<sup>1</sup>Department of Meteorology, University of Reading

<sup>2</sup>Department of Applied Mathematics and Theoretical Physics, University of Cambridge

### LC1 equatorward breaking

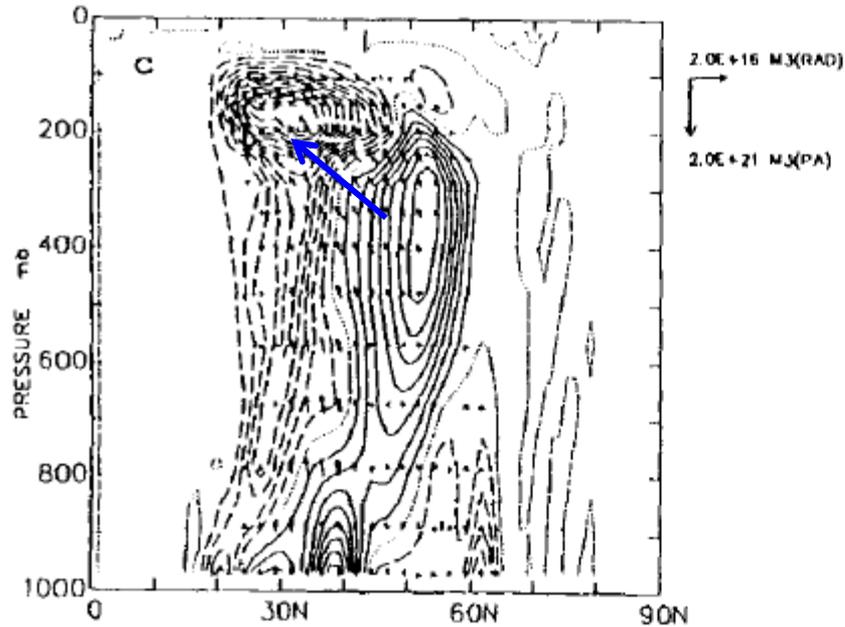


### LC2 poleward breaking

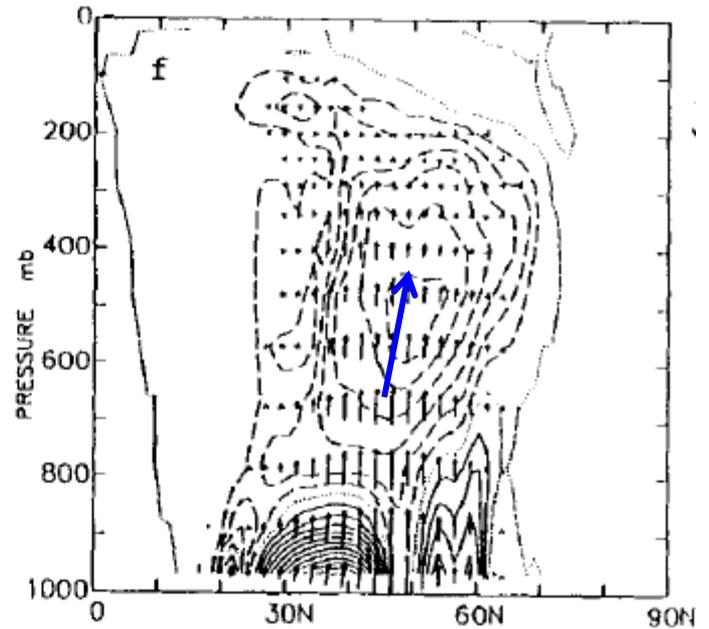


# Idealized baroclinic wave life cycles

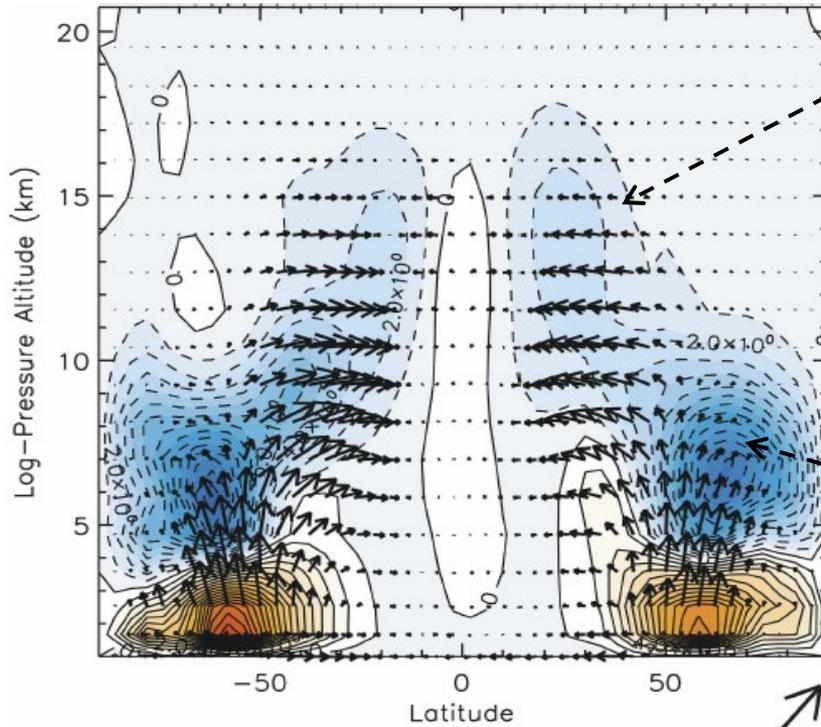
equatorward propagation (LC1)



poleward propagation (LC2)

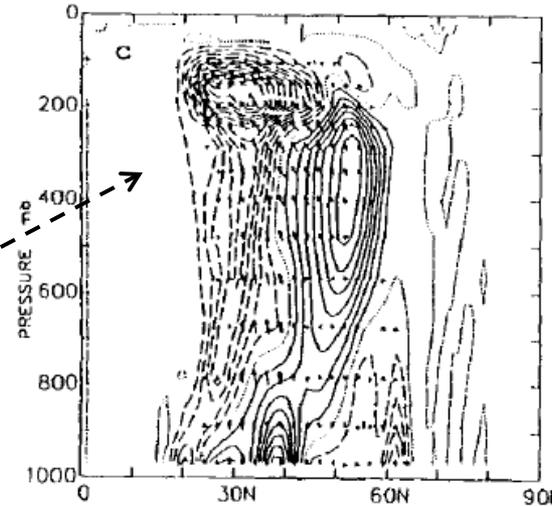


# Climatology

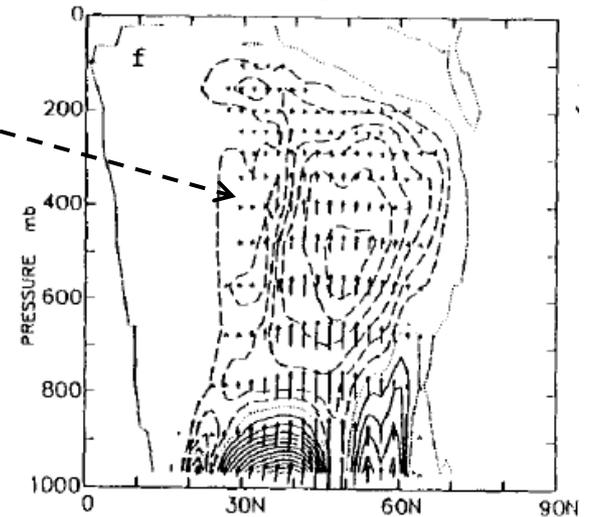


Garcia and Randel, 2008, J. Atmos. Sci.

LC1



LC2



Thorncroft et al., 1993, Q.J.R. Meteorol. Soc.

# Using phase speed spectra to diagnose critical layer interactions

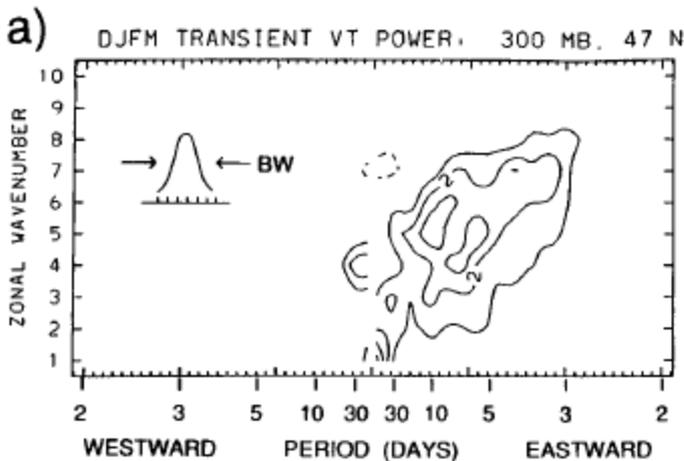
wave flux  
co-spectra as a function  
of zonal wavenumber  
and phase speed

$$K_{n,c} = K_{k,\omega} \cdot \left( \frac{n}{a \cos \phi} \right).$$

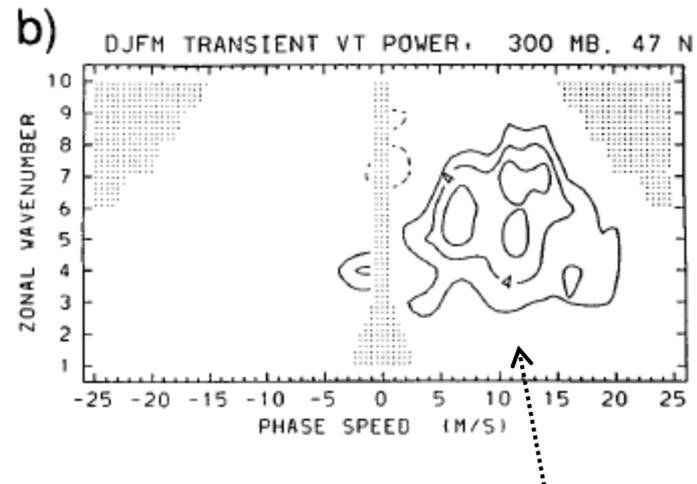
traditional  
wavenumber vs.  
frequency

Randel and Held 1991

wavenumber vs. frequency



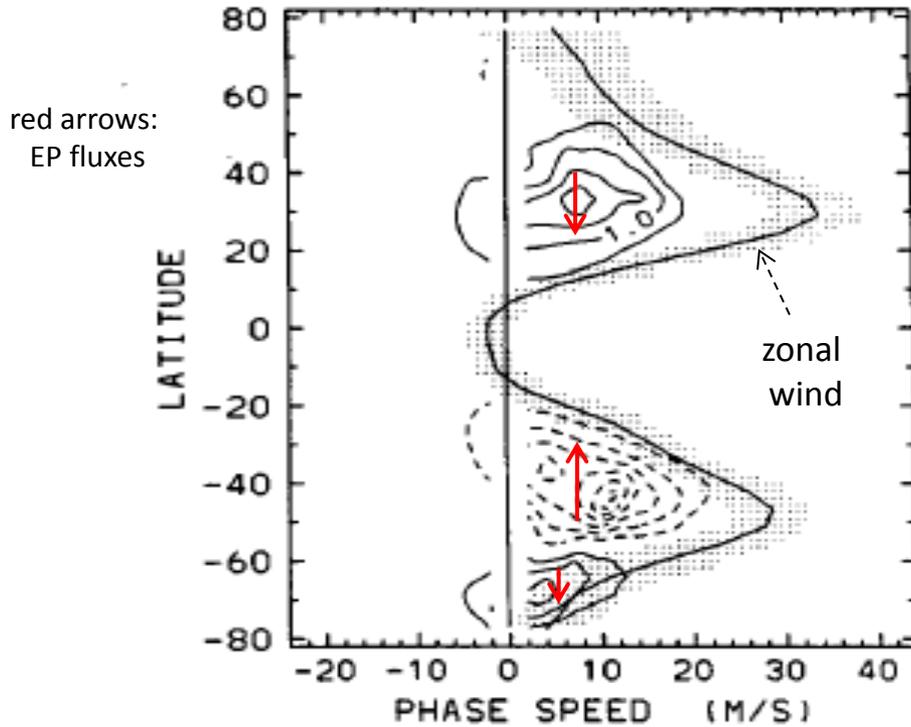
wavenumber vs. phase speed



Rossby waves move eastward  
at ~5-15 m/s

Integrate over wavenumber to derive eddy flux phase speed spectra

eddy momentum flux  $u'v'$  300 hPa

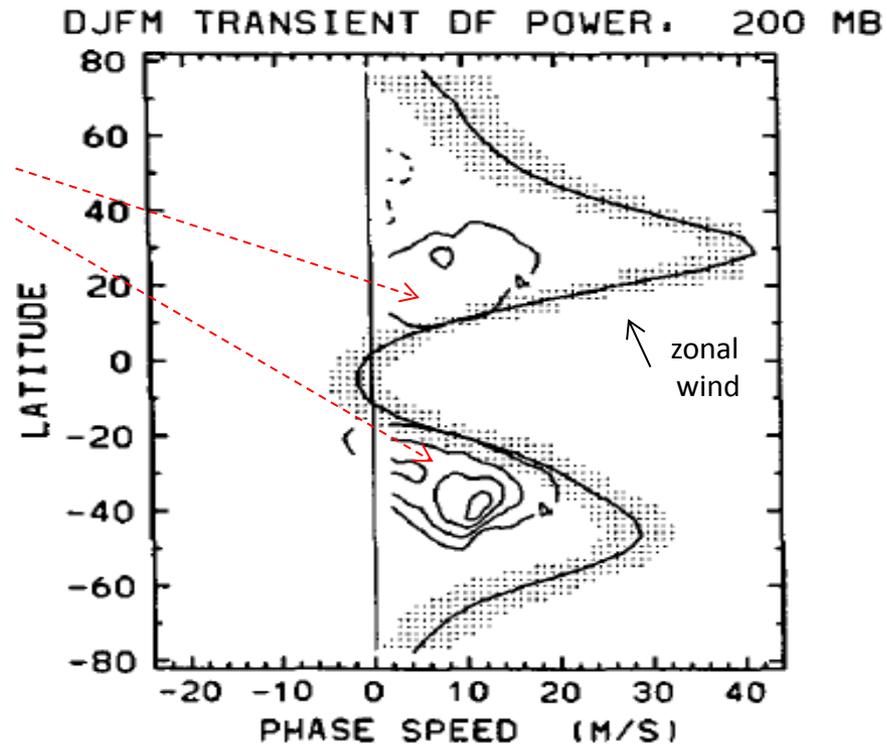


- EP fluxes: propagation to near critical lines ( $c = U$ )
- evidence for critical layer behavior

Randel and Held 1991

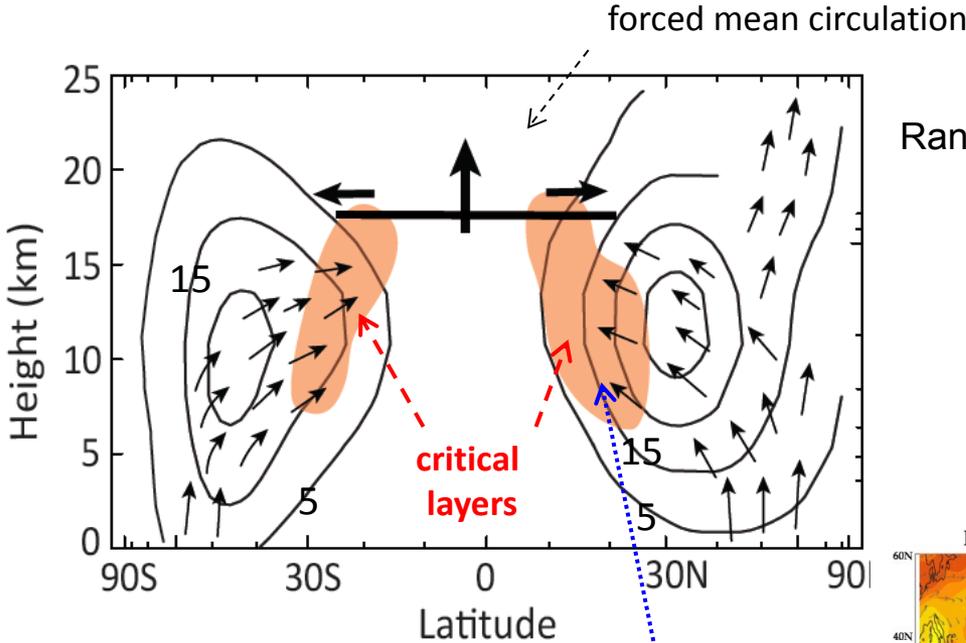
# EP flux divergence phase speed spectra

EP flux divergence  
poleward of critical lines

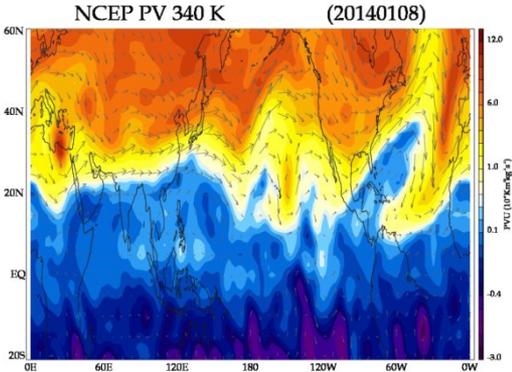


Randel and Held 1991

# Subtropical critical layers for Rossby waves with phase speeds $\sim 5-15$ m/s



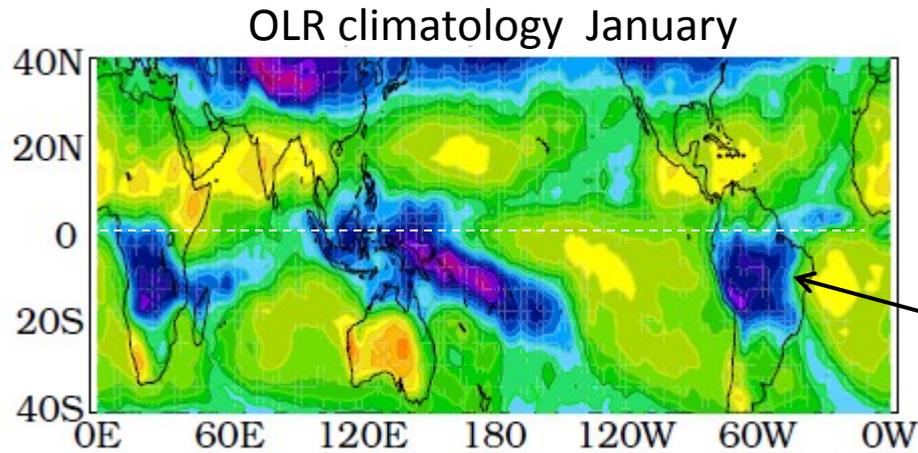
Randel and Jensen, 2013, Nat. Geosci.



## Key points:

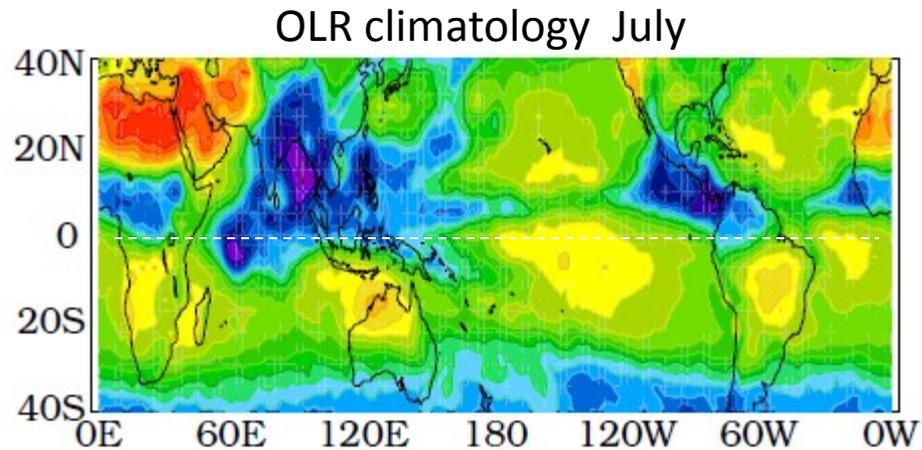
- Baroclinic wave life cycles: baroclinic growth and barotropic decay
- Two idealized types of life cycles: equatorward and poleward wave breaking (LC1 and LC2)
- Consistent with tropospheric EP flux circulation statistics
- Phase speed spectra: clear evidence for critical layers in subtropics (important influence of extratropical waves on tropical circulations)

Large-scale tropical circulations are forced by latent heating from deep convection



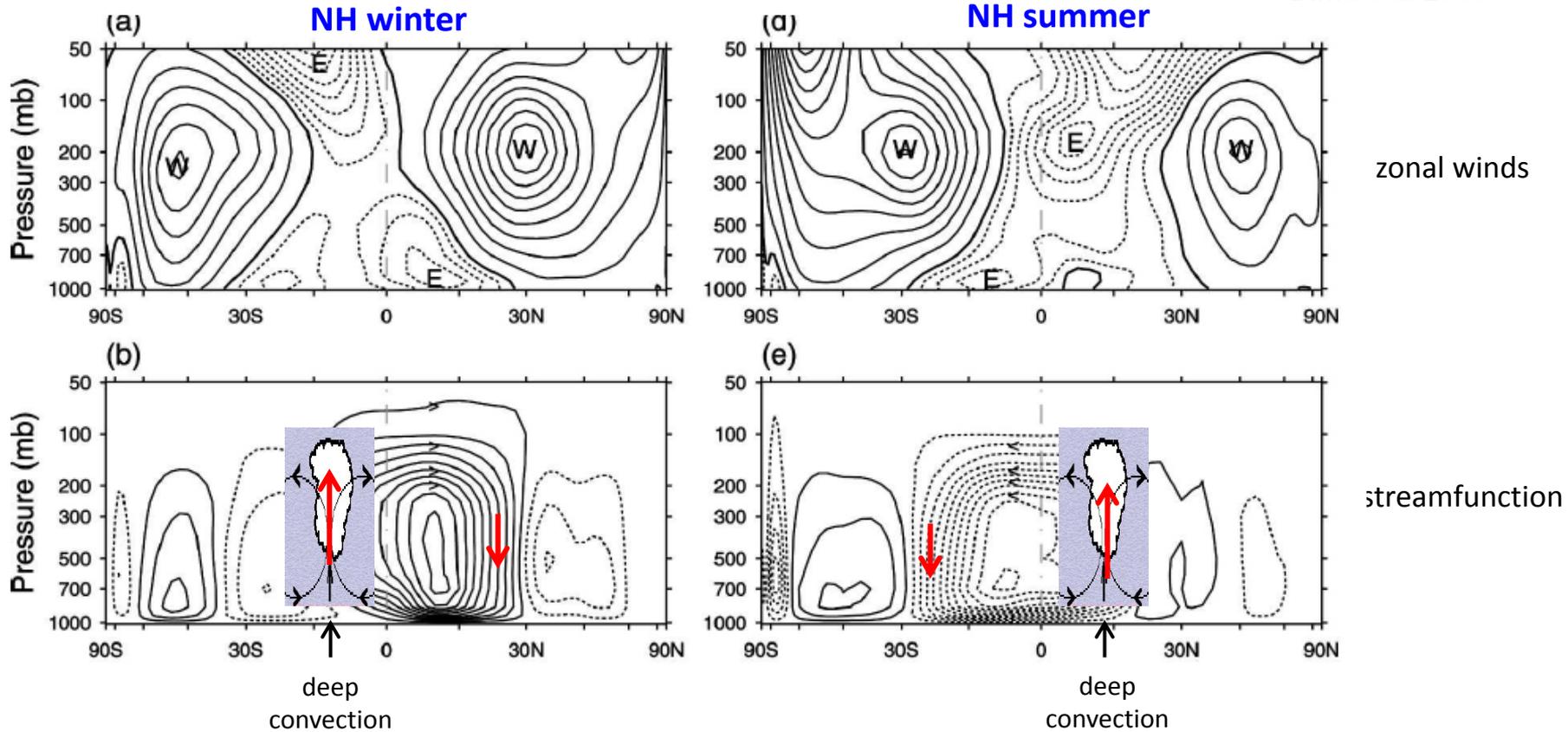
Outgoing Longwave Radiation (OLR) is a useful proxy for deep convection

high clouds ~  
deep convection

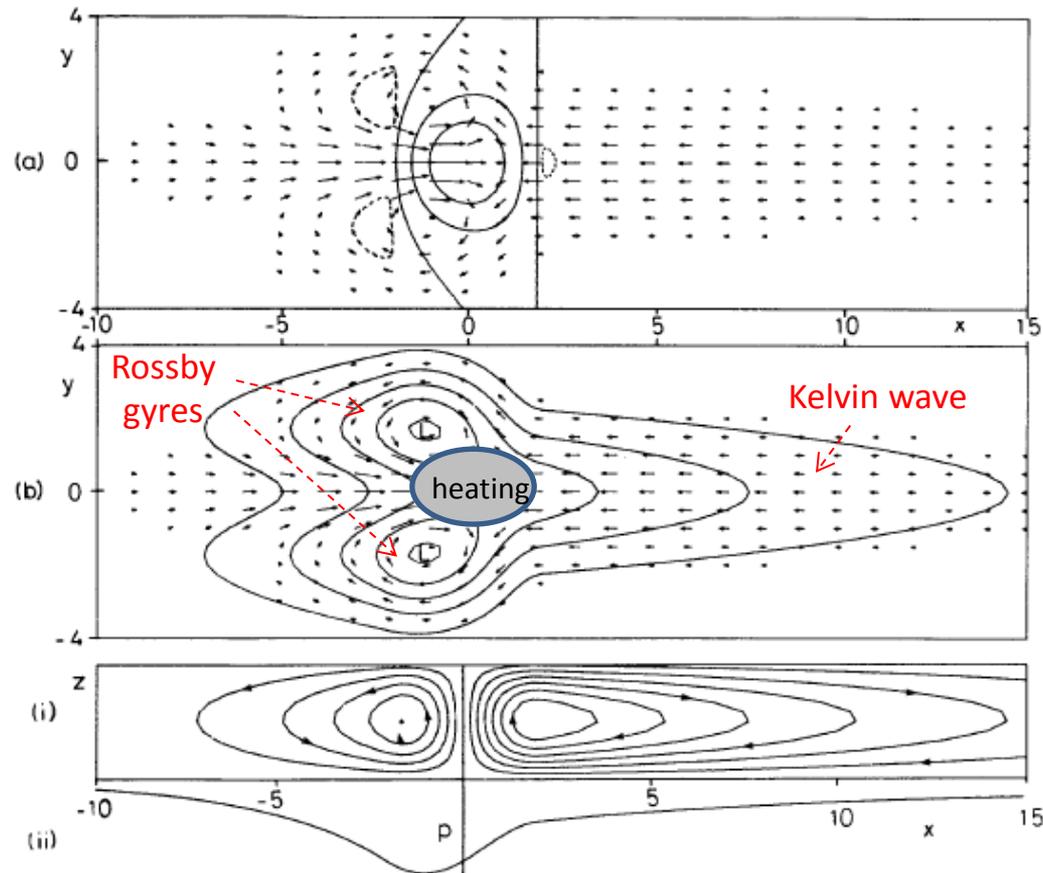


# Seasonal variations in tropical overturning circulation (Hadley cell)

Dima et al 2005



# Dynamical response to low frequency convective forcing



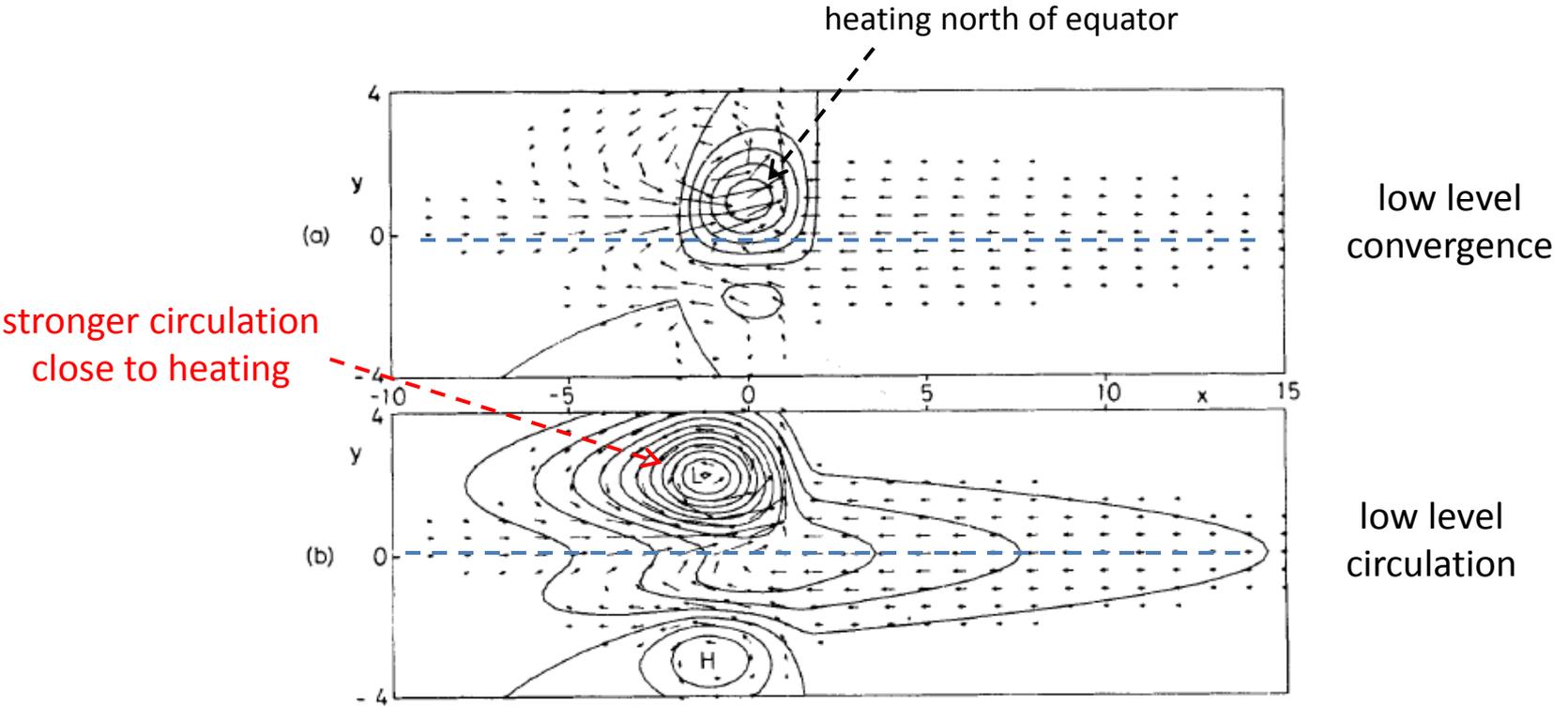
low level  
convergence

low level  
circulation

vertical structure

Gill, 1980

# Dynamical response to heating centered north of equator

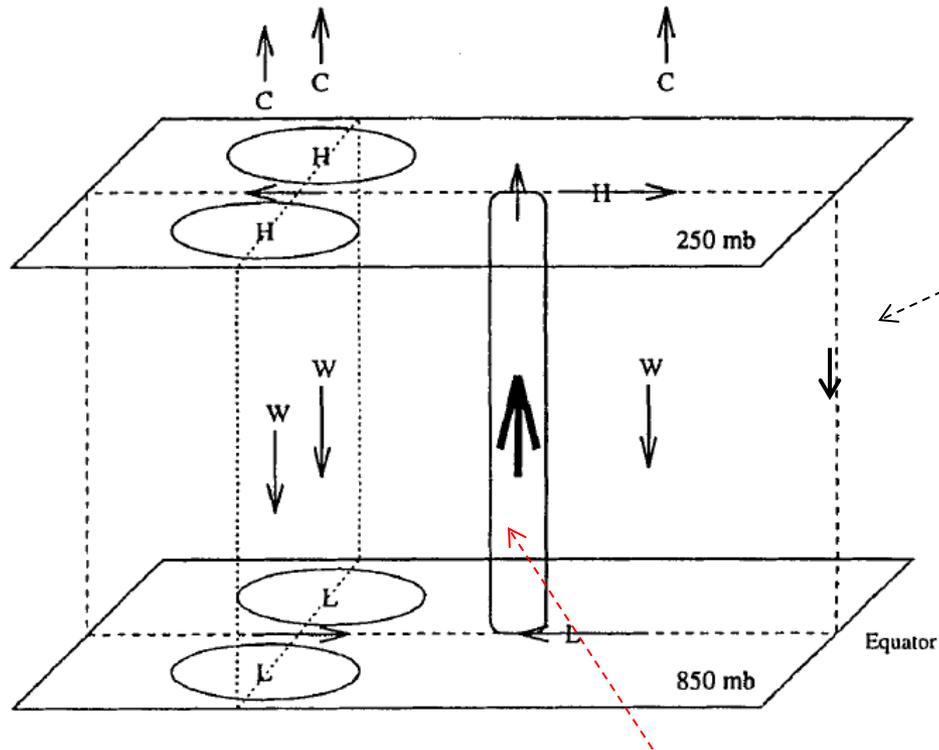


# The tropical tropopause

By E. J. HIGHWOOD\* and B. J. HOSKINS  
*University of Reading, UK*

vertical structure: out of phase between  
lower troposphere and upper troposphere

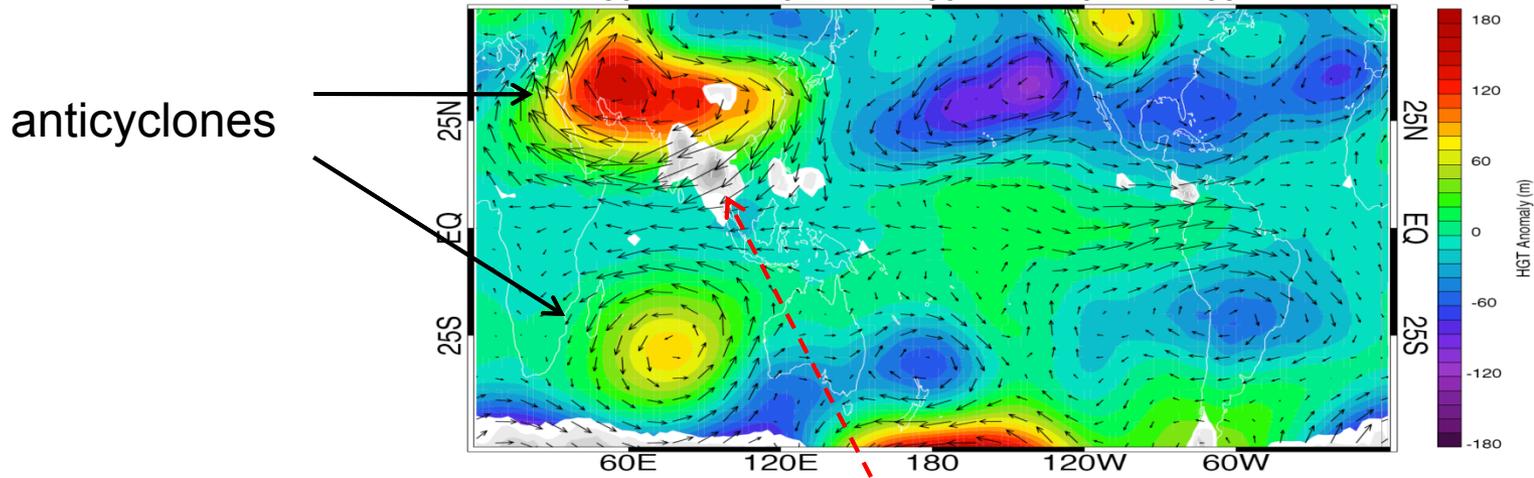
QJRMS 1998



deep convection / heating

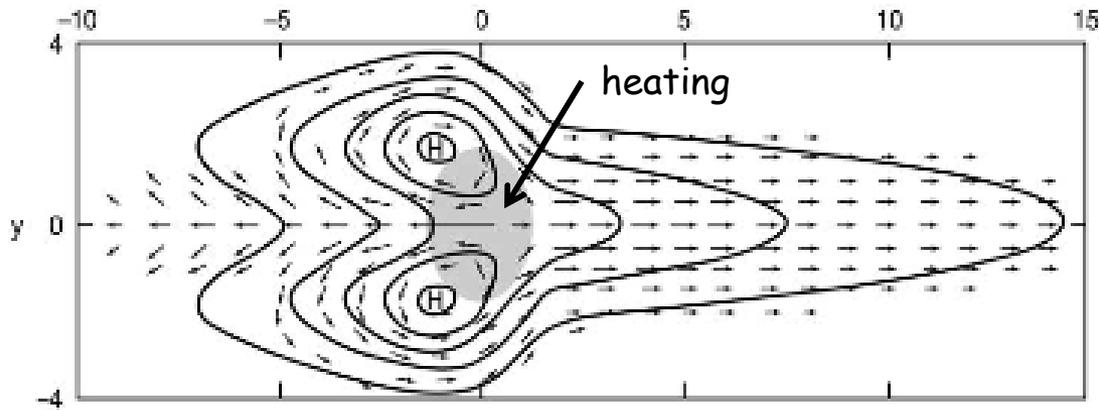
# Tropical heating produces subtropical anticyclones in the UTLS

JJA geopotential height and winds 100 hPa



Convection (heating)

Park et al 2007



Matsuno-Gill Solution

## Key points:

- Organized deep convection (latent heating) drives large-scale tropical circulations
- Seasonal movement between solstices (SH – NH subtropics)
- Hadley and Walker overturning circulations
- Matsuno-Gill dynamical response to local heating: subtropical Rossby waves and equatorial Kelvin waves
- Subtropical anticyclones in UTLS (especially Asian monsoon during NH summer)

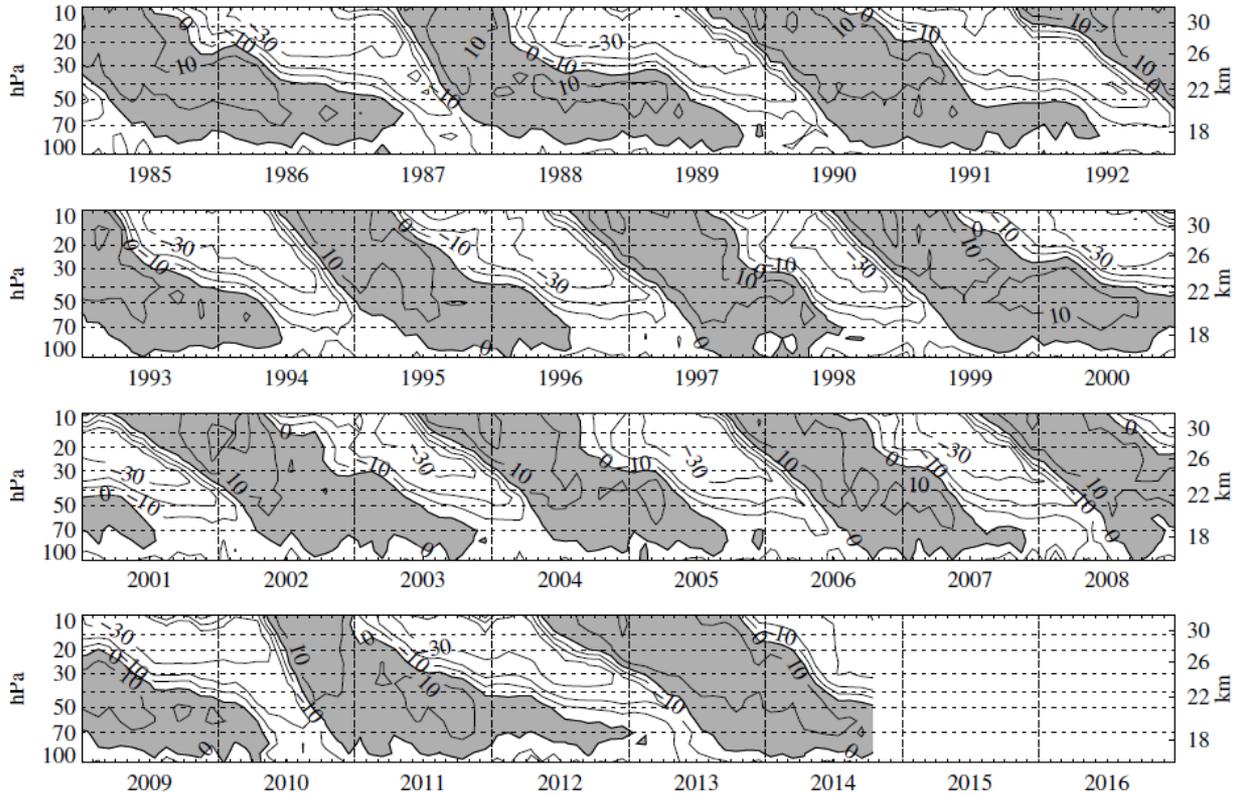
The Quasi-Biennial Oscillation (QBO) and El Niño Southern Oscillation (ENSO) are important modes of circulation in the stratosphere, troposphere and UTLS

- QBO: approximate 28-month oscillation in tropical stratosphere, forced by vertically propagating waves interacting with mean flow
- ENSO: ~2-6 year time scale oscillations of convection, oceanic and atmospheric circulations. Wave effects extend into high latitudes and into tropical lower stratosphere.
- Both QBO and ENSO influence temperature, circulation and transport in UTLS

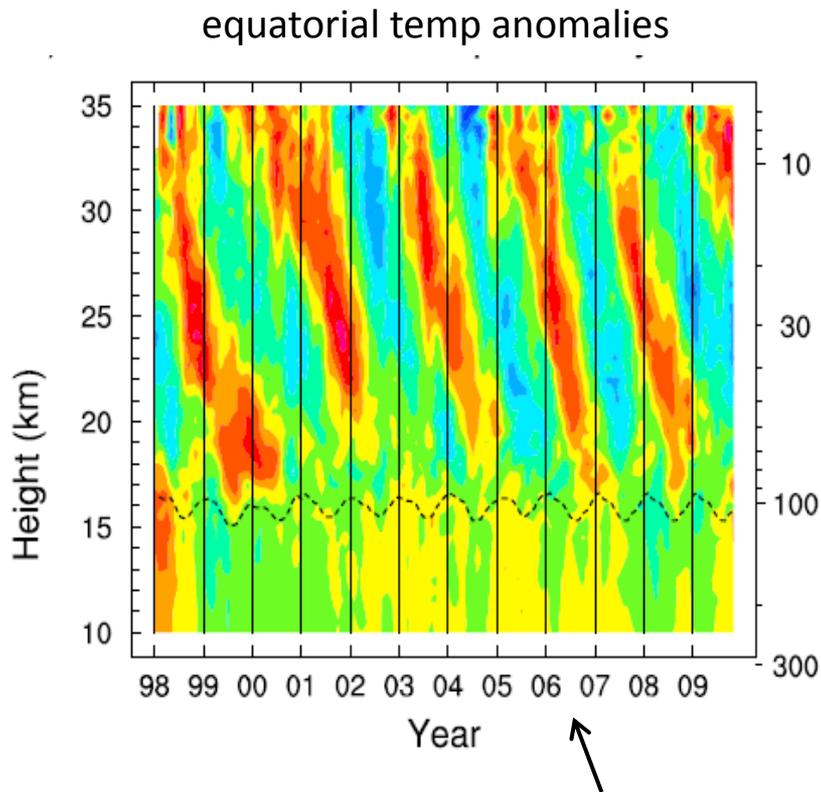
# Brief observational characteristics of QBO

easterly descent is regular;  
sometimes the westerly phase 'stalls'

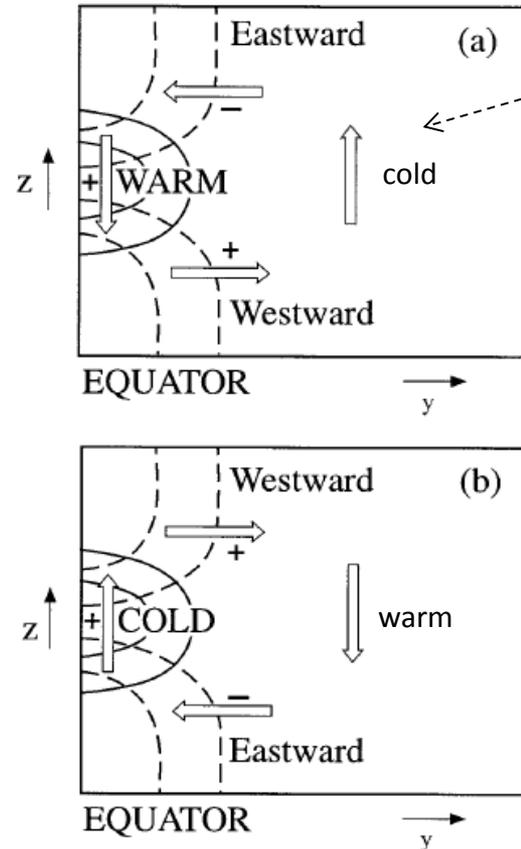
Zonal winds at Singapore (1° N) - standard QBO reference



# Meridional circulations and temperatures linked with the QBO



- temp anomalies at equator up to +/- 4 K
- extend downward to near tropopause

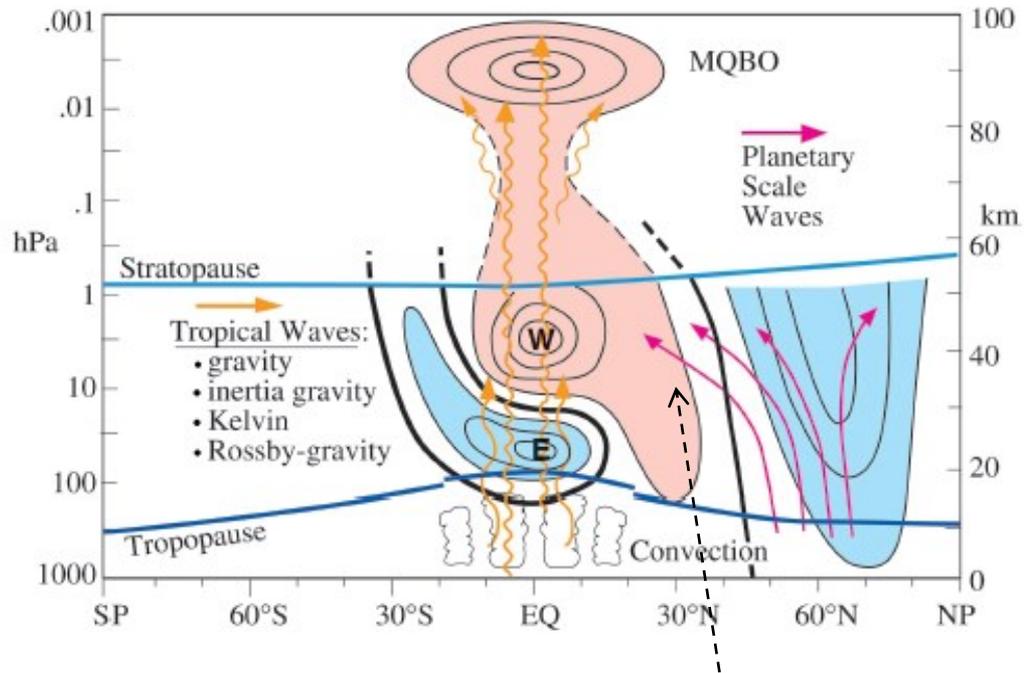


extension to middle latitudes

Plumb and Bell, 1982b, Q.J.R. Meteorol. Soc.

# Meridional structure and global influences of the QBO

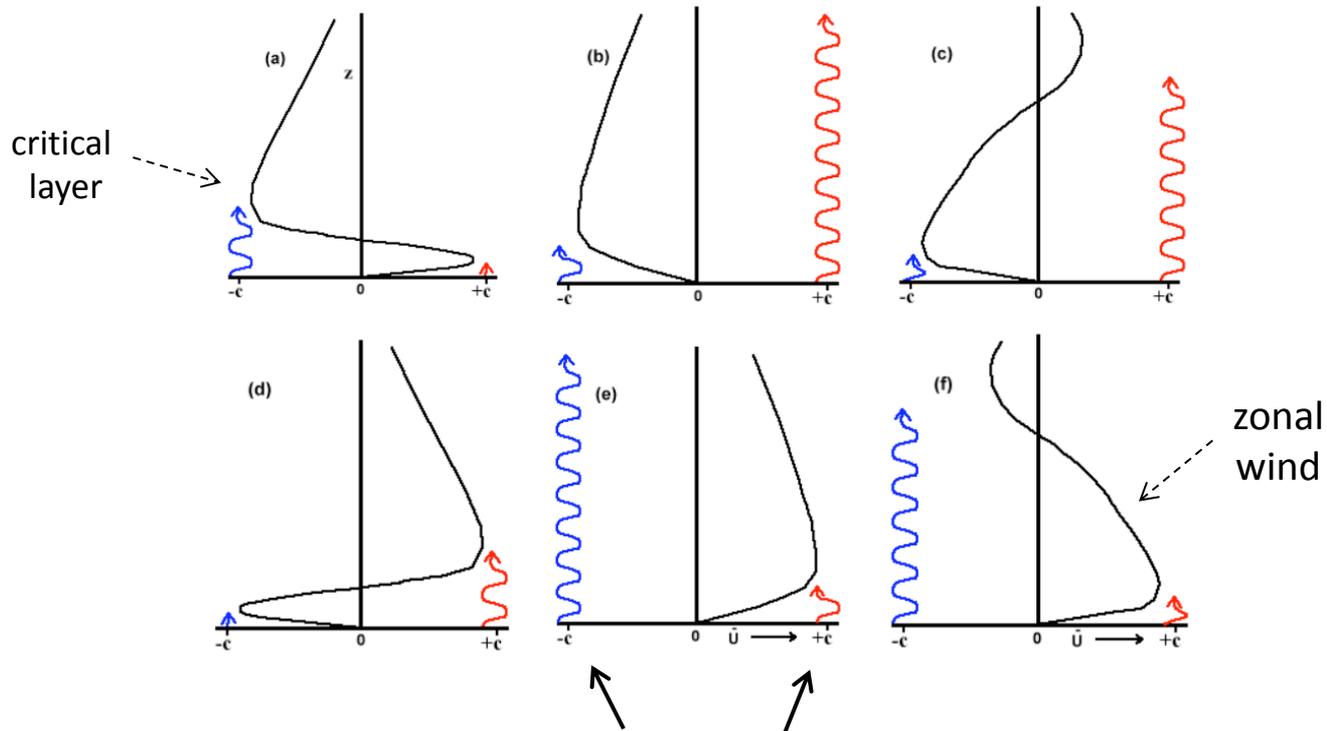
Baldwin et al 1998



influence on high latitudes  
via subtropical critical lines

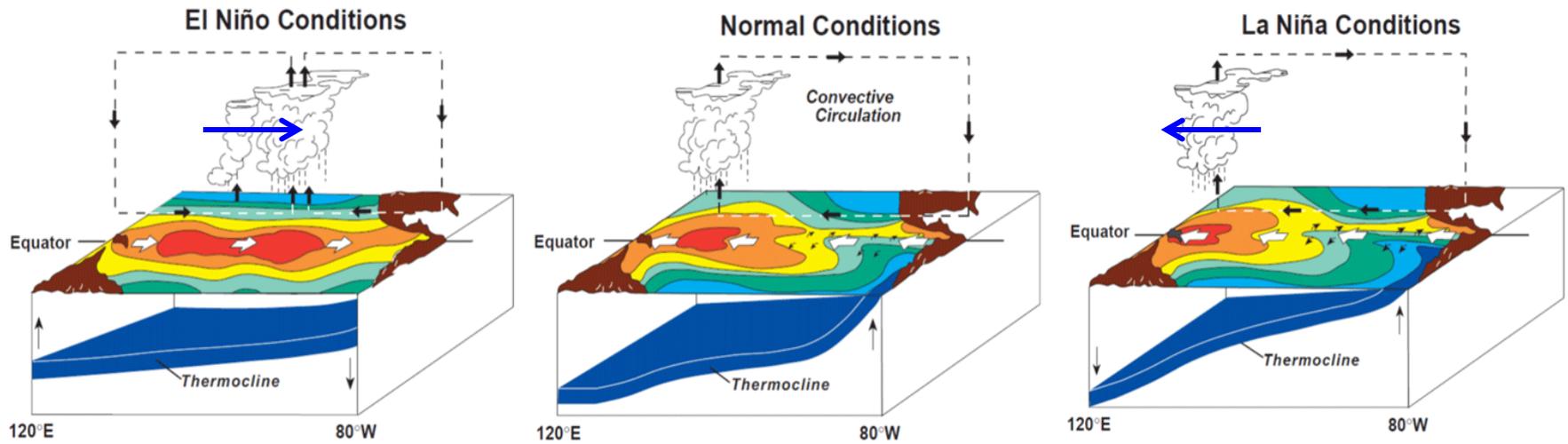
- waves propagate vertically until reaching critical level ( $U=c$ )
- wave absorption produces zonal acceleration (towards phase speed of wave)

Lindzen-Holton mechanism for the QBO

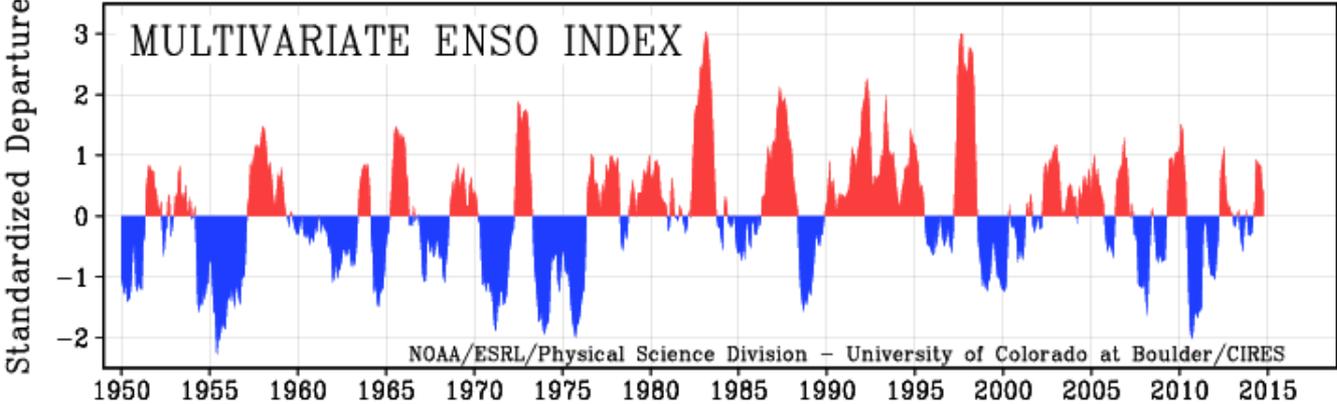


eastward and westward propagating waves forced from troposphere

# ENSO – large-scale shifts in tropical convection, winds, temperatures

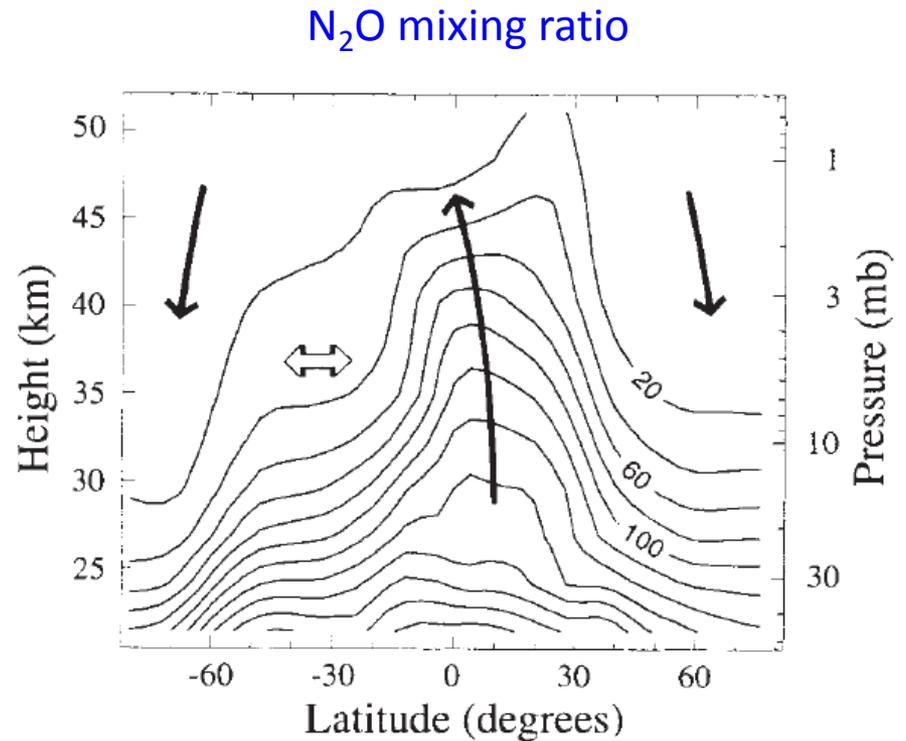


# Standard time series for ENSO

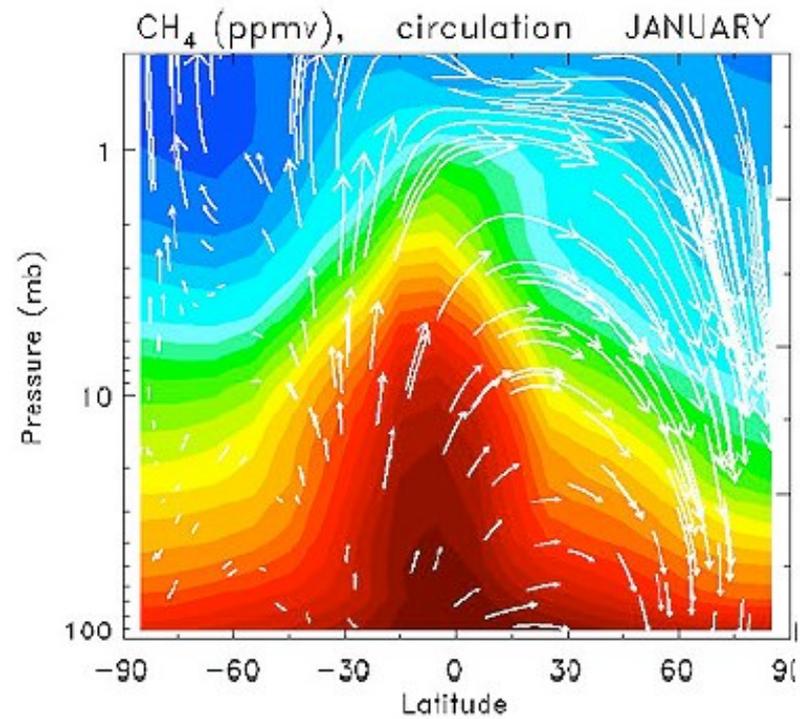


## Stratospheric tracer transport: evidence from satellite observations

- $N_2O$  is a 'tropospheric source gas'
- destroyed by photolysis (radiation) in upper stratosphere
- Source of reactive nitrogen ( $NO$ ,  $NO_2$ ) in upper stratosphere; important for stratospheric ozone
- Behavior reflects Brewer-Dobson circulation and eddy mixing

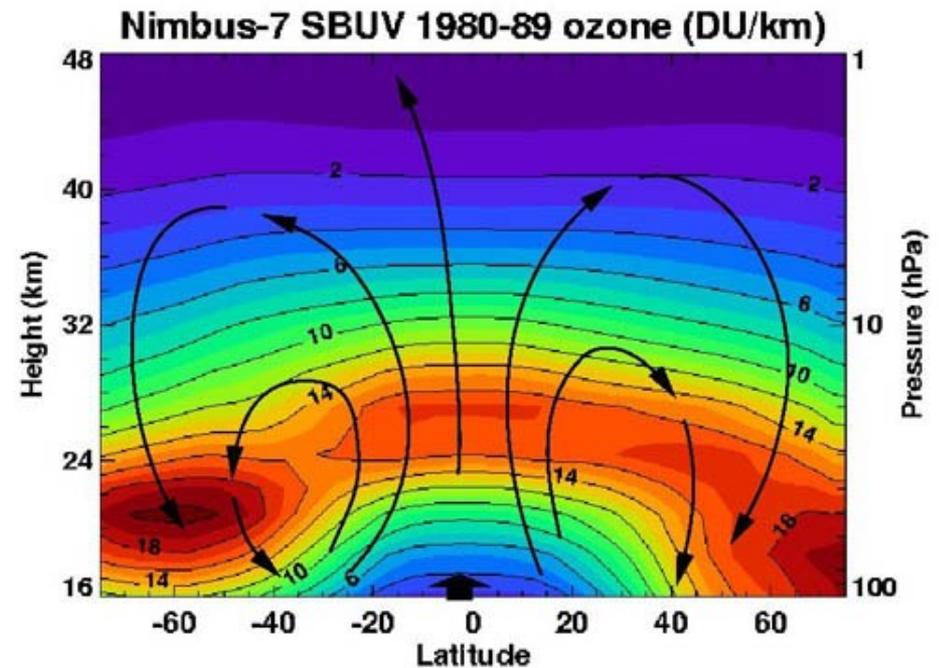


Methane  $\text{CH}_4$  is another tropospheric source gas, oxidized to  $\text{H}_2\text{O}$  in upper stratosphere

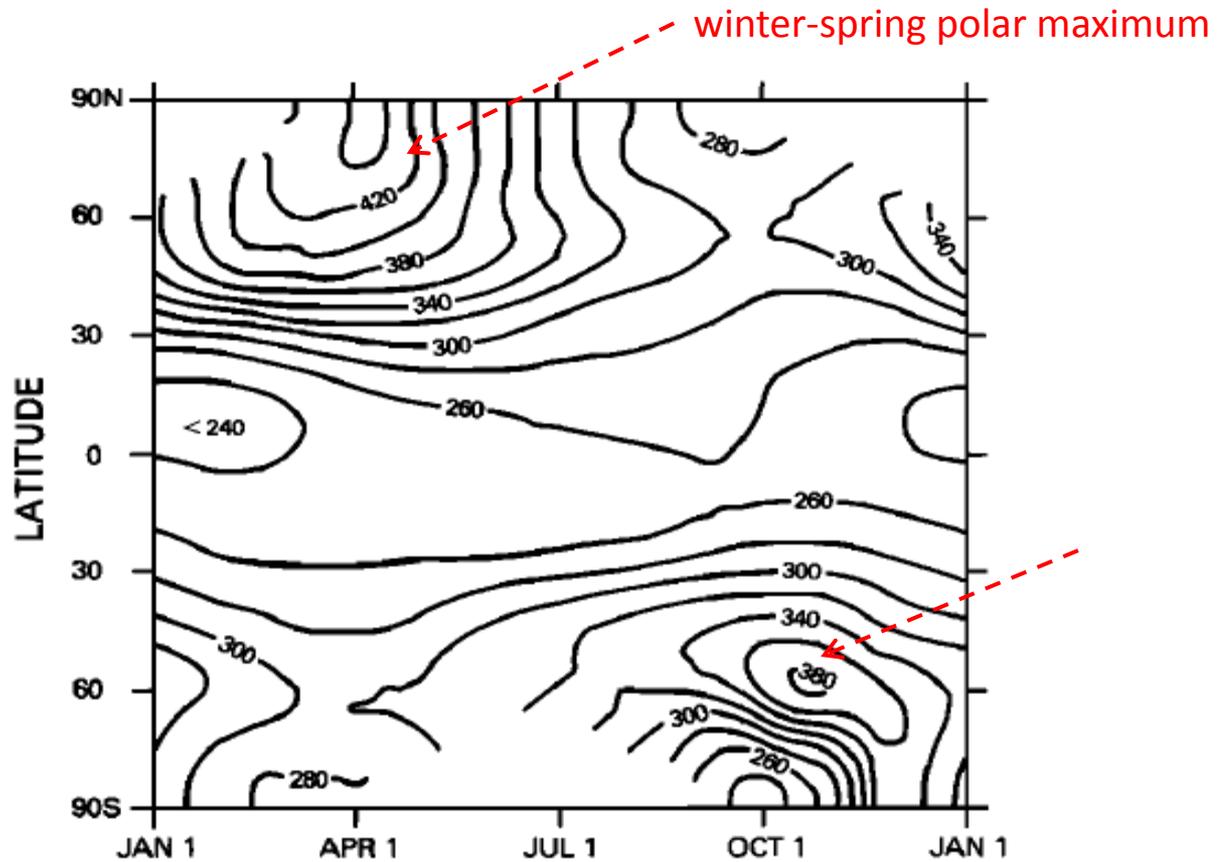


## Observed ozone and Brewer-Dobson circulation

- ozone is made in the tropical stratosphere
- Short lifetime in upper stratosphere
- Long lifetime in lower stratosphere
- transport causes high latitude maximum during winter / spring



Seasonal cycle of column ozone reflects Brewer-Dobson circulation

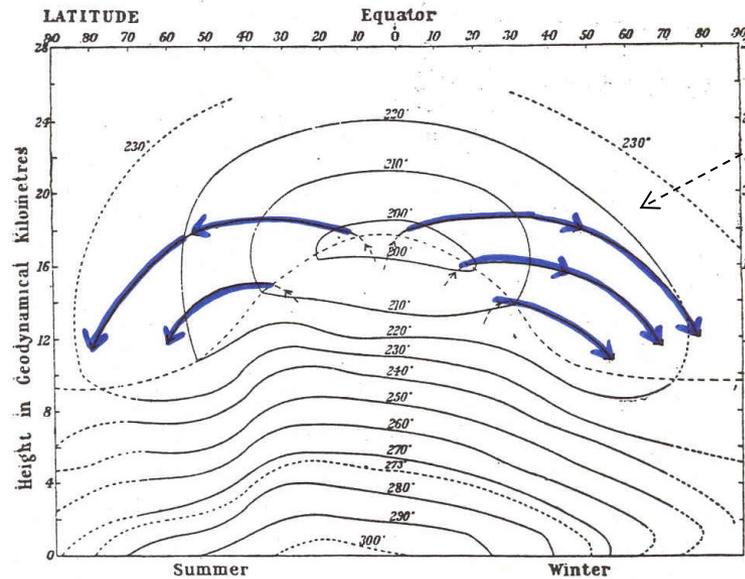


Bowman and Krueger, 1982

# EVIDENCE FOR A WORLD CIRCULATION PROVIDED BY MEASUREMENTS OF HELIUM AND WATER VAPOUR DISTRIBUTION IN THE STRATOSPHERE

QJRM, 1949

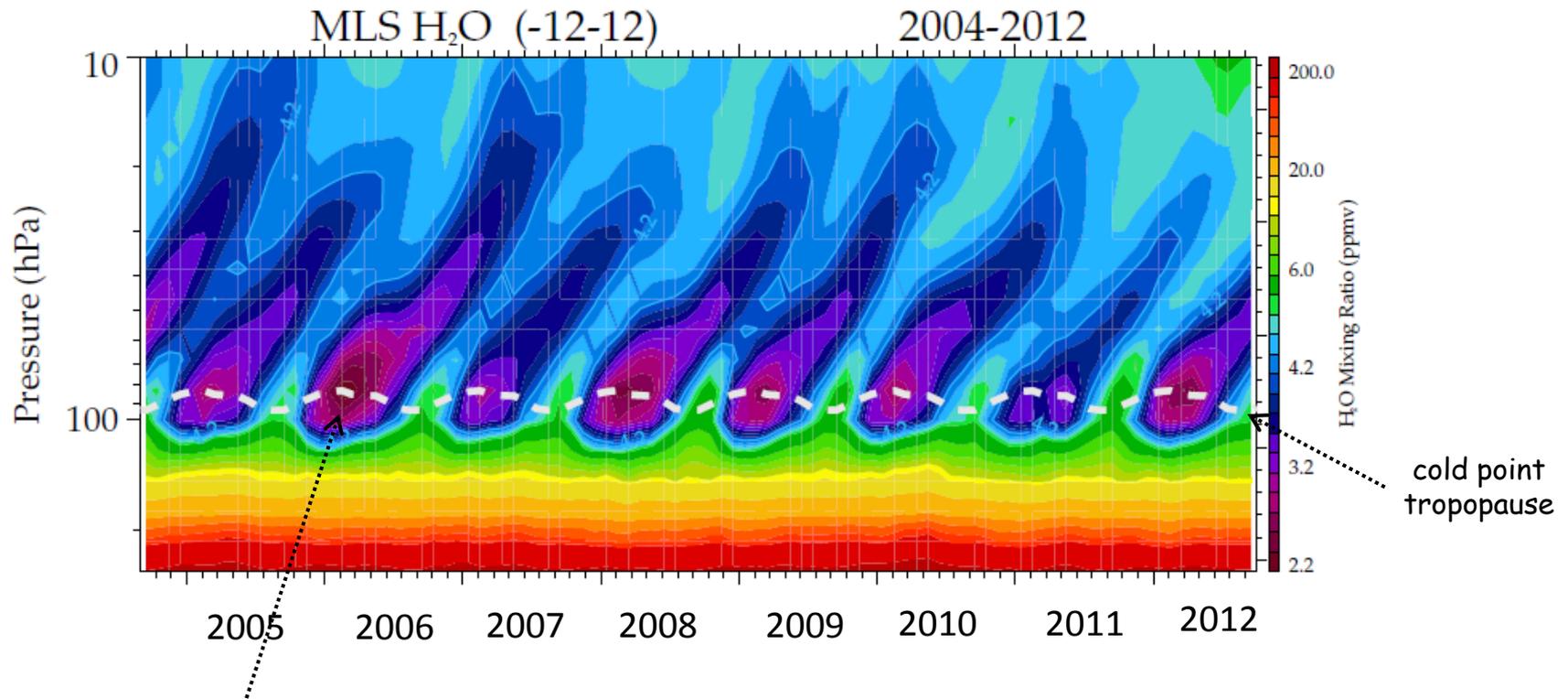
By A. W. BREWER, M.Sc., A.Inst.P.



The stratosphere is extremely dry because air is dehydrated passing the cold tropical tropopause

FIG. 5. Isotherms over the Globe  
A supply of dry air is maintained by a slow mean circulation from the equatorial tropopause.

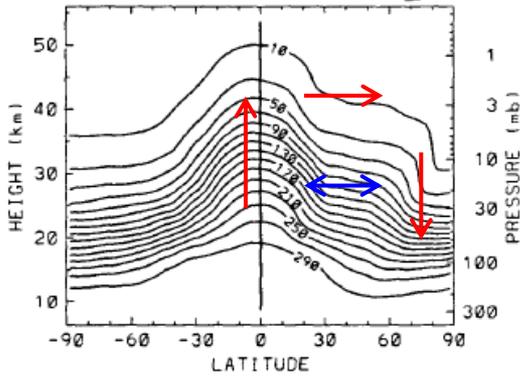
# Tropical tape recorder observed by MLS 2004-2012



- annual cycle in tropopause temperature imparts annual cycle in H<sub>2</sub>O
- upward propagation with Brewer-Dobson circulation

# tracer zonal mean transport budget

model N<sub>2</sub>O (ppbv)



mean advection

eddy transport

$$\frac{\partial \bar{\chi}}{\partial t} = -\bar{v}^* \frac{1}{a} \frac{\partial \bar{\chi}}{\partial \phi} - \bar{w}^* \frac{\partial \bar{\chi}}{\partial z} + \nabla \cdot \mathbf{M} + P - L$$

$$M_y = -e^{-z/H} \left( \overline{v' \chi'} - \frac{\overline{v' T'}}{S} \bar{\chi}_z \right)$$

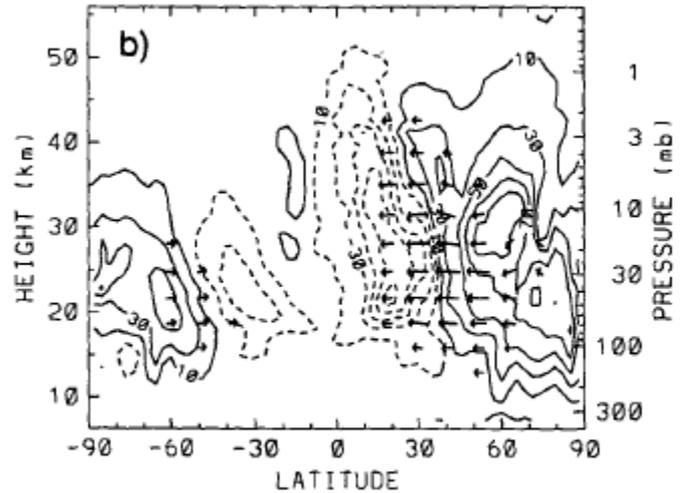
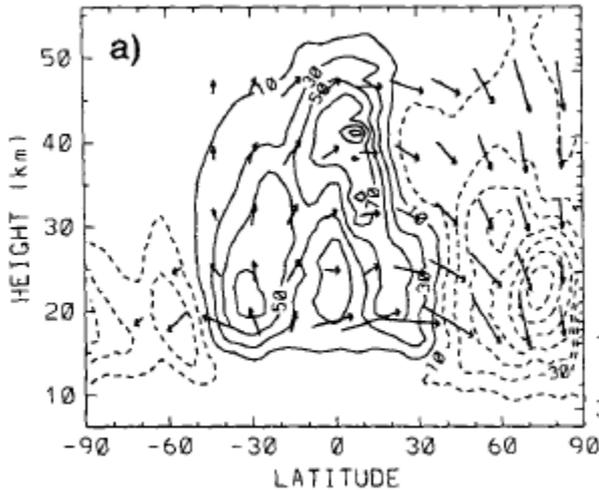
$$M_z = -e^{-z/H} \left( \overline{w' \chi'} + \frac{\overline{v' T'}}{S} \bar{\chi}_y \right)$$

these terms generally balance each other

mean advection

eddy transport

contours:  
N<sub>2</sub>O  
tendency  
(ppbv/100 days)



Tracer transport equation similar to thermodynamic equation:

tracer

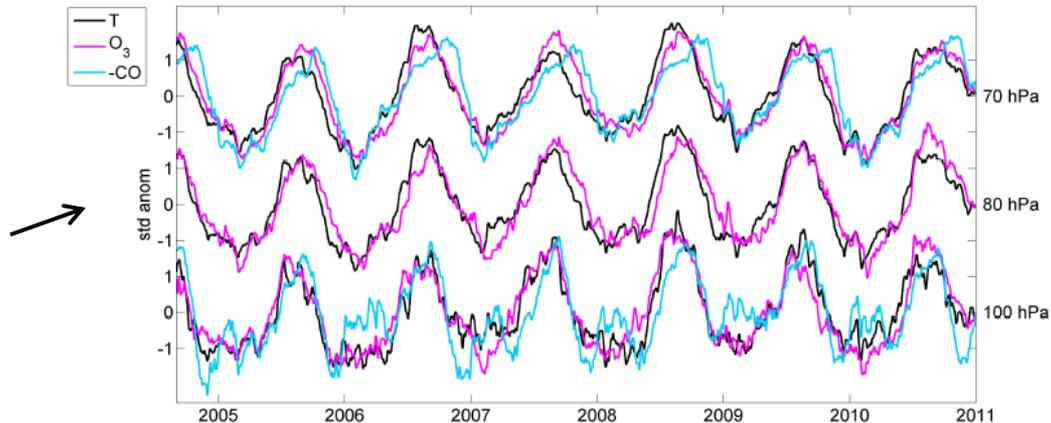
$$\frac{\partial \bar{\chi}}{\partial t} = -\bar{v}^* \frac{1}{a} \frac{\partial \bar{\chi}}{\partial \phi} - \bar{w}^* \frac{\partial \bar{\chi}}{\partial z} + \nabla \cdot \mathbf{M} + P - L$$

temperature

$$\frac{\partial \bar{T}}{\partial t} + \bar{v}^* \frac{1}{a} \frac{\partial \bar{T}}{\partial \phi} + \bar{w}^* S = \bar{Q},$$

This is why temperature and tracers are sometimes highly correlated:

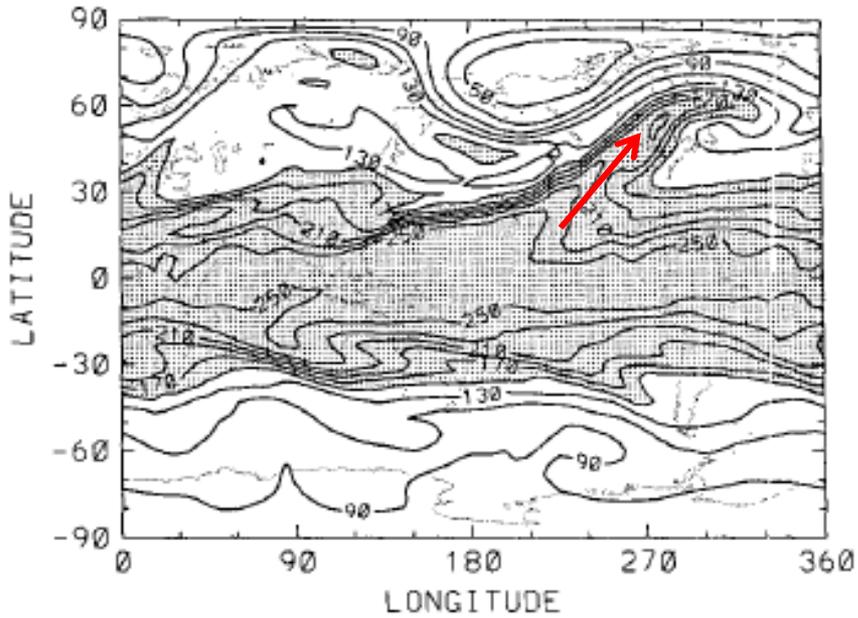
for example,  
T, O<sub>3</sub> and CO  
in tropical  
stratosphere  
(Abalos et al 2012)



# Examples of stratospheric wave mixing

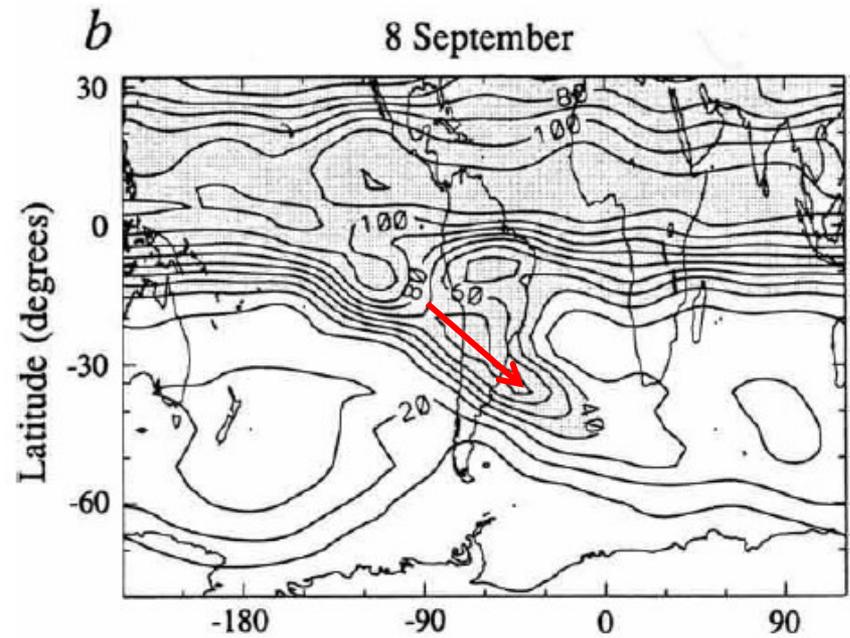
**model**

CCM2 N<sub>2</sub>O February 20 30 mb



N<sub>2</sub>O near 35 km from CLAES instrument on UARS

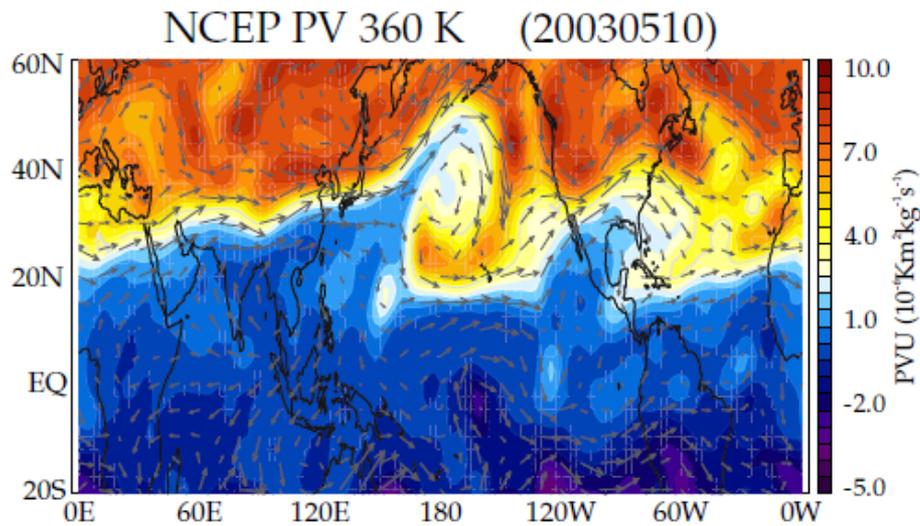
**observations**



Randel et al 1993

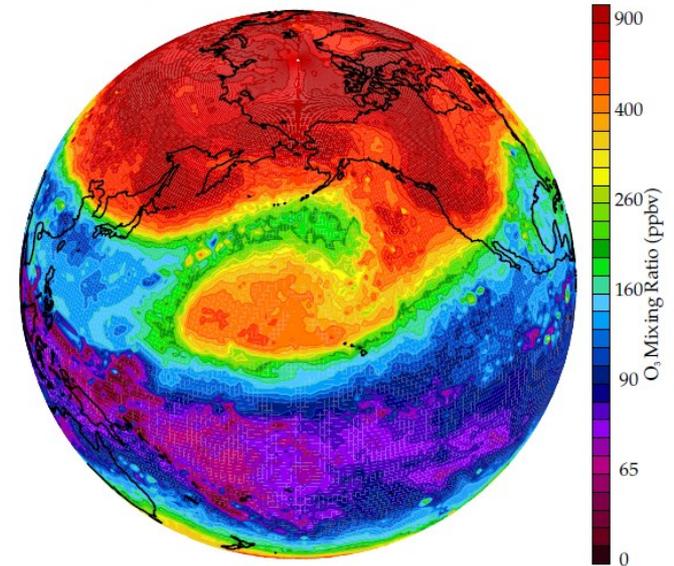
# Mixing across tropopause linked to Rossby wave breaking

## potential vorticity



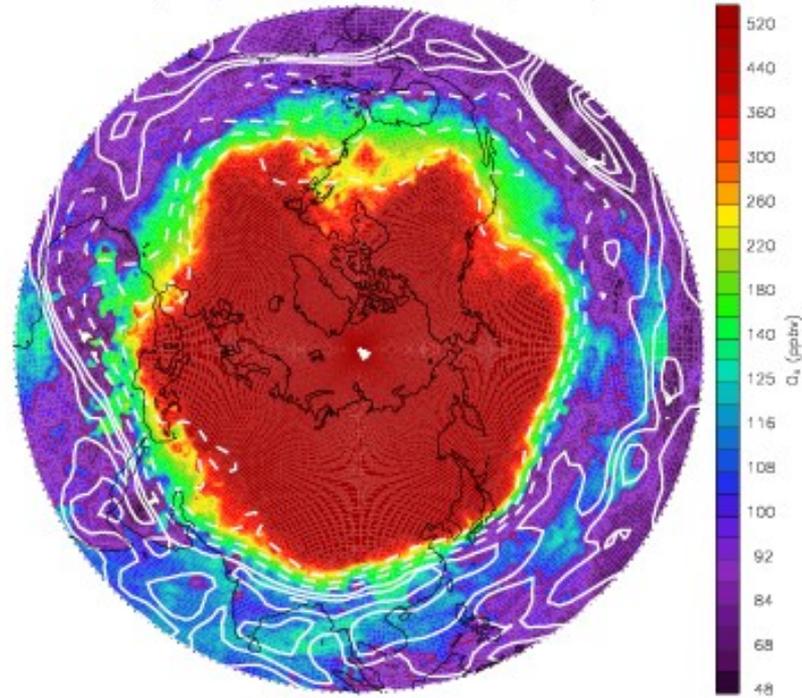
## ozone derived from AIRS

AIRS O<sub>3</sub> 360K (MAY/10/2003)



# Rossby wave variability reflected in ozone near tropopause

AIRS O<sub>3</sub> (NH) at 360K (JAN/01/2003)



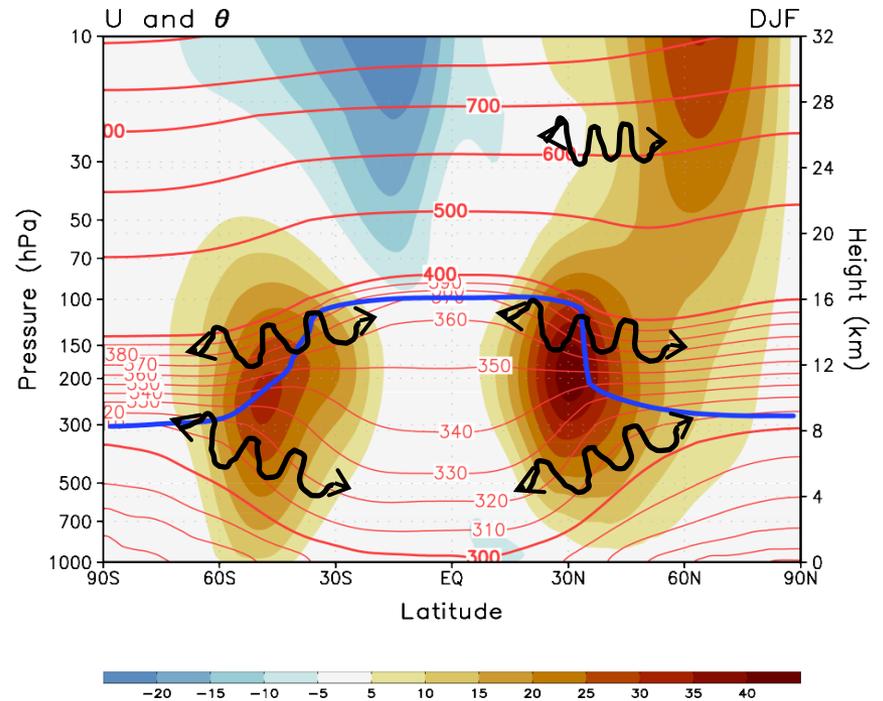
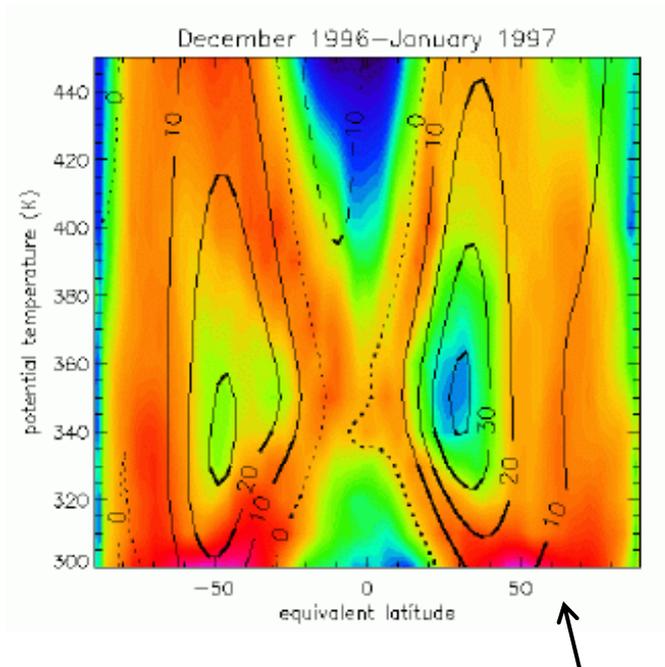
# Effective diffusivity as a diagnostic of atmospheric transport

## 2. Troposphere and lower stratosphere

JGR 2000

Peter Haynes and Emily Shuckburgh

Estimates of mixing based on stretching of PV contours in trajectory calculations



important points:

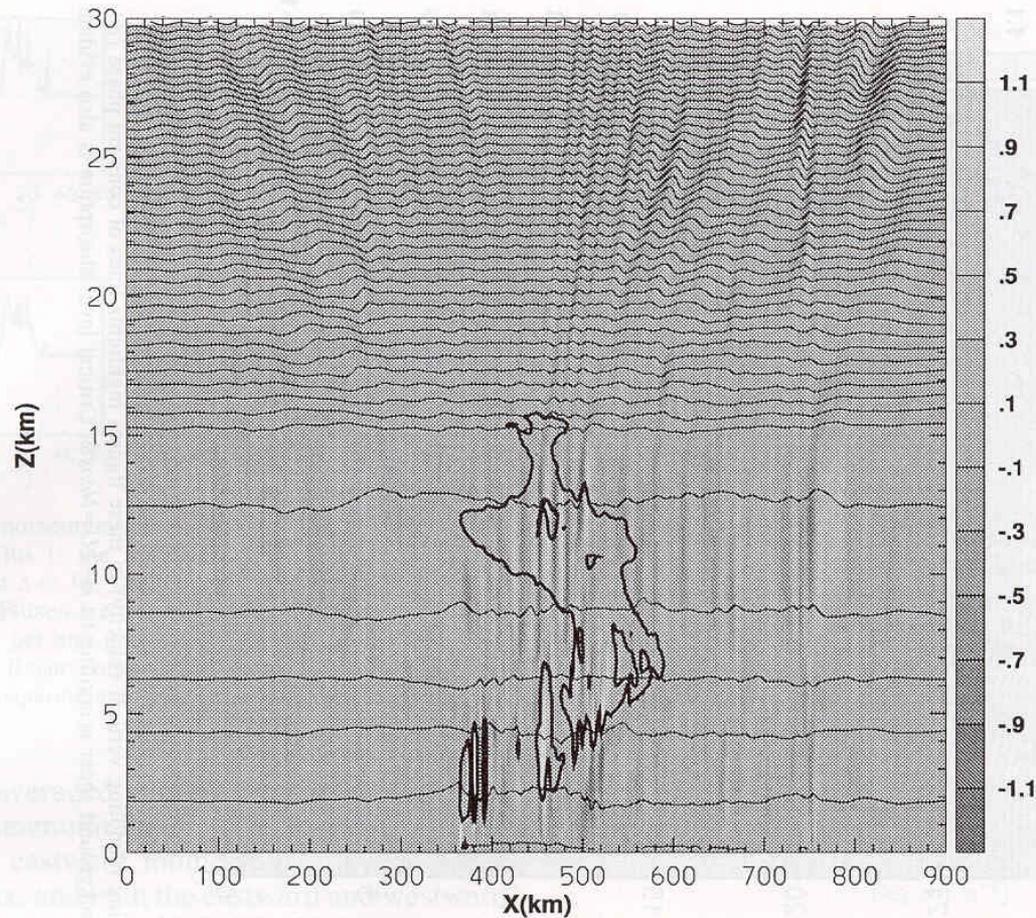
- mixing on flanks of jet (near critical lines for  $c \sim 10$  m/s)
- small mixing across jet core (jet cores are mixing barriers)

## Key points:

- Stratospheric transport: Brewer-Dobson circulation and wave mixing (clear behavior for tropospheric source gases)
- Stratospheric ozone: produced in tropical stratosphere, transported to high latitudes (reflects seasonal Brewer-Dobson circulation)
- Stratospheric water vapor: dehydration near tropical cold point, strong seasonal cycle ('tape recorder')
- Tracer budgets: mean advection and eddy transports (tied to Rossby waves and critical layers)

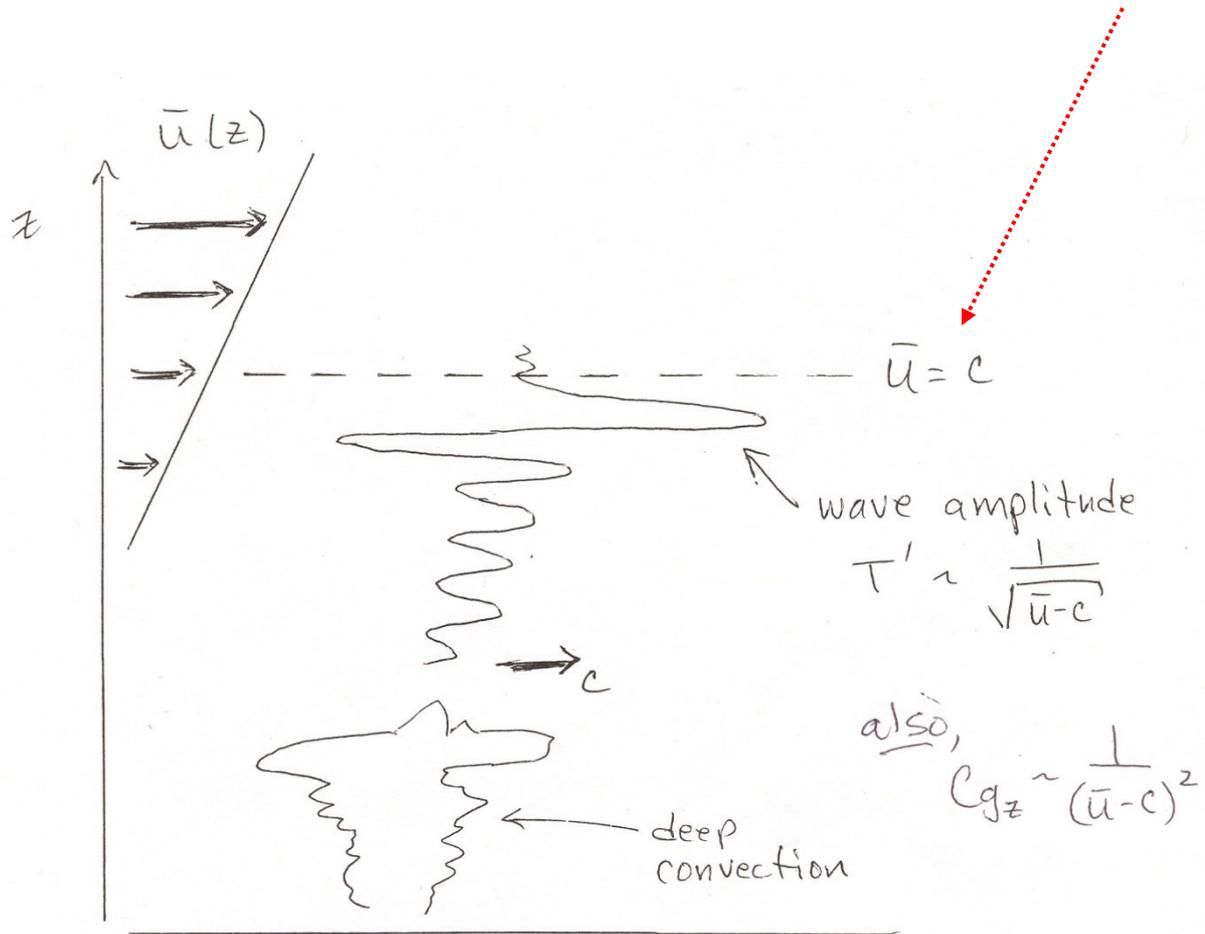
Extra slides

# Model simulation of gravity waves forced by deep convection

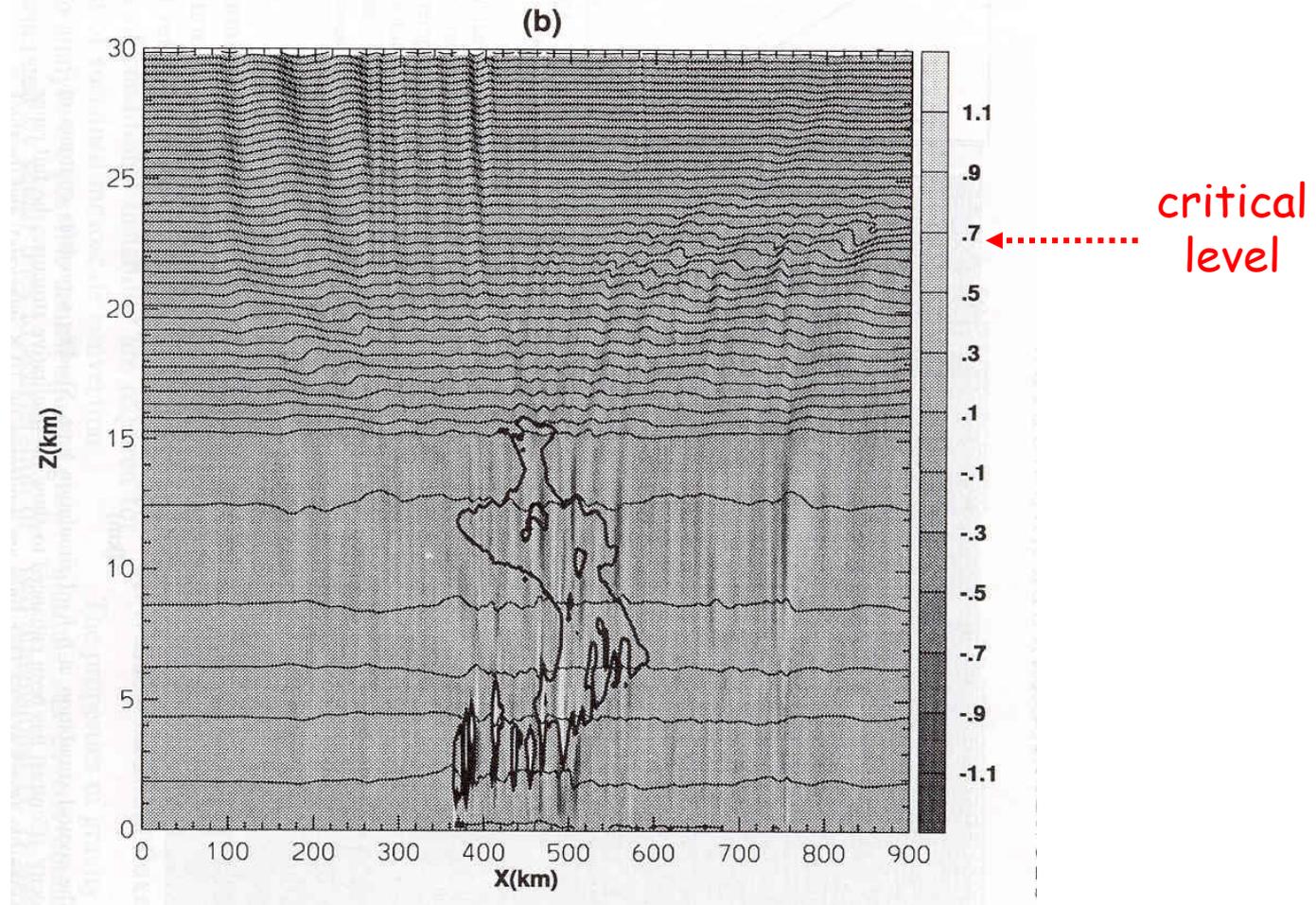


Alexander and Holton, 2000

# Gravity waves interacting with a **critical level**



# Gravity waves interacting with a critical level



# Climatology of Intrusions into the Tropical Upper Troposphere

Darryn W. Waugh

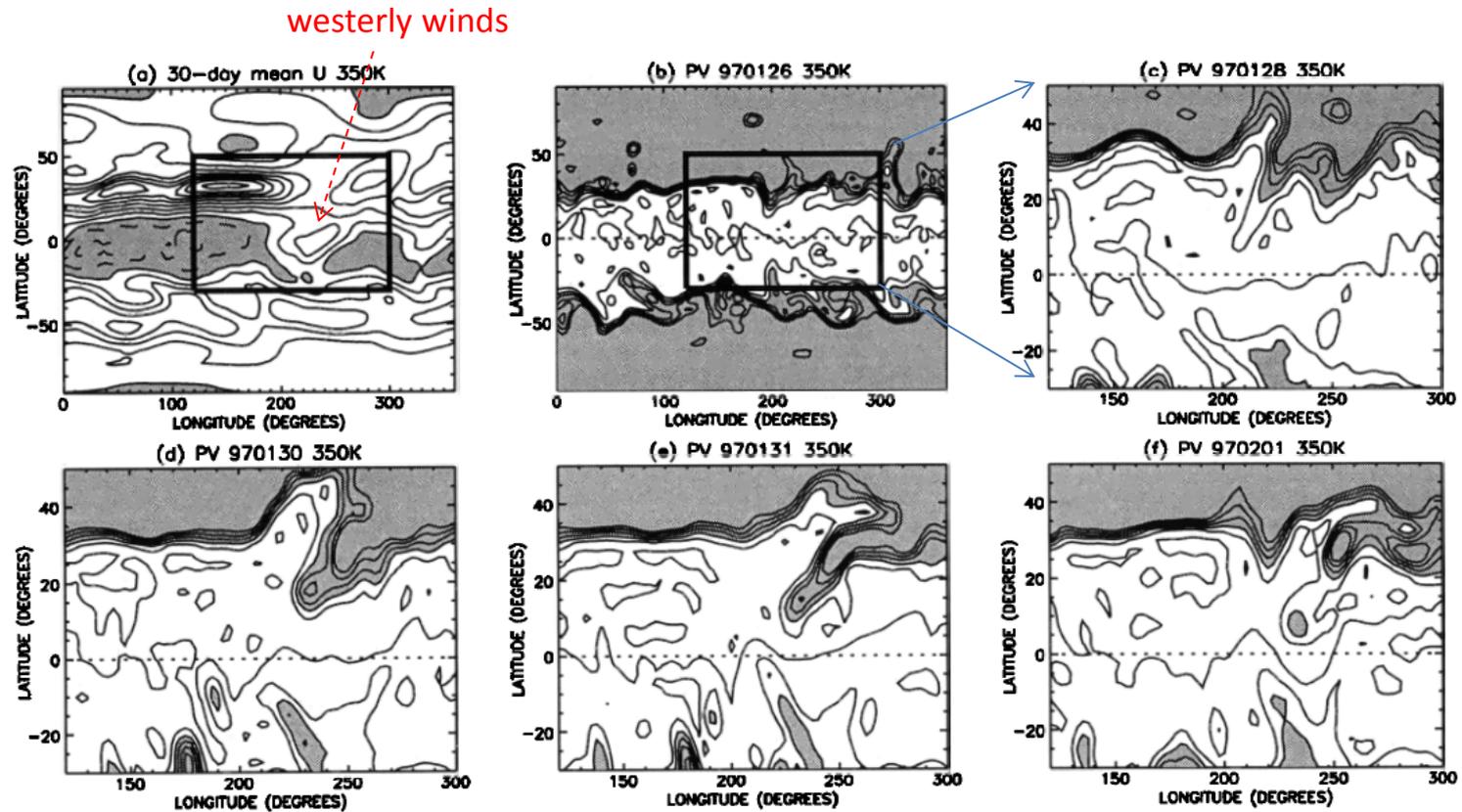
Department of Earth and Planetary Science, Johns Hopkins University, Baltimore, MD.

Lorenzo M. Polvani

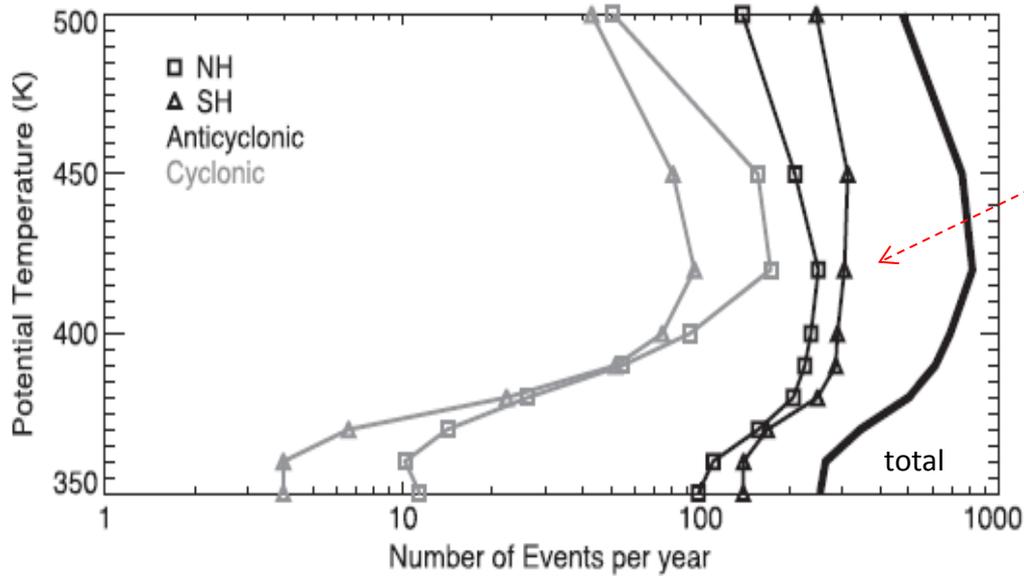
Department of Applied Physics and Applied Mathematics, Columbia University, New York, NY.

Latitudinal propagation depends on background zonal winds:

- Rossby wave propagation through westerly wind 'ducts'



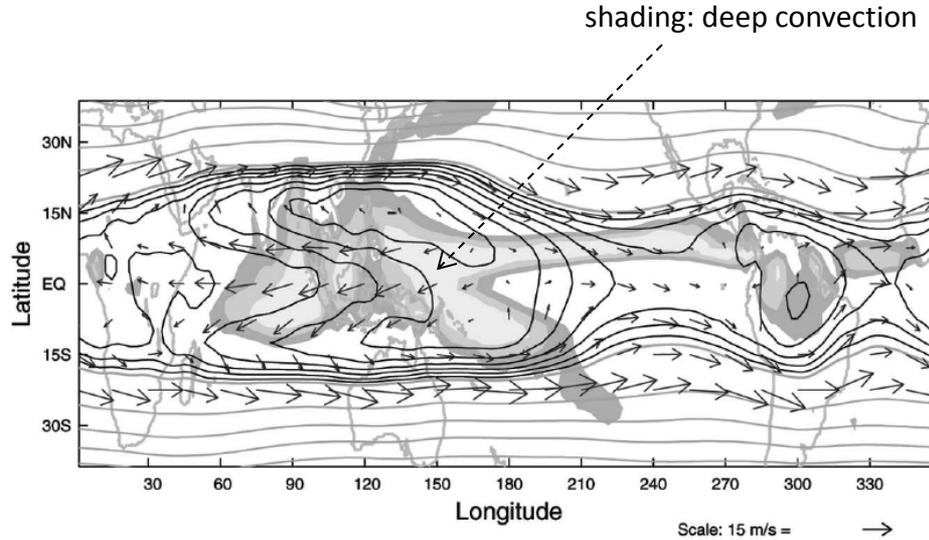
# climatology of Rossby wave breaking events



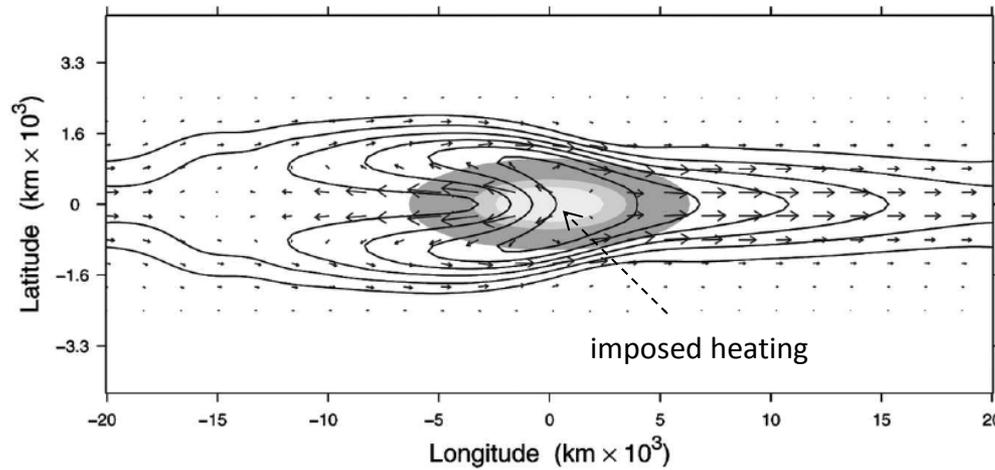
Homeyer and Bowman 2013

# observations vs. model

upper troposphere  
geopotential height  
and winds



annual mean  
observations



nonlinear  
shallow water  
model

Dima et al 2005

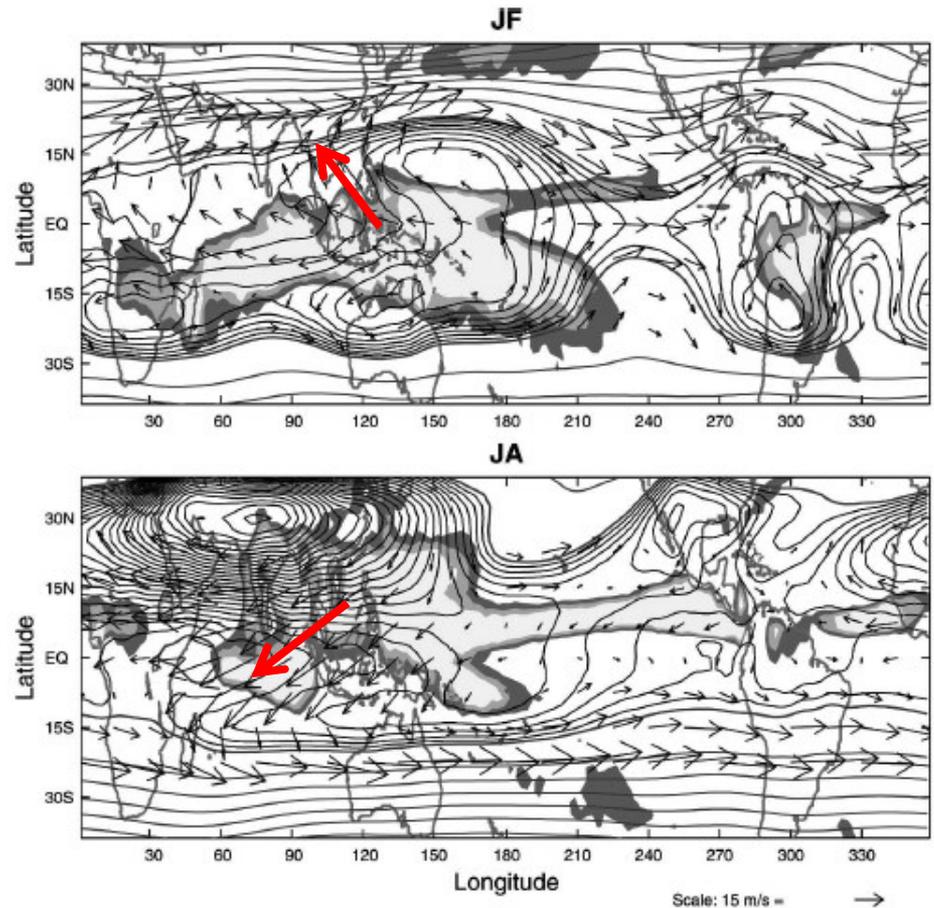
# Seasonal variation of tropical waves

tropical zonal mean  
momentum balance:

$$\frac{\partial[u]}{\partial t} \equiv [v] \left( f - \frac{1}{\cos\phi} \frac{\partial[u] \cos\phi}{\partial y} \right) - [\omega] \frac{\partial[u]}{\partial p}$$

$$- \frac{1}{\cos^2\phi} \frac{\partial[u^*v^*] \cos^2\phi}{\partial y} - \frac{\partial[u^*\omega^*]}{\partial p} - [F_x].$$

eddy momentum  
fluxes balance  
Hadley flow

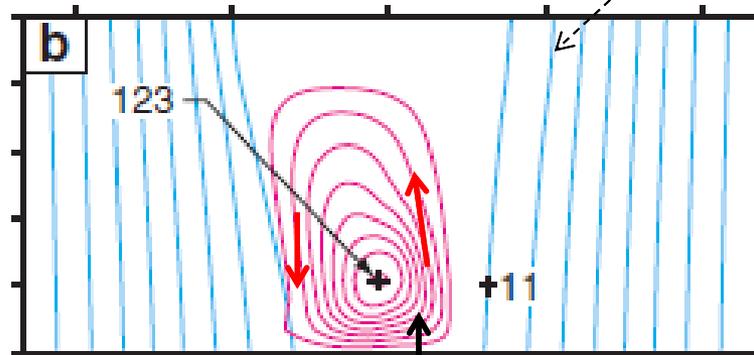


Dima et al 2005

# Hadley cell interactions with extratropical eddies: (a complicated subject)

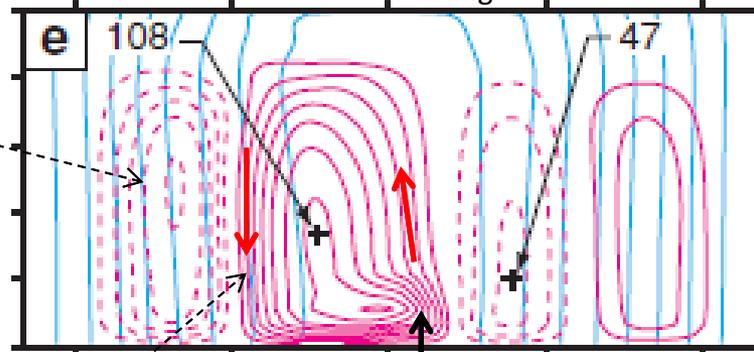
Idealized model simulations:

angular momentum contours



**zonally symmetric dynamics**

heating



**include extratropical eddies**

midlatitude  
'Ferrel cell' linked  
to midlatitude eddies

extension of Hadley  
cell to higher latitudes

**Latitude**

Schneider 2006

# tracer zonal mean transport budget

eddy transport

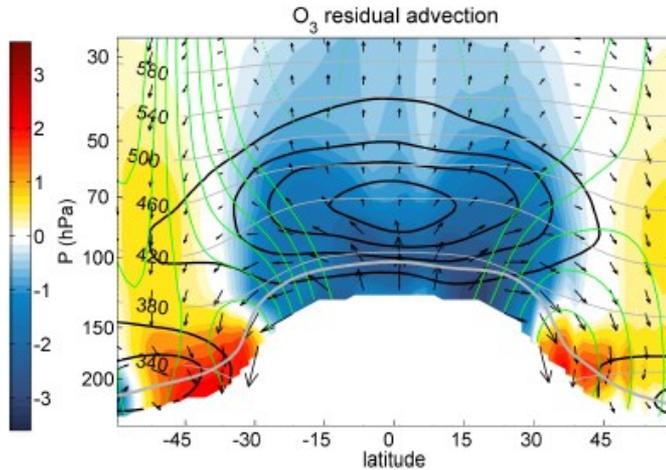
$$\frac{\partial \bar{\chi}}{\partial t} = -\bar{v}^* \frac{1}{a} \frac{\partial \bar{\chi}}{\partial \phi} - \bar{w}^* \frac{\partial \bar{\chi}}{\partial z} + \nabla \cdot \mathbf{M} + P - L$$

tropical ozone budget  
from WACCM

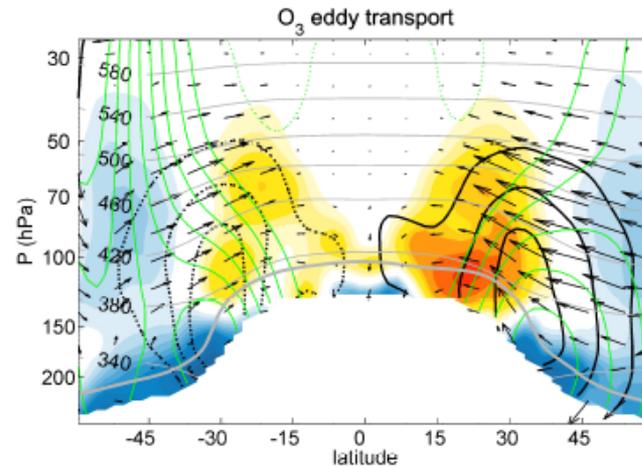
$$M_y = -e^{-z/H} \left( \overline{v' \chi'} - \frac{\overline{v' T'}}{S} \bar{\chi}_z \right)$$

$$M_z = -e^{-z/H} \left( \overline{w' \chi'} + \frac{\overline{v' T'}}{S} \bar{\chi}_y \right)$$

mean advection



eddy transport



ozone  
tendency  
in % per day

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